

The background of the slide is a deep space image filled with numerous galaxies and stars. The galaxies are of various shapes and sizes, some appearing as bright, fuzzy clouds, while others are more distant and smaller. The stars are scattered across the dark field, with some showing prominent diffraction spikes. The overall color palette is dominated by blacks and greys, with highlights of white, yellow, and blue from the celestial objects.

Glimpsing the

Glimpsing the
Big Bang?

Martin Hendry
University of Glasgow

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Gravitational waves: have US scientists heard echoes of the big bang?

Discovery of gravitational waves by Bicep telescope at south pole could give scientists insights into how universe was born

Stuart Clark


The Guardian, Friday 14 March 2014 15:33 GMT

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Primordial gravitational waves would provide evidence of inflation in the moments after the big bang. Photograph: Alamy

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BICEP2 I: DETECTION OF *B*-mode POLARIZATION AT DEGREE ANGULAR SCALES

BICEP2 COLLABORATION - P. A. R. ADE¹, R. W. AIKIN², D. BARKATS³, S. J. BENTON⁴, C. A. BISCHOFF⁵, J. J. BOCK^{2,6}, J. A. BREVIK², I. BUDER⁵, E. BULLOCK⁷, C. D. DOWELL⁶, L. DUBAND⁸, J. P. FILIPPINI², S. FLIESCHER⁹, S. R. GOLWALA², M. HALPERN¹⁰, M. HASSELFIELD¹⁰, S. R. HILDEBRANDT^{2,6}, G. C. HILTON¹¹, V. V. HRISTOV², K. D. IRWIN^{12,13,11}, K. S. KARKARE⁵, J. P. KAUFMAN¹⁴, B. G. KEATING¹⁴, S. A. KERNASOVSKIY¹², J. M. KOVAC^{5,16}, C. L. KUO^{12,13}, E. M. LEITCH¹⁵, M. LUEKER², P. MASON², C. B. NETTERFIELD⁴, H. T. NGUYEN⁶, R. O'BRIENT⁶, R. W. OGBURN IV^{12,13}, A. ORLANDO¹⁴, C. PRYKE^{9,7,16}, C. D. REINTSEMA¹¹, S. RICHTER⁵, R. SCHWARZ⁹, C. D. SHEEHY^{9,15}, Z. K. STANISZEWSKI^{2,6}, R. V. SUDIWALA¹, G. P. TEPLY², J. E. TOLAN¹², A. D. TURNER⁶, A. G. VIEREGG^{5,15}, C. L. WONG⁵, AND K. W. YOON^{12,13}

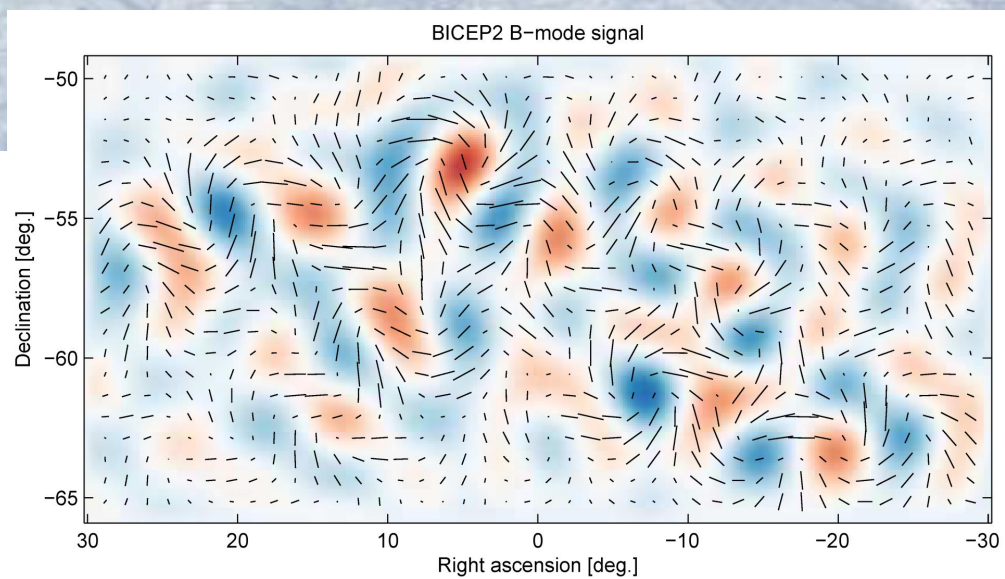
to be submitted to a journal TBD

ABSTRACT

We report results from the BICEP2 experiment, a Cosmic Microwave Background (CMB) polarimeter specifically designed to search for the signal of inflationary gravitational waves in the *B*-mode power spectrum around $\ell \sim 80$. The telescope comprised a 26 cm aperture all-cold refracting optical system equipped with a focal plane of 512 antenna coupled transition edge sensor (TES) 150 GHz bolometers each with temperature sensitivity of $\approx 300 \mu\text{K}_{\text{CMB}} \sqrt{\text{s}}$. BICEP2 observed from the South Pole for three seasons from 2010 to 2012. A low-foreground region of sky with an effective area of 380 square degrees was observed to a depth of 87 nK-degrees in Stokes *Q* and *U*. In this paper we describe the observations, data reduction, maps, simulations and results. We find an excess of *B*-mode power over the base lensed- Λ CDM expectation in the range $30 < \ell < 150$, inconsistent with the null hypothesis at a significance of $> 5\sigma$. Through jackknife tests and simulations based on detailed calibration measurements we show that systematic contamination is much smaller than the observed excess. We also estimate potential foreground signals and find that available models predict these to be considerably smaller than the observed signal. These foreground models possess no significant cross-correlation with our maps. Additionally, cross-correlating BICEP2 against 100 GHz maps from the BICEP1 experiment, the excess signal is confirmed with 3σ significance and its spectral index is found to be consistent with that of the CMB, disfavoring synchrotron or dust at 2.3σ and 2.2σ , respectively. The observed *B*-mode power spectrum is well-fit by a lensed- Λ CDM + tensor theoretical model with tensor/scalar ratio $r = 0.20^{+0.07}_{-0.05}$, with $r = 0$ disfavored at 7.0σ . Subtracting the best available estimate for foreground dust modifies the likelihood slightly so that $r = 0$ is disfavored at 5.9σ .

Subject headings: cosmic background radiation — cosmology: observations — gravitational waves — inflation — polarization

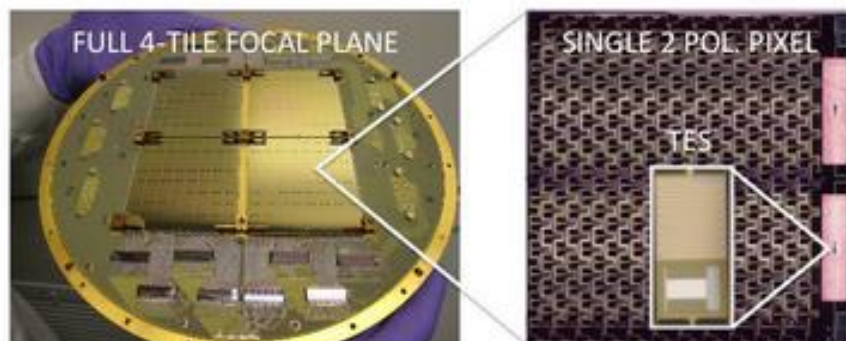




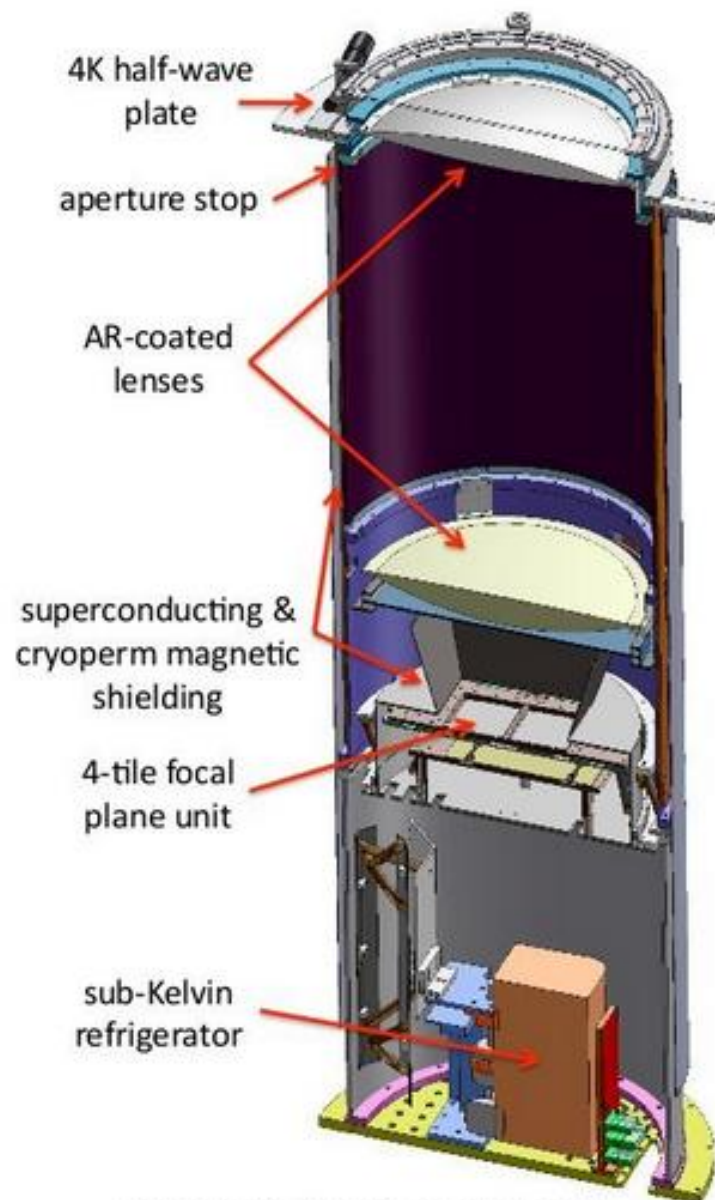
Instrumentation

BICEP2 observes on degree angular scales deep into the same patch of sky as BICEP1, employing cutting-edge technology to pack many more detectors onto the same size focal plane as BICEP1, thus dramatically increasing sensitivity. BICEP2 aims to measure the polarization of the cosmic microwave background (CMB) with better sensitivity than ever before.

The BICEP2 telescope is a small aperture, refracting telescope, allowing for precise control of systematic effects. The telescope optics are cooled to 4K. BICEP2 relies on the same principles as BICEP1, but with new detectors which allow for closer packing onto the focal plane. BICEP2 has 512 antenna-coupled TES bolometers in the focal plane. The focal plane is kept at 250 mK for CMB observations. Gaining detector sensitivity is critical in the search for "B-mode" polarization, and BICEP2 has an order of magnitude greater mapping speed than BICEP1.



The 512-detector BICEP2 focal plane and a zoom in of a dual polarization pixel with TES bolometer shown.

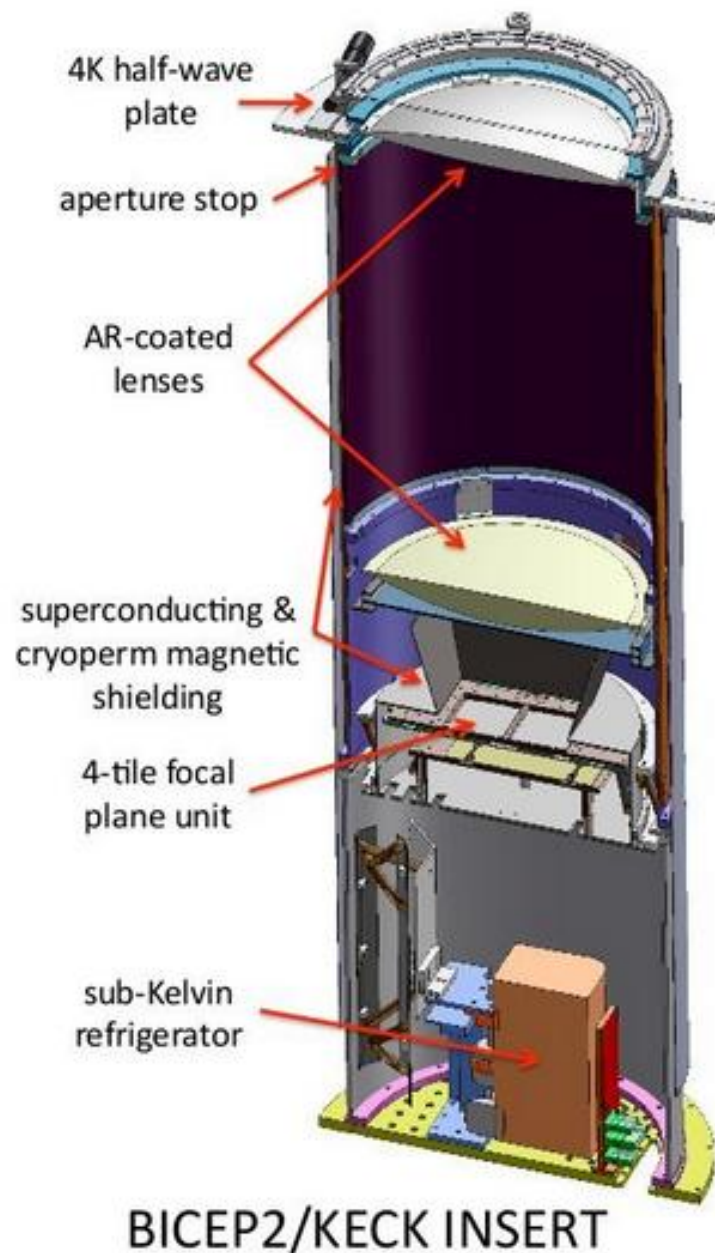
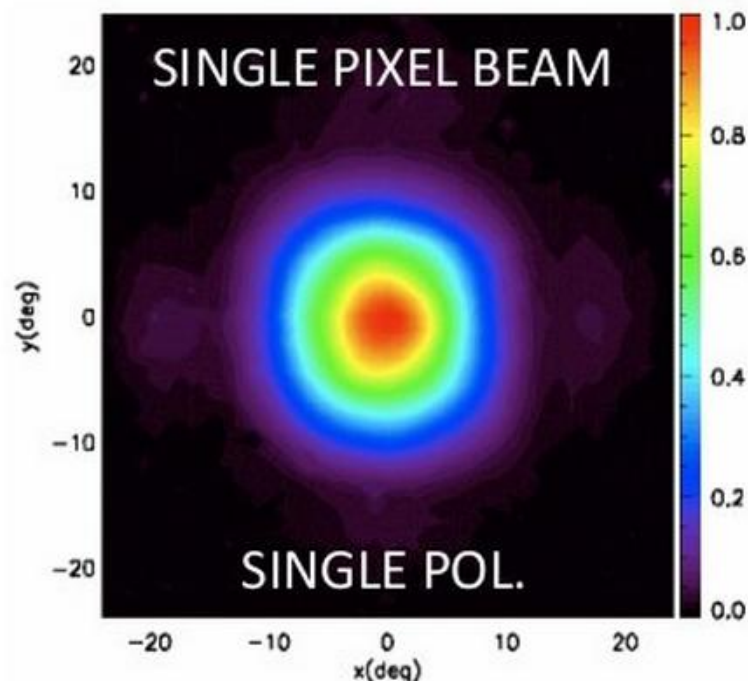


BICEP2/KECK INSERT

Instrumentation

BICEP2 observes on degree angular scales deep into the same patch of sky as BICEP1, employing cutting-edge technology to pack many more detectors onto the same size focal plane as BICEP1, thus dramatically increasing sensitivity. BICEP2 aims to measure the polarization of the cosmic microwave background (CMB) with better sensitivity than ever before.

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Primordial gravitational wave discovery heralds 'whole new era' in physics

Gravitational waves could help unite general relativity and quantum mechanics to reveal a 'theory of everything'

Stuart Clark

The Guardian, Monday 17 March 2014 18.08 GMT

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Scientists detected telltale signs of gravitational waves using the Bicep2 telescope (far left) at the south pole. Photograph: Keith Vanderlinde/NSF

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Gravitational waves discovery: 'We have a first tantalising glimpse of the cosmic birth'



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Former Doves frontman Jimi Goodwin performs a live version of Terracotta Warrior at the Guardian's

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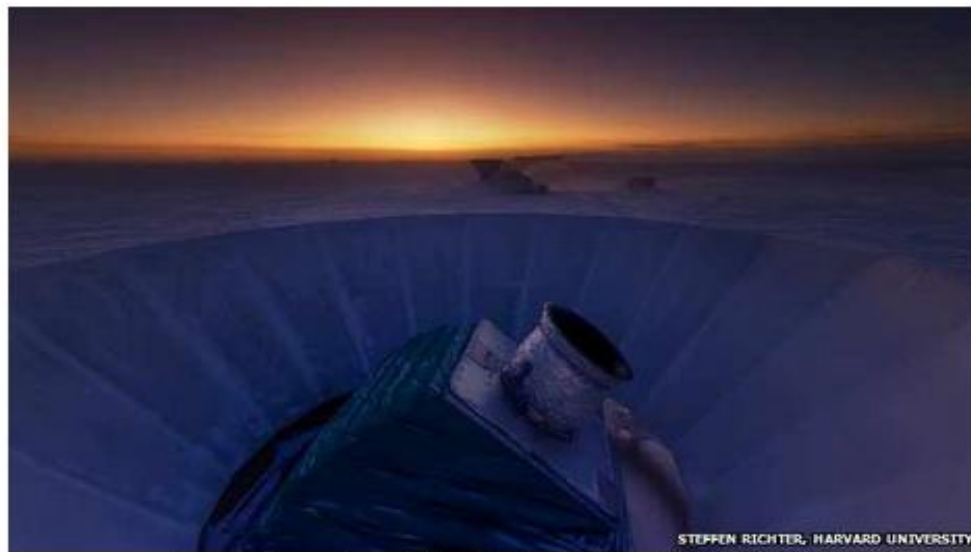
17 March 2014 Last updated at 14:46



Cosmic inflation: 'Spectacular' discovery hailed

By Jonathan Amos

Science correspondent, BBC News



STEFFEN RICHTER, HARVARD UNIVERSITY

The measurements were taken using the BICEP2 instrument at the South Pole Telescope facility

Scientists say they have extraordinary new evidence to support a Big Bang Theory for the origin of the Universe.

Researchers believe they have found the signal left in the sky by the super-rapid expansion of space that must have occurred just fractions of a second after everything came into being.

It takes the form of a distinctive twist in the oldest light detectable with telescopes.

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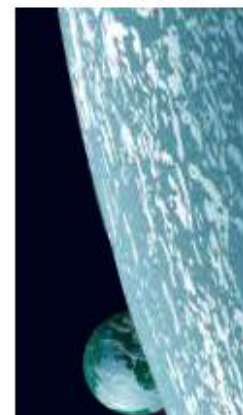
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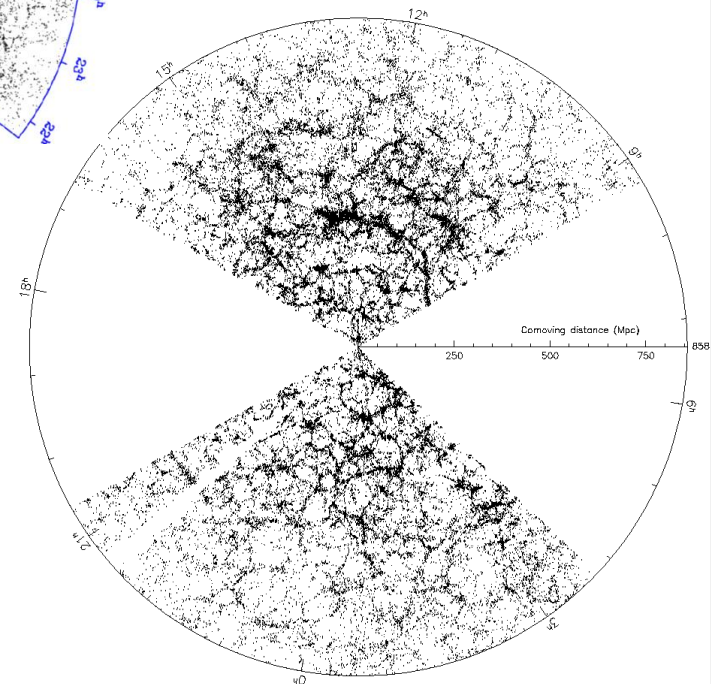
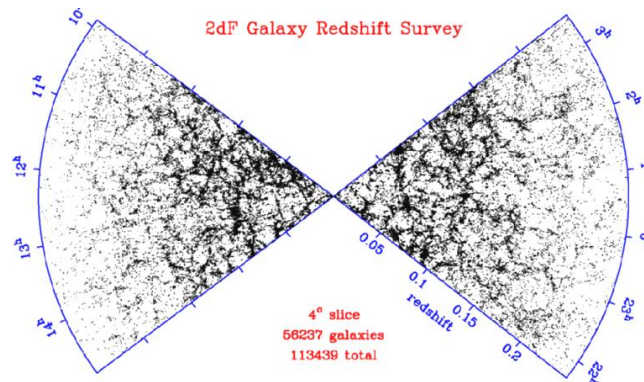


Modelling the Universe

Cosmological principle: universe is **homogeneous** and **isotropic** on large scales ($> 100 - 300$ Mpc).

Good evidence for this from:

- galaxy surveys
- isotropy of the CMBR



Modelling the Universe

Background cosmological model described by the **Robertson-Walker metric**

$$ds^2 = -dt^2 + R(t)^2 \left[\frac{dr^2}{1 - kr^2} + r^2 d\Omega^2 \right]$$

$R(t)$ = cosmic scale factor

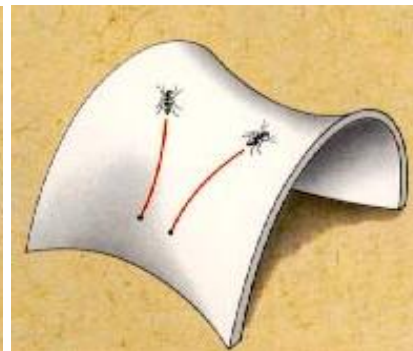
$$k = \text{curvature constant} = \begin{cases} -1, & \text{open} \\ 0, & \text{flat} \\ +1, & \text{closed} \end{cases}$$

$$\frac{R(t)}{R_0} = \frac{1}{1+z}$$

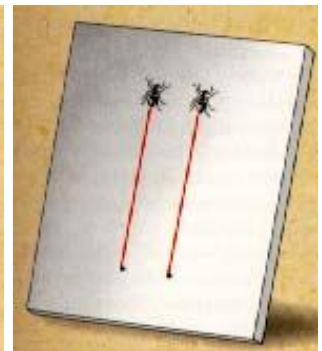
$$z = \frac{\lambda_{\text{obs}} - \lambda_{\text{emit}}}{\lambda_{\text{emit}}} = \text{redshift}$$



Closed



Open



Flat

Friedmann's Equations:

$$H^2 = \left(\frac{\dot{R}}{R} \right)^2 = \frac{8\pi G \rho}{3} + \frac{\Lambda}{3} - \frac{k}{R^2}$$

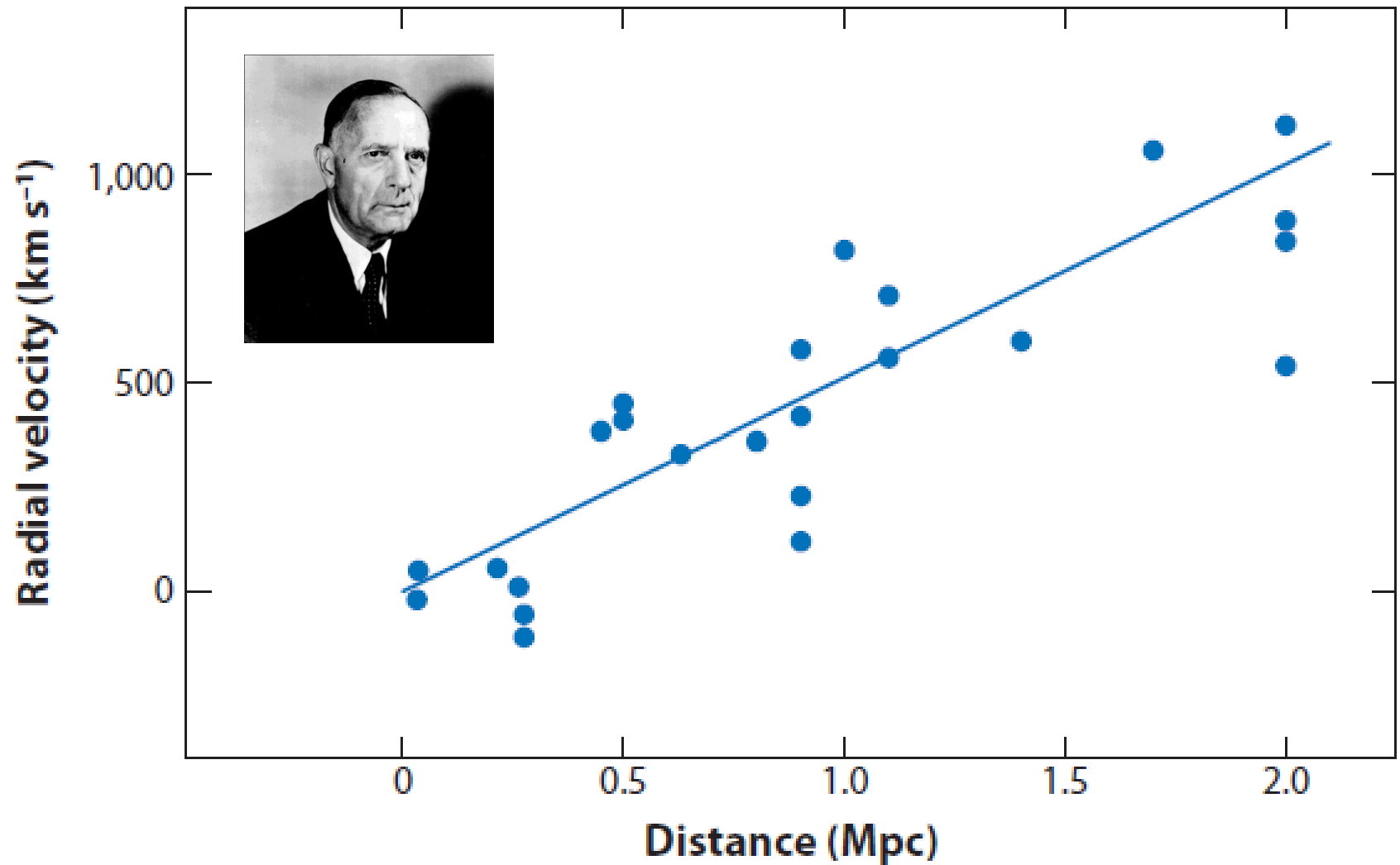
$$\frac{\ddot{R}}{R} = -\frac{4\pi G}{3}(\rho + 3P) + \frac{\Lambda}{3}$$

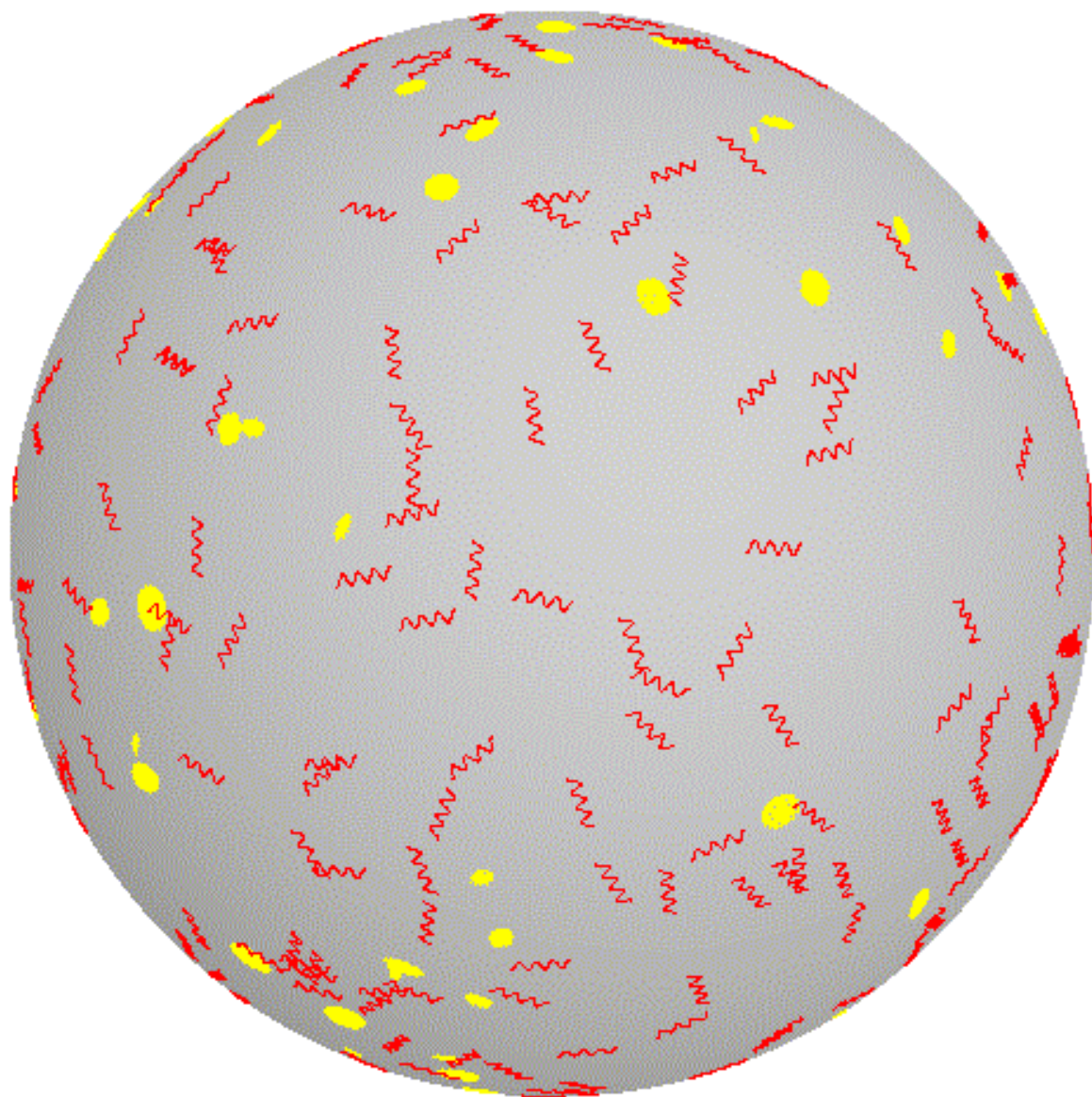
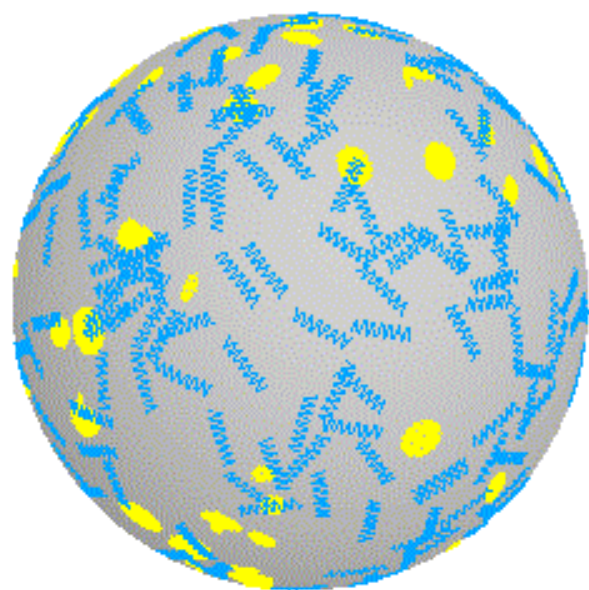
Can tune Λ to give $\dot{R} = 0$ for all t but *unstable*

(and *Hubble expansion* made idea redundant)

But Lambda term could still be non-zero anyway !

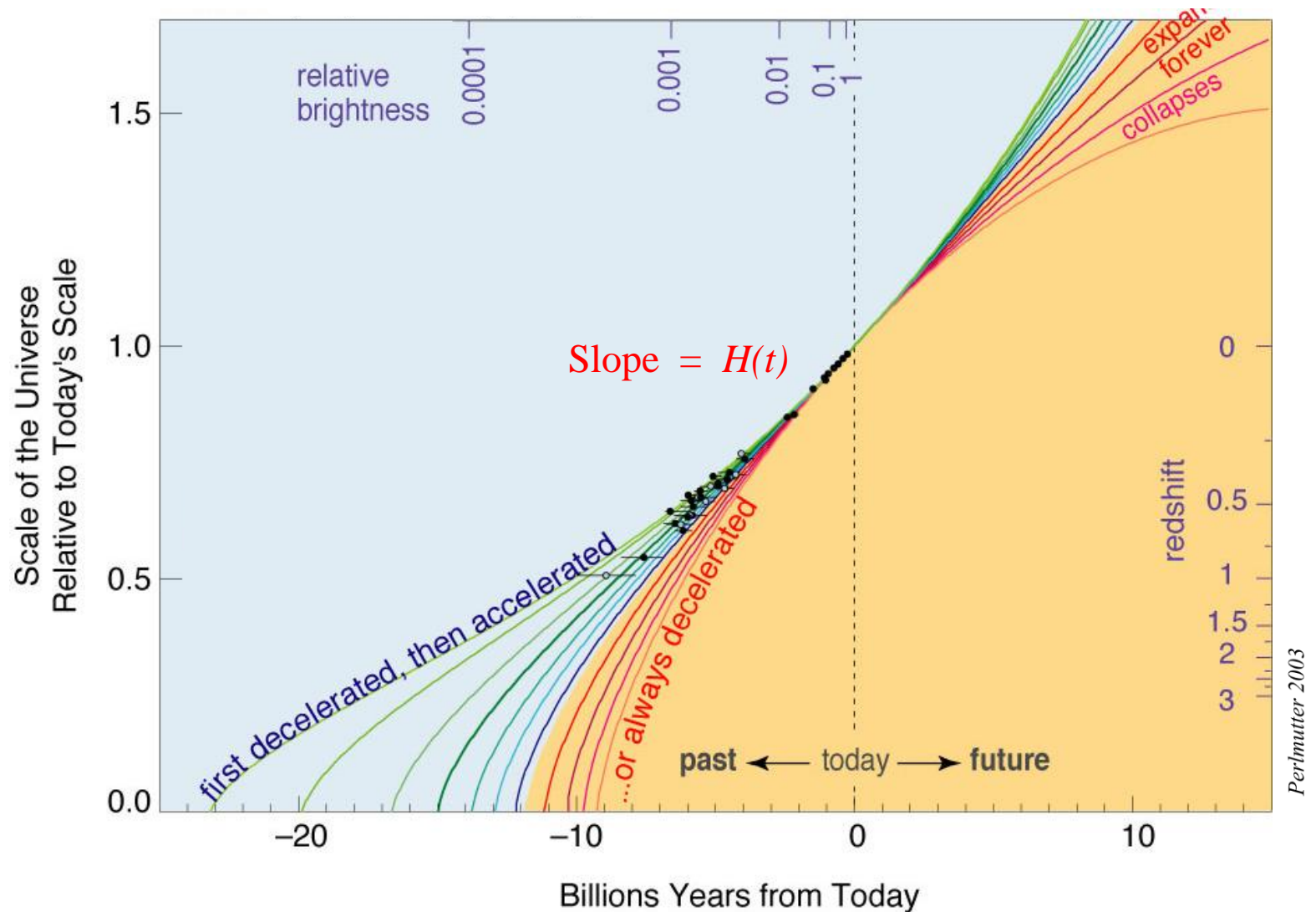
Adapted from Hubble (1929)





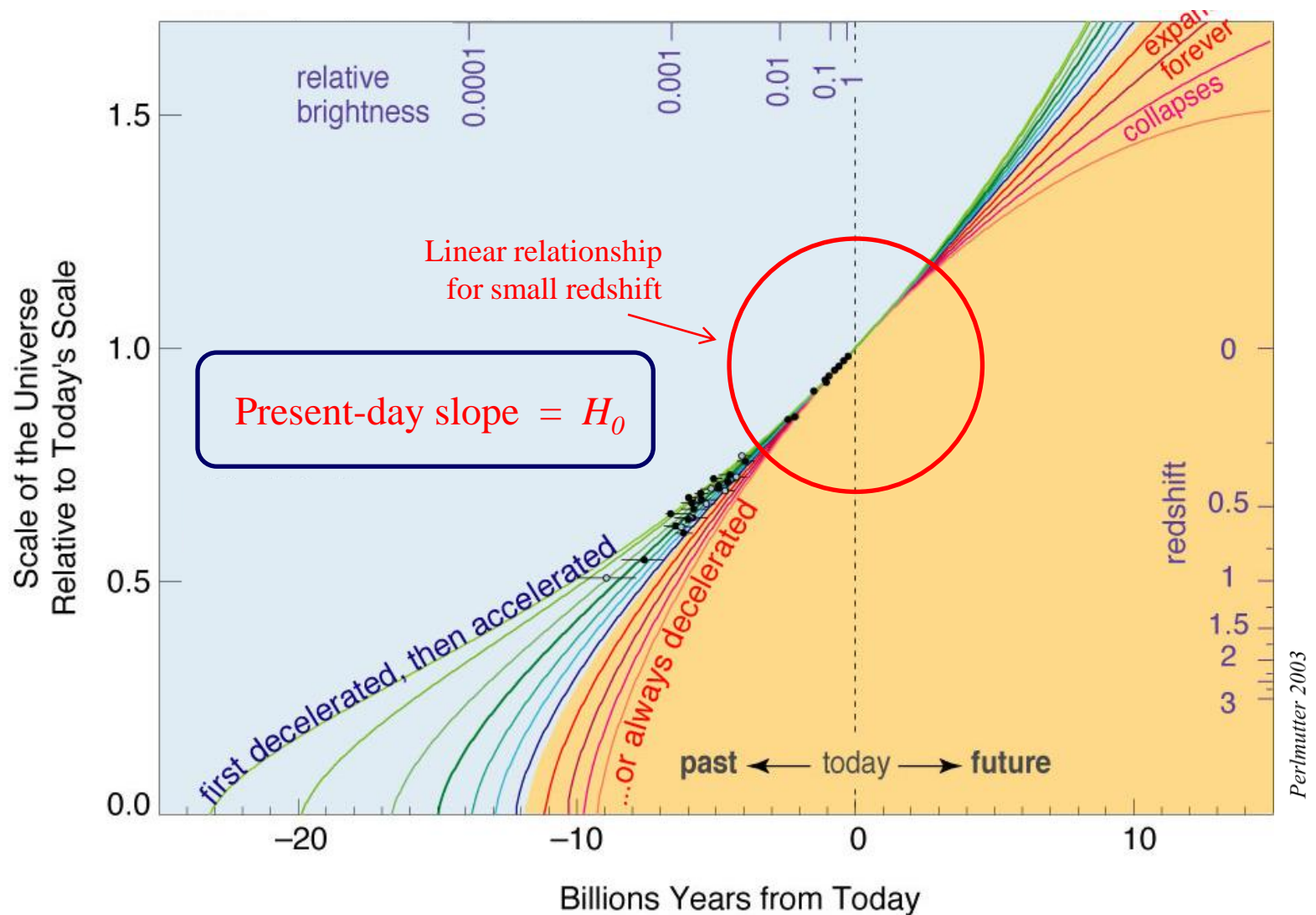
Evolution of scale factor given by Friedmann's equation

$$H^2 = \left(\frac{\dot{R}}{R} \right)^2 = \frac{8\pi G\rho}{3} + \frac{\Lambda}{3} - \frac{k}{R^2}$$



Evolution of scale factor given by Friedmann's equation

$$H^2 = \left(\frac{\dot{R}}{R} \right)^2 = \frac{8\pi G\rho}{3} + \frac{\Lambda}{3} - \frac{k}{R^2}$$



Re-expressing Friedmann's Equations

For $\Lambda = 0$

$$H^2 = \frac{8\pi G\rho}{3} - \frac{k}{R^2} \quad \Rightarrow \quad k = 0 \Leftrightarrow \rho = \left[\frac{8\pi G}{3H^2} \right]^{-1} = \rho_{\text{crit}}$$

Define

$$\Omega_m = \rho / \rho_{\text{crit}} = \frac{8\pi G\rho}{3H^2} \quad \Omega_\Lambda = \frac{\Lambda}{3H^2} \quad \Omega_k = -\frac{k}{R^2 H^2}$$

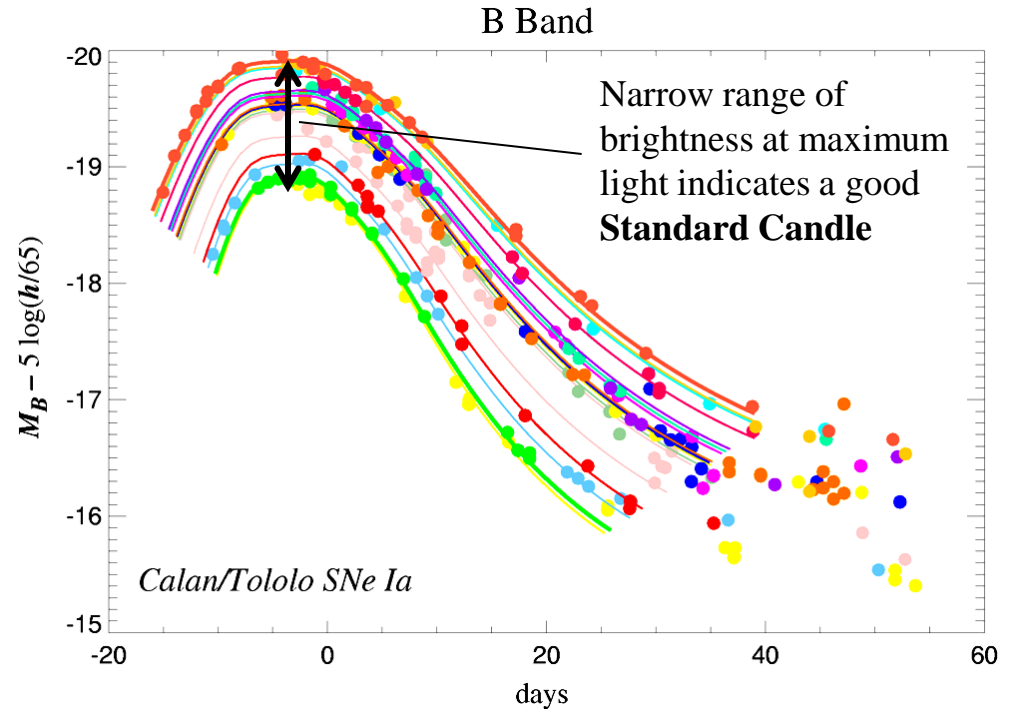
It follows that, at *any time*

$$\Omega_m + \Omega_\Lambda + \Omega_k = 1$$

Is the Universe speeding up or slowing down?



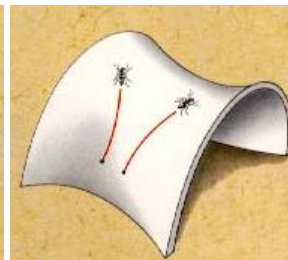
We can answer this question using **type Ia supernovae**



Geometry of the universe affects the relationship between redshift and luminosity distance of distant supernovae



Closed

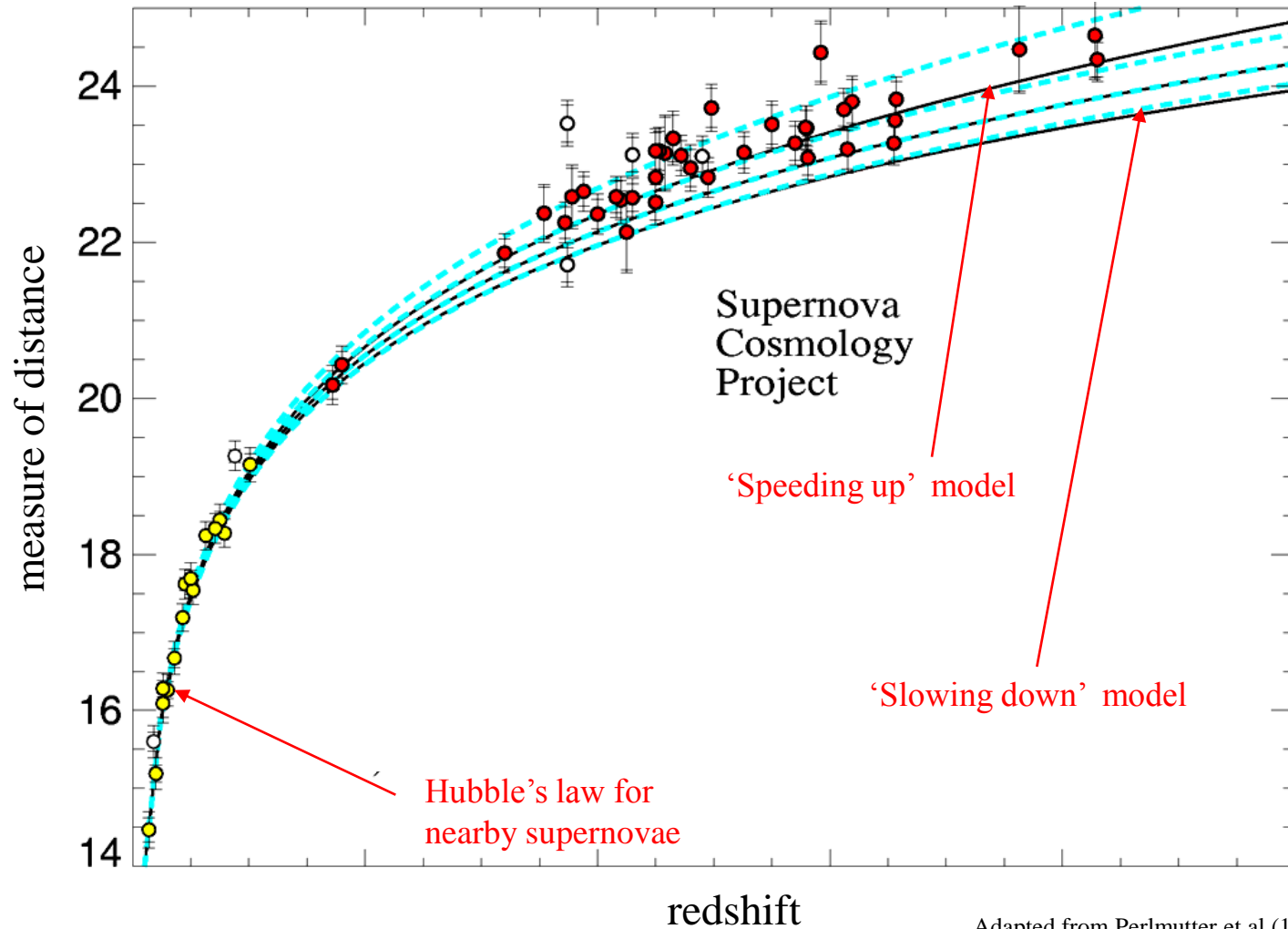


Open

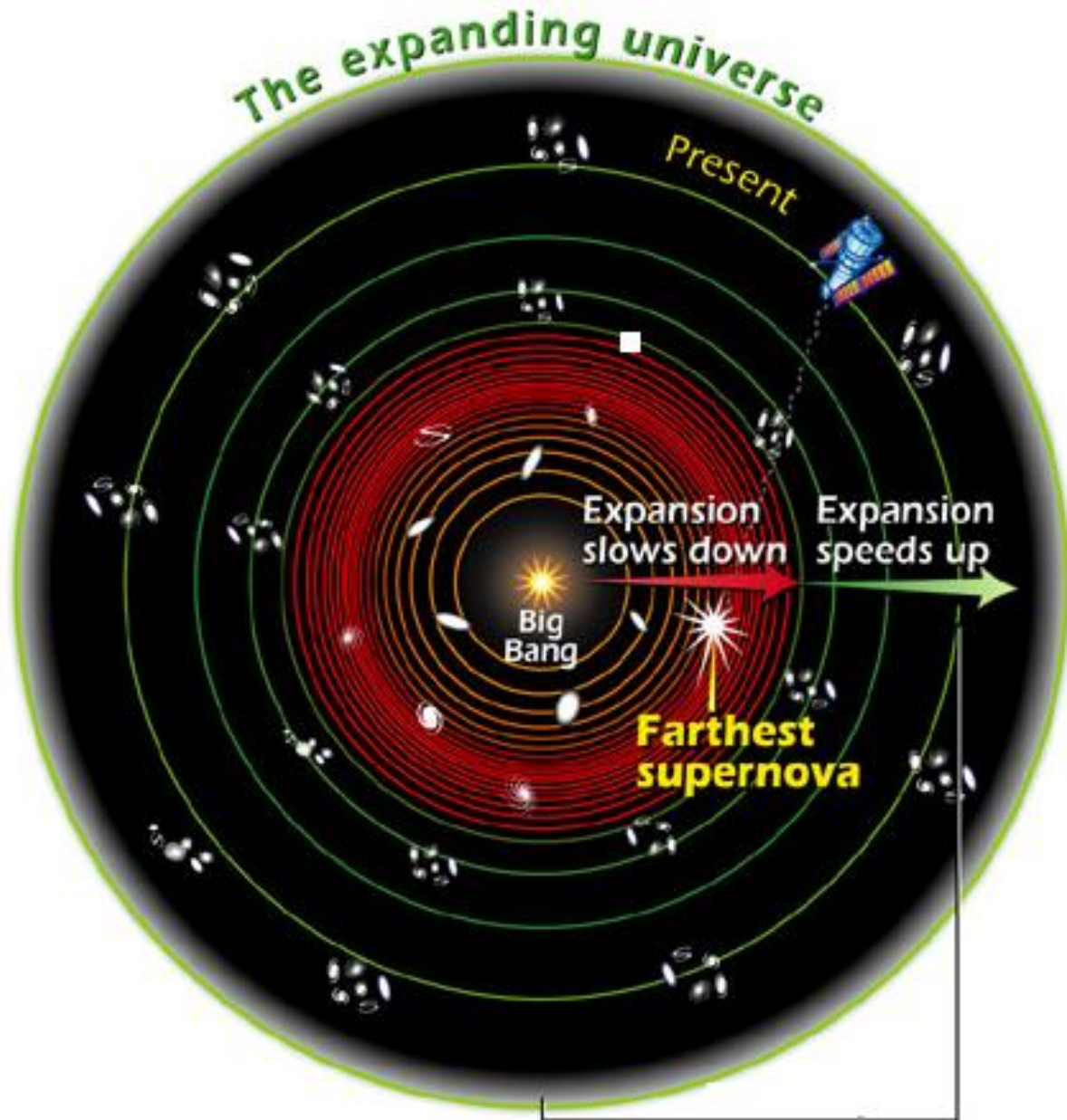


Flat

Hubble diagram of distant supernovae



Adapted from Perlmutter et al (1998)



The Nobel Prize in Physics 2011
Saul Perlmutter, Brian P. Schmidt, Adam G. Riess

The Nobel Prize in Physics 2011	▼
Nobel Prize Award Ceremony	▼
Saul Perlmutter	▼
Brian P. Schmidt	▼
Adam G. Riess	▼



Photo: U. Montan

Saul Perlmutter



Photo: U. Montan

Brian P. Schmidt



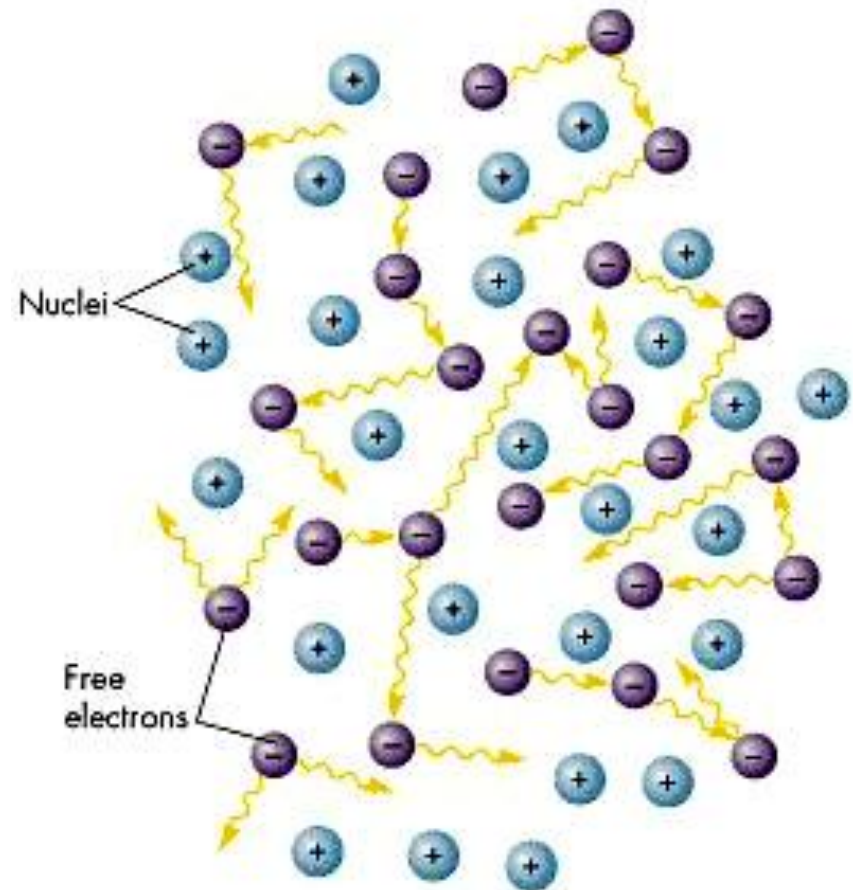
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Adam G. Riess

Dark Dark Energy

Early Universe too hot for
neutral atoms to exist

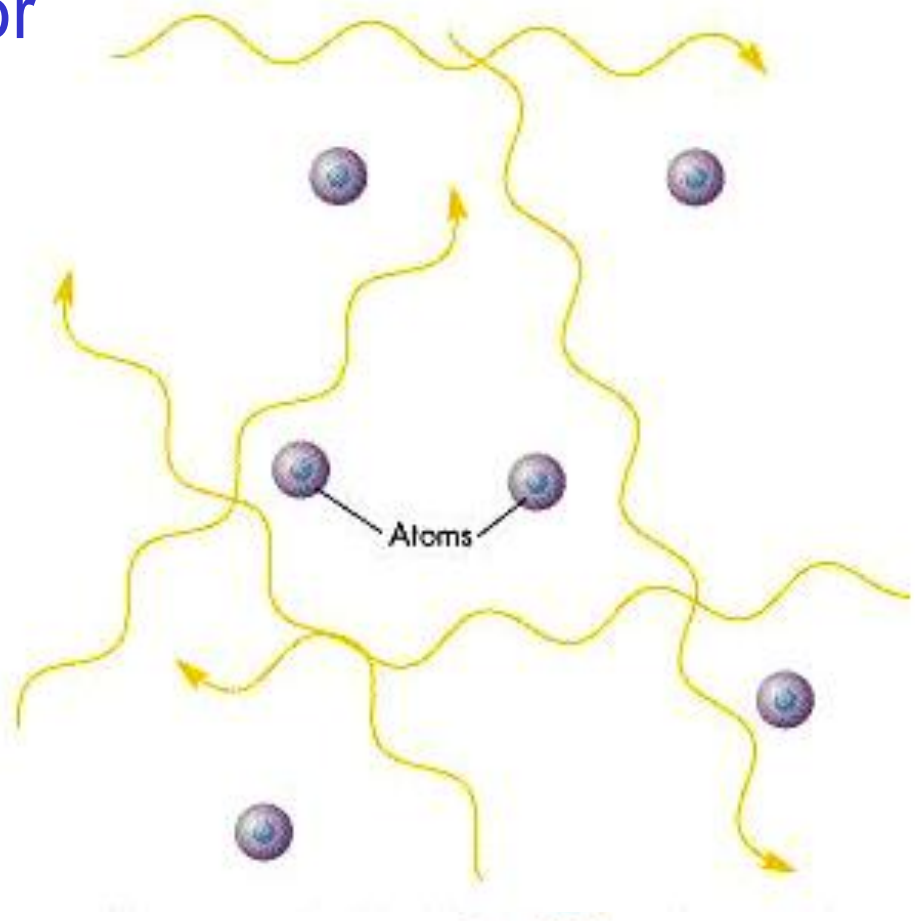
Free electrons scatter
light (as in a fog)



Early Universe too hot for
neutral atoms to exist

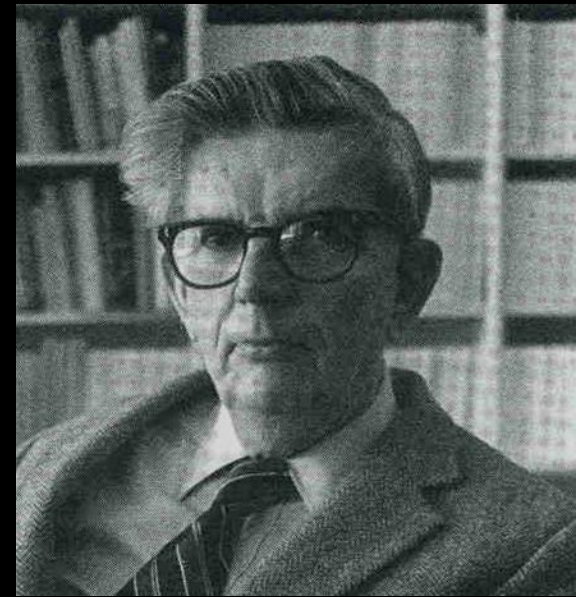
Free electrons scatter
light (as in a fog)

After ~380,000 years,
Universe cool enough
for neutral hydrogen to
form: the fog clears!



**Background radiation
predicted in 1950s and 1960s
by Gamov, Dicke, Peebles.**

**Discovered in 1965 by
Penzias and Wilson**



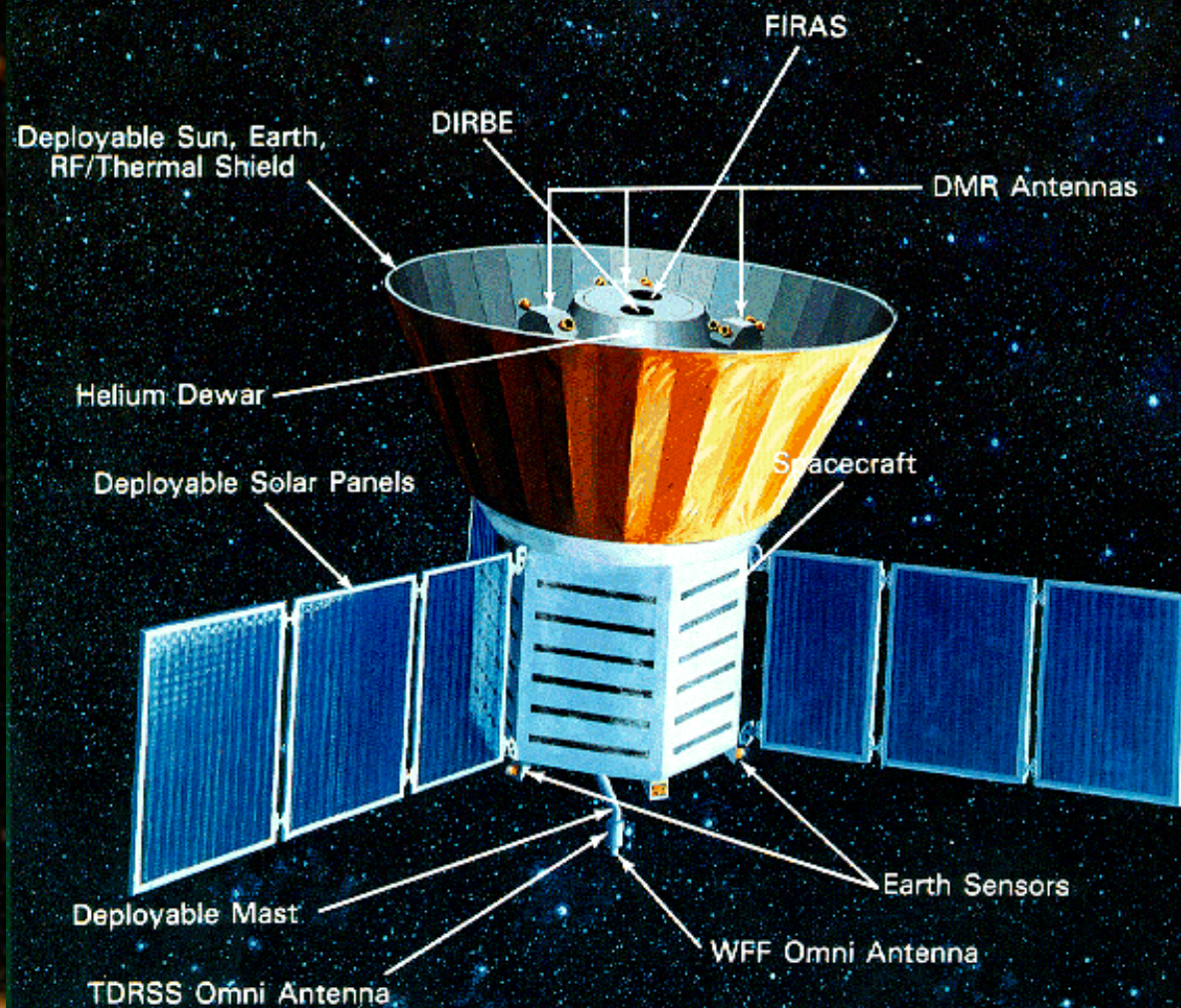
Robert Dicke



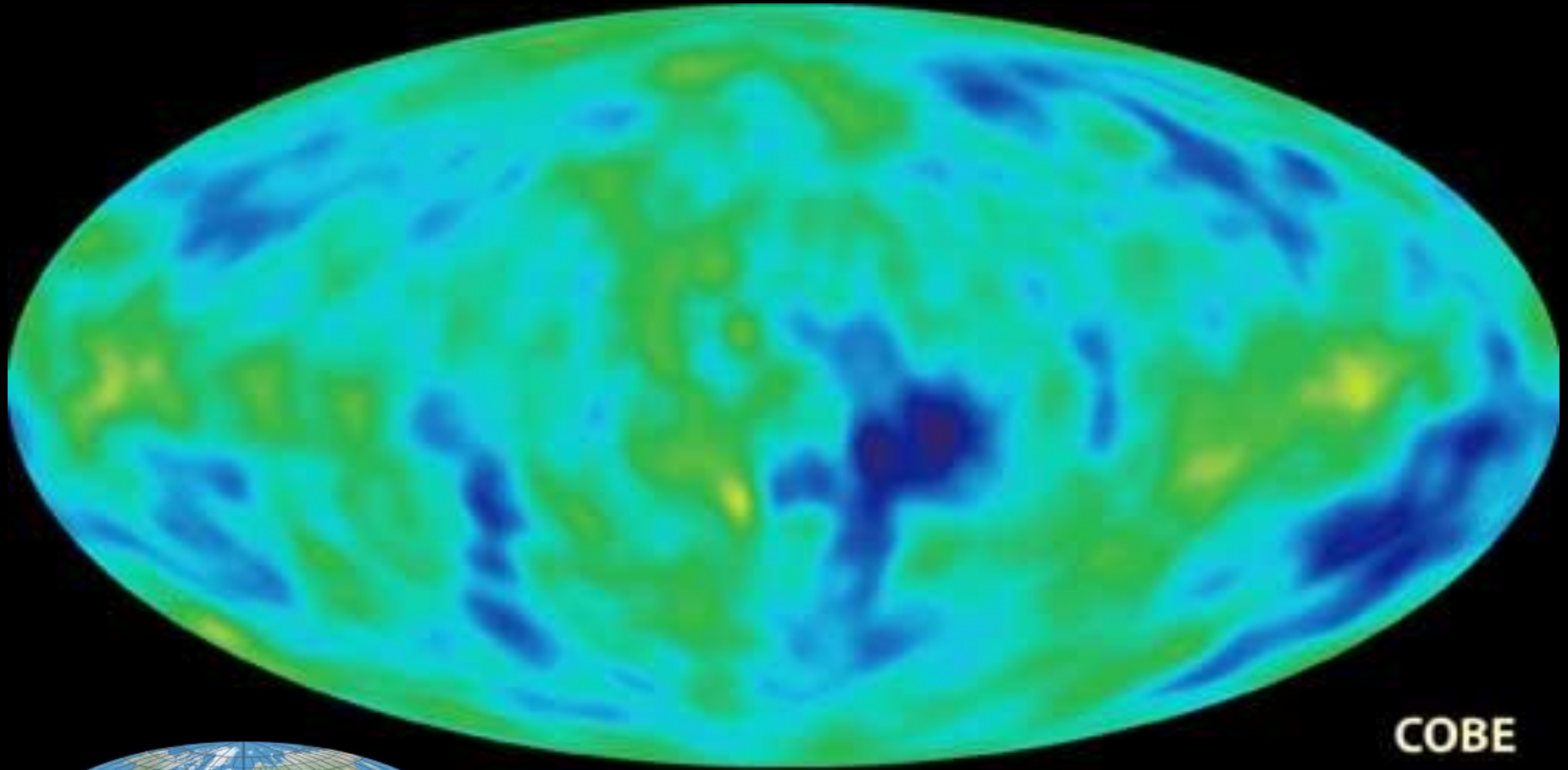
Arno Penzias and Robert Wilson



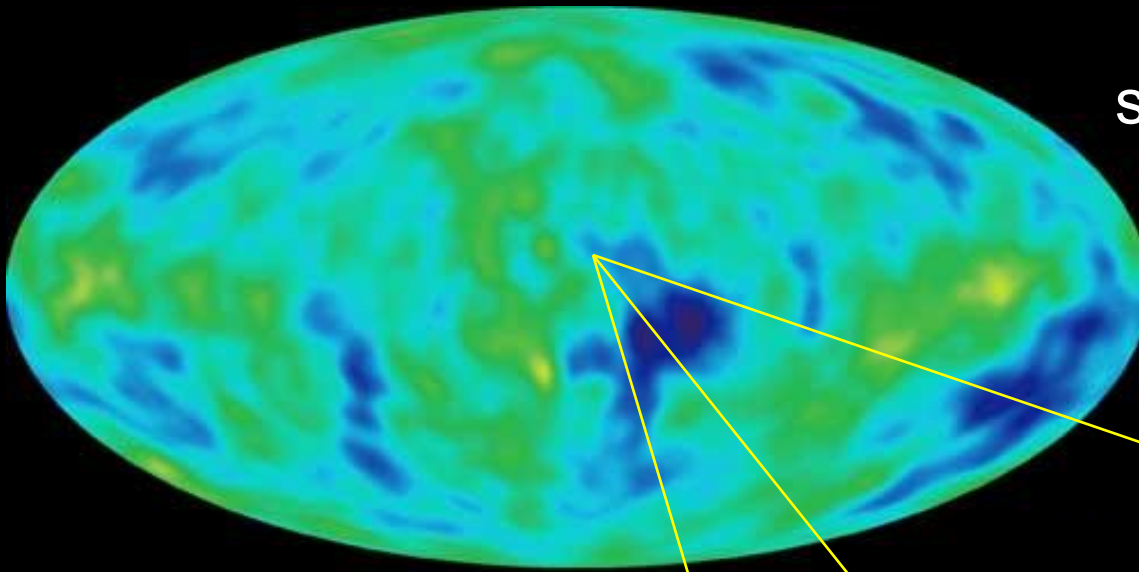
Jim Peebles



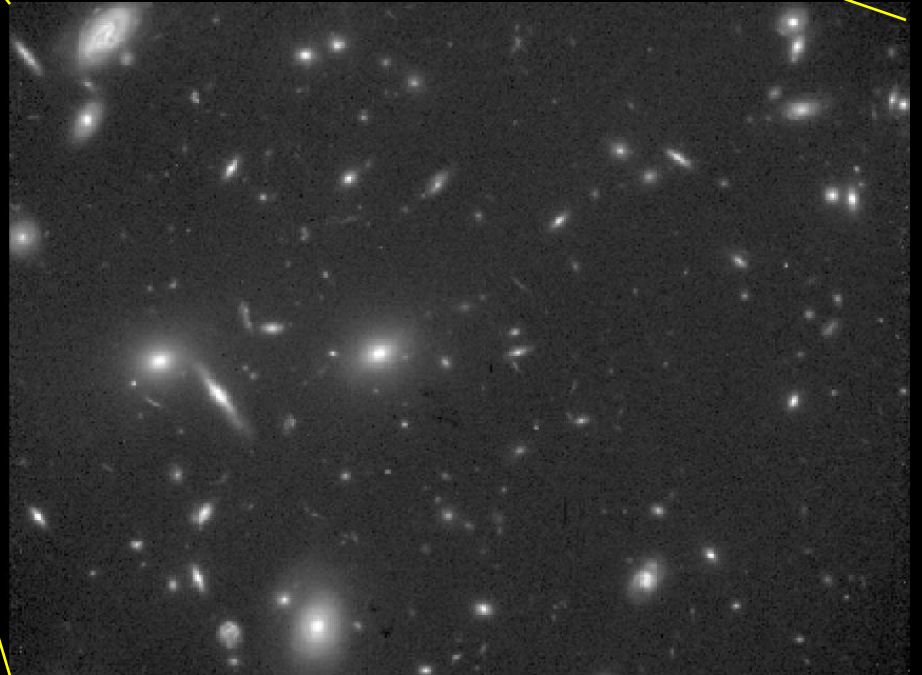
COBE map of temperature across the sky



CMBR 'ripples' are the seeds of today's galaxies

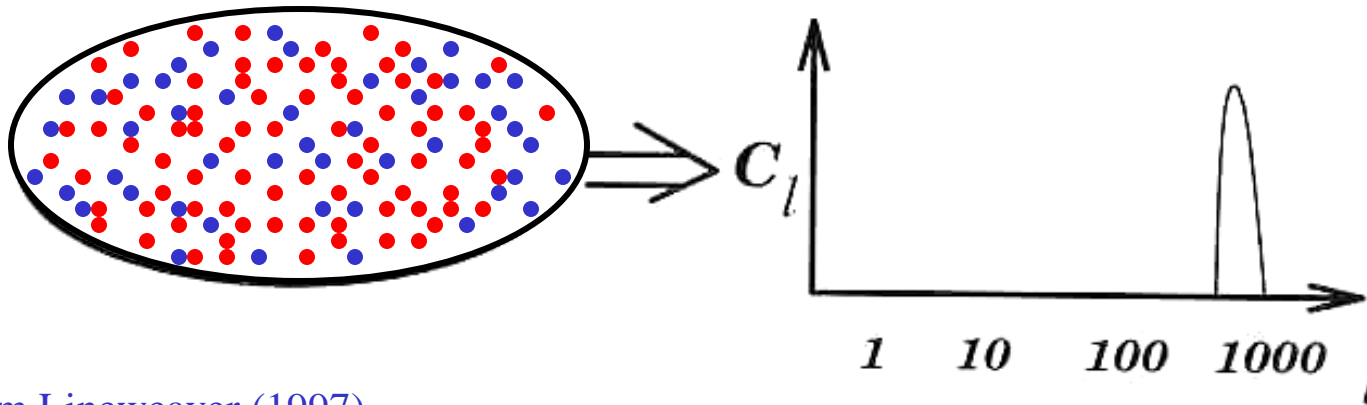
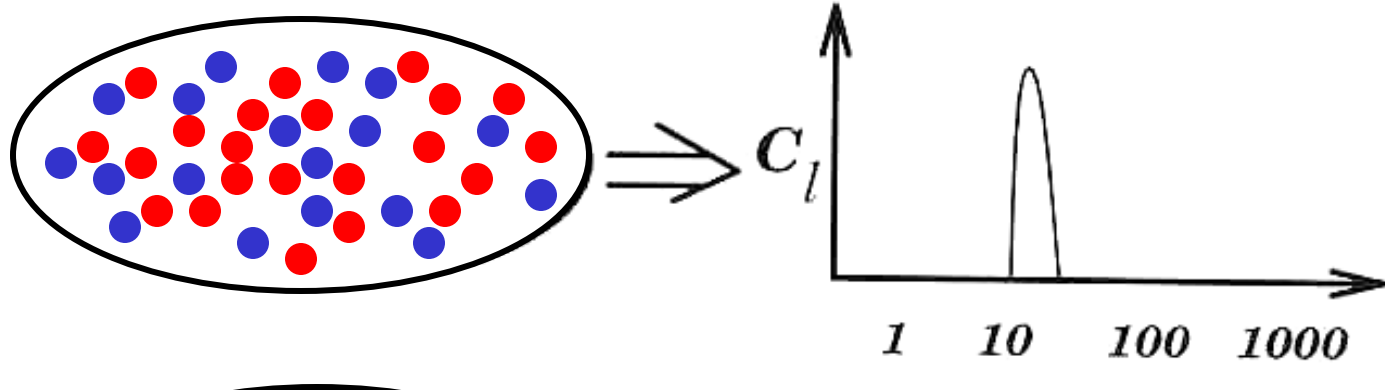
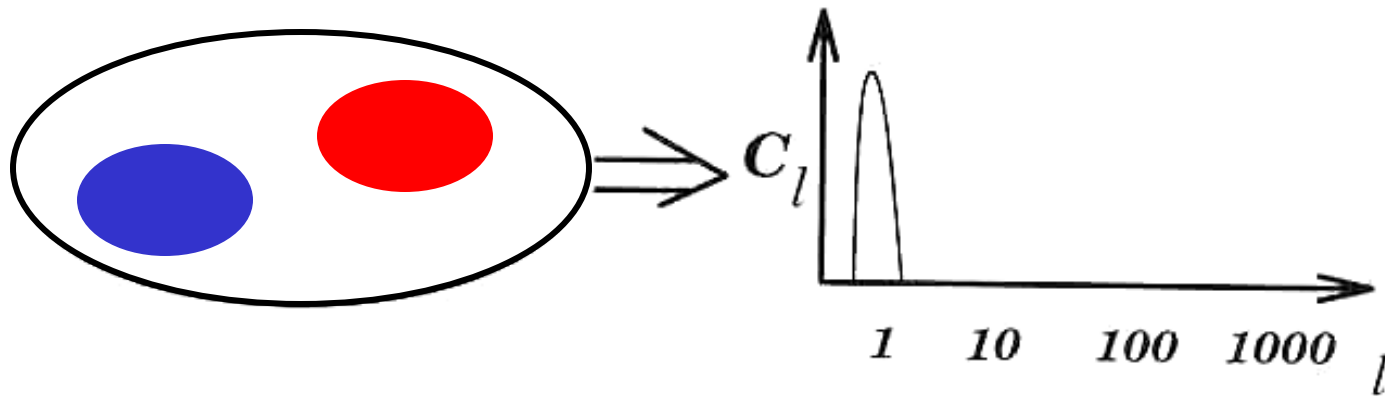


Galaxy formation is highly sensitive to the pattern of CMBR temperature



Maps

Power Spectra





COBE

W-band temperature anisotropy

WMAP

Internal Linear Combination of 5 bands, smoothed

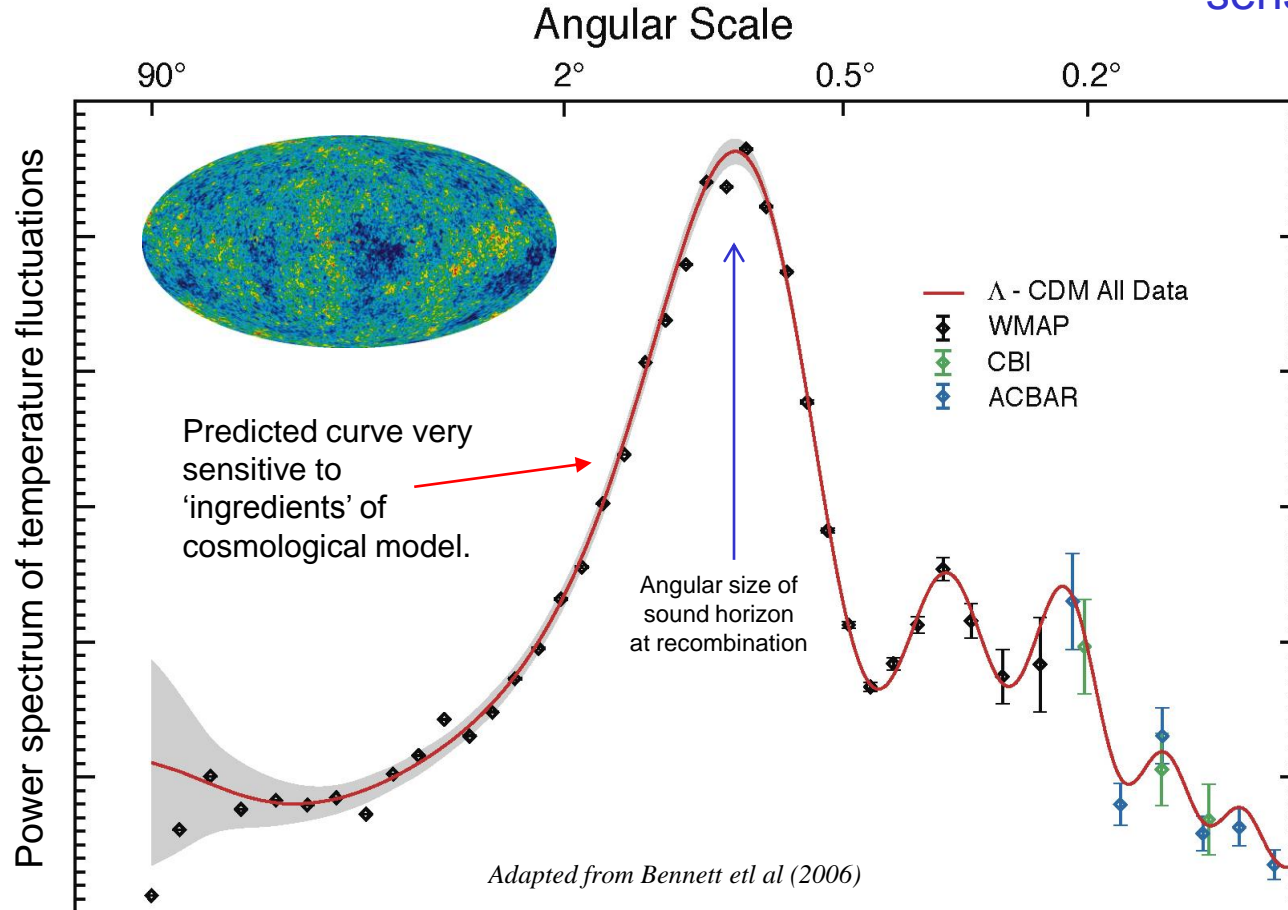
Simulated temperature anisotropy

PLANCK

Simulated temperature and polarization anisotropy

Cosmological parameters from the CMBR

Amplitude of Doppler peaks
sensitive to density of matter



Friedmann equation

$$H^2 = \left(\frac{\dot{R}}{R} \right)^2 = \frac{8\pi G\rho}{3} + \frac{\Lambda}{3} - \frac{k}{R^2}$$

For $\Lambda = 0$

$$k = 0 \Leftrightarrow \rho = \left[\frac{8\pi G}{3H^2} \right]^{-1} = \rho_{\text{crit}}$$

$$\Omega_m = \frac{\rho}{\rho_{\text{crit}}} = \frac{8\pi G\rho}{3H^2}$$

i.e. $\rho \propto H^2$

CMBR spectrum tightly constrained but H_0 degenerate with other parameters

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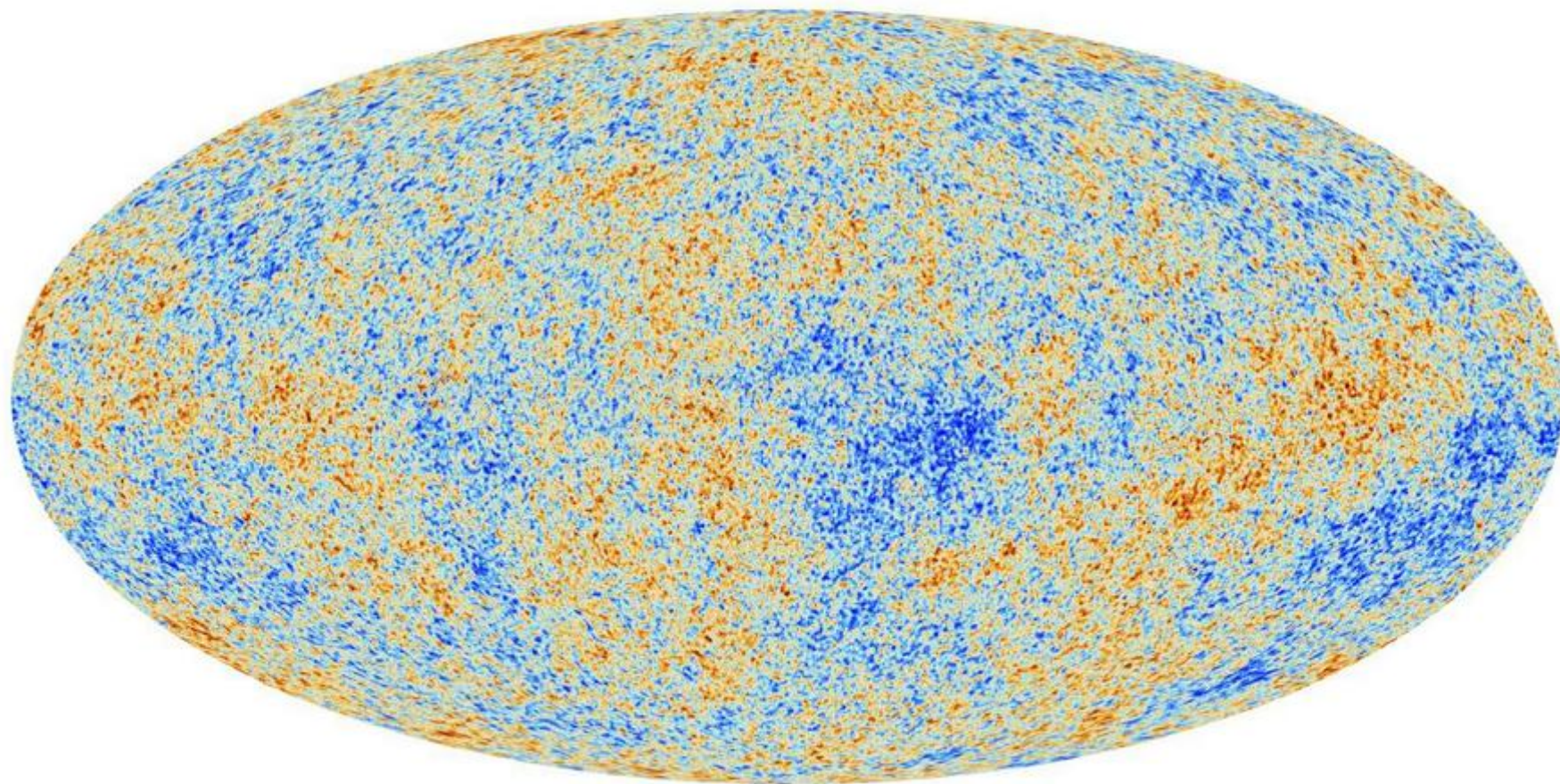
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21 March 2013 Last updated at 10:00

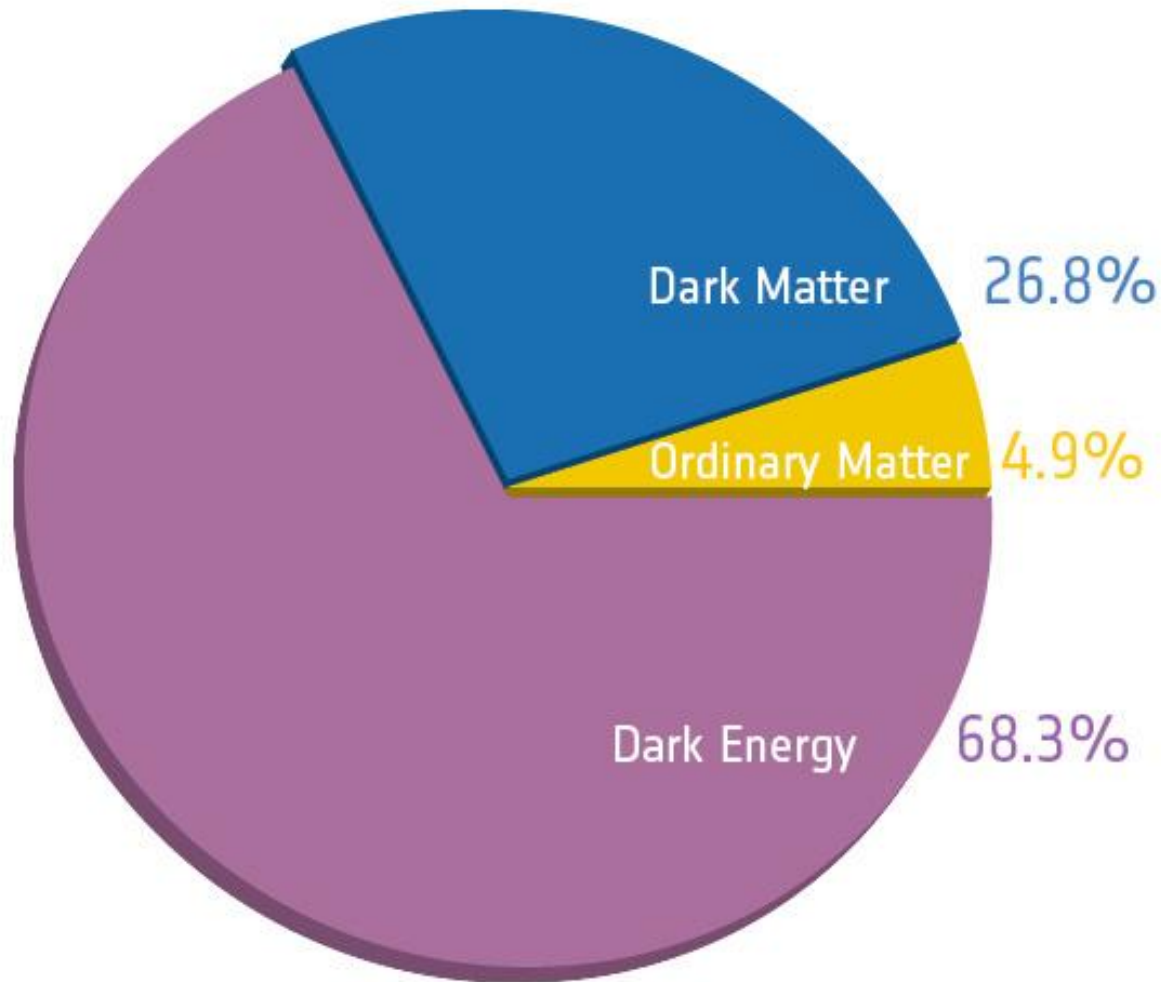
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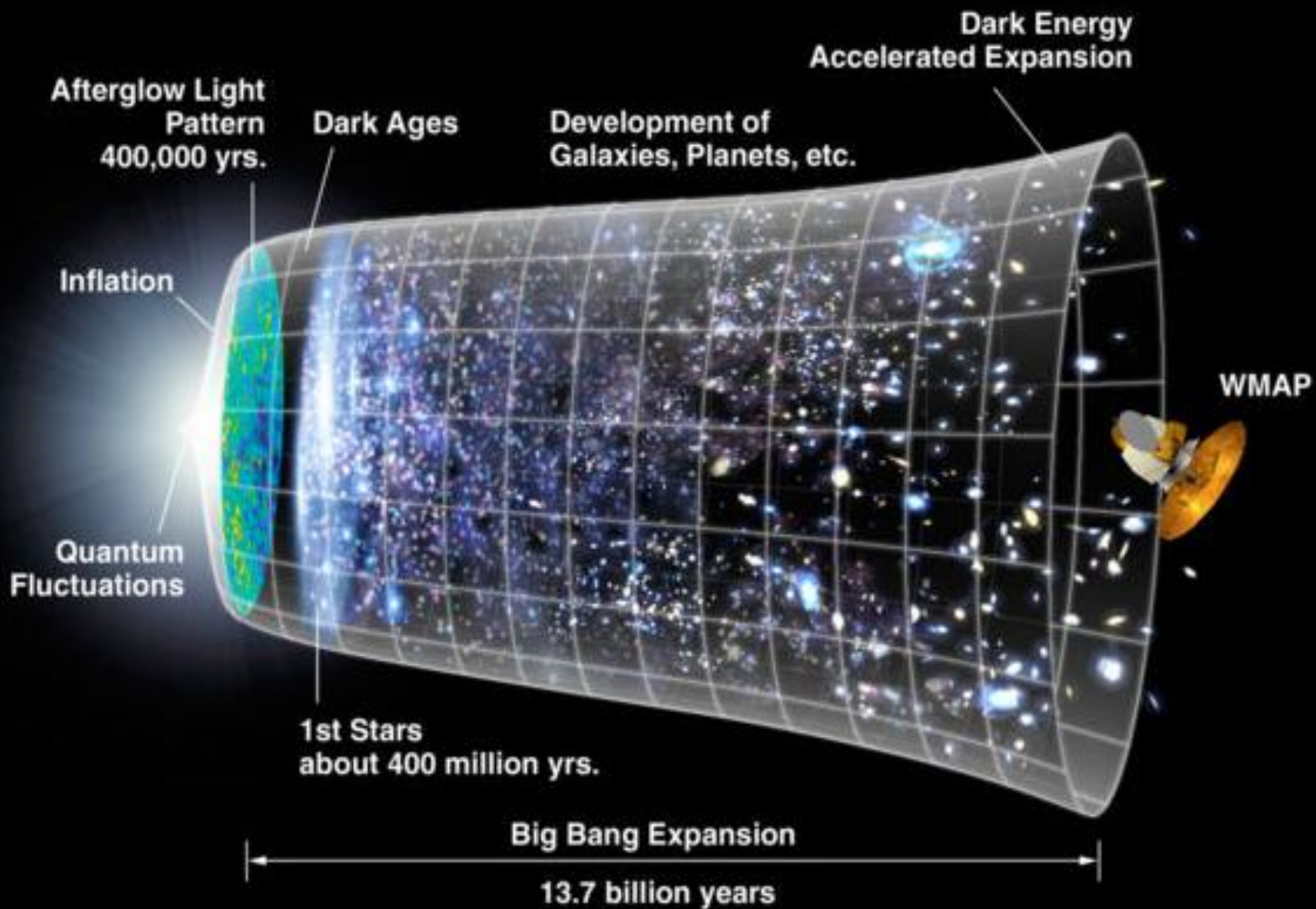
Planck satellite: Maps detail Universe's ancient light

A map tracing the "oldest light" in the sky has been produced by Europe's Planck Surveyor satellite. Its pattern confirms the Big Bang theory for the origin of the Universe but subtle, unexpected details will require scientists to adjust some of their ideas.



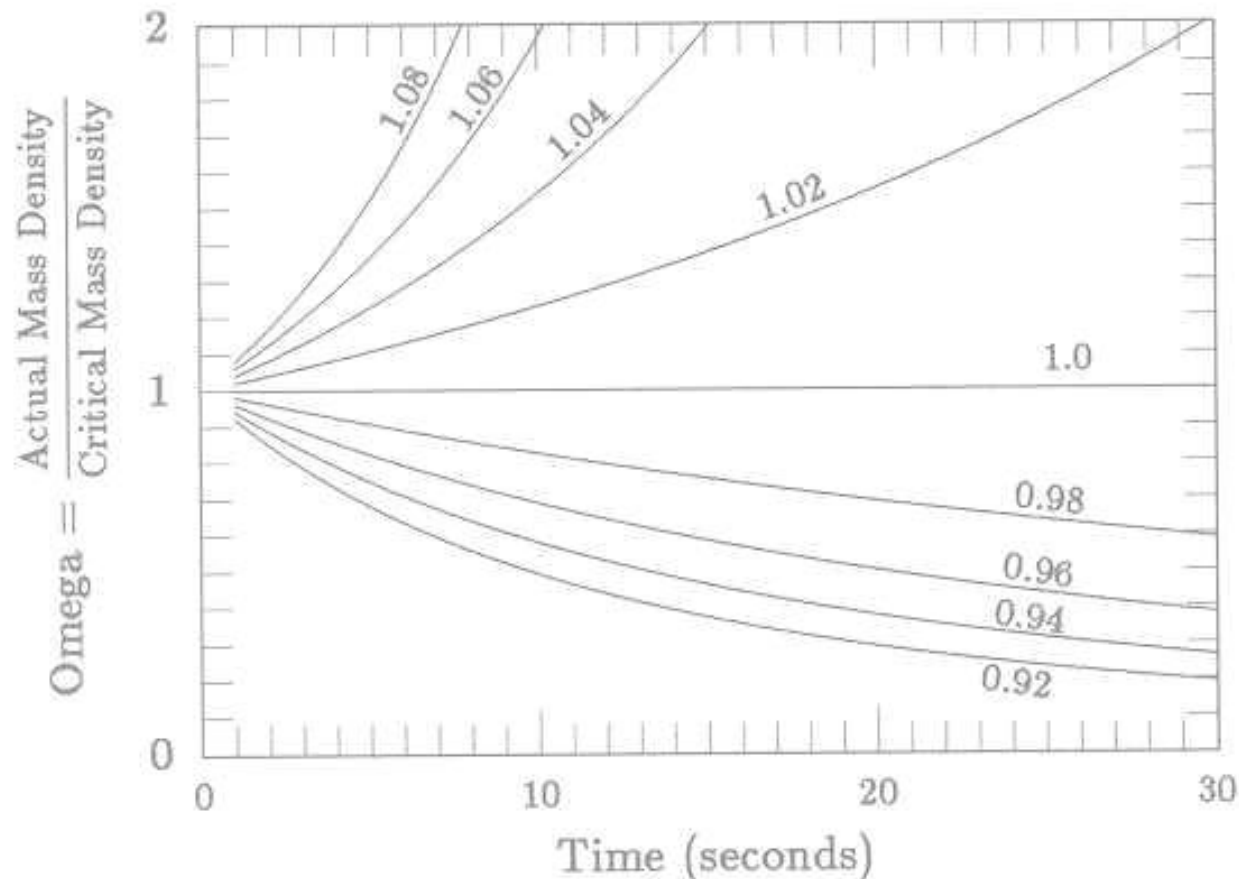
State of the Universe: Mar 2013





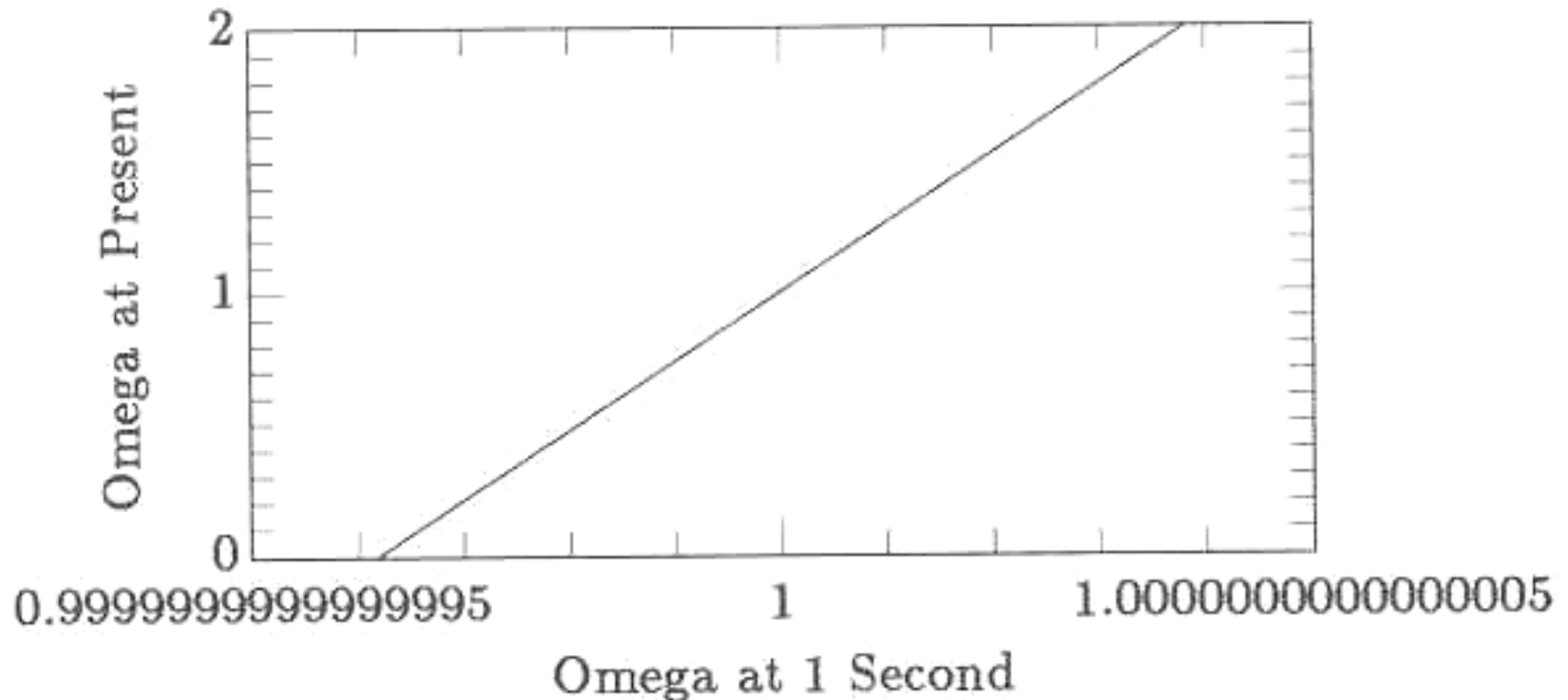
Inflation for astronomers

We have been considering $\Omega_{k0} = 0$ but suppose that in the past $\Omega_k \neq 0$. From the Friedmann equations it would then be very difficult to explain why it is so close to zero today.

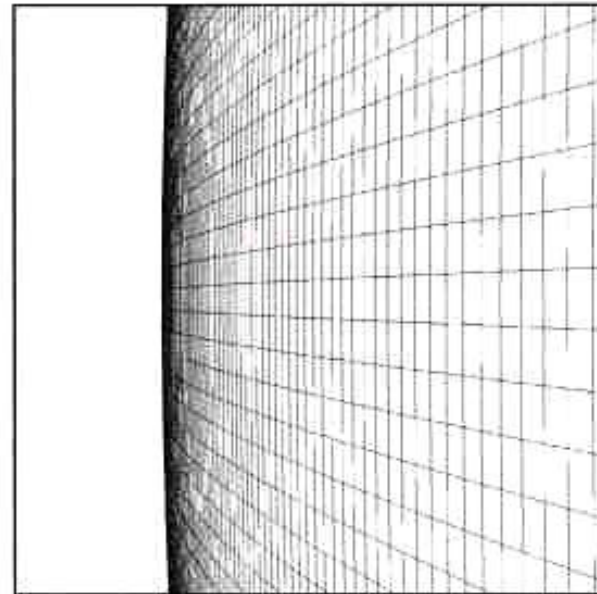
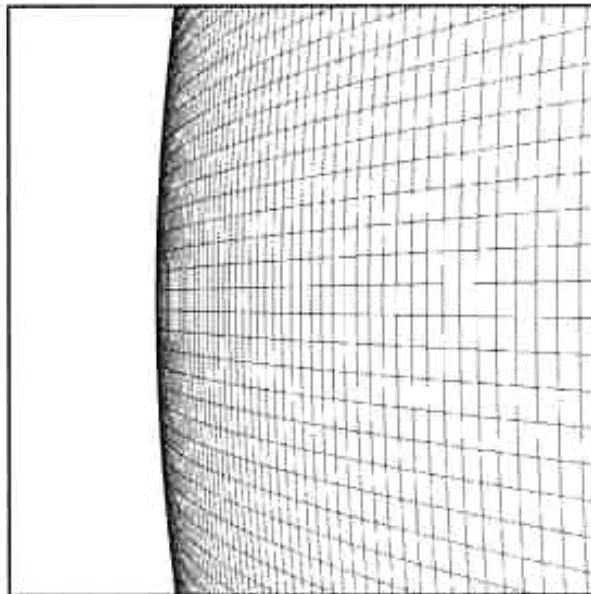
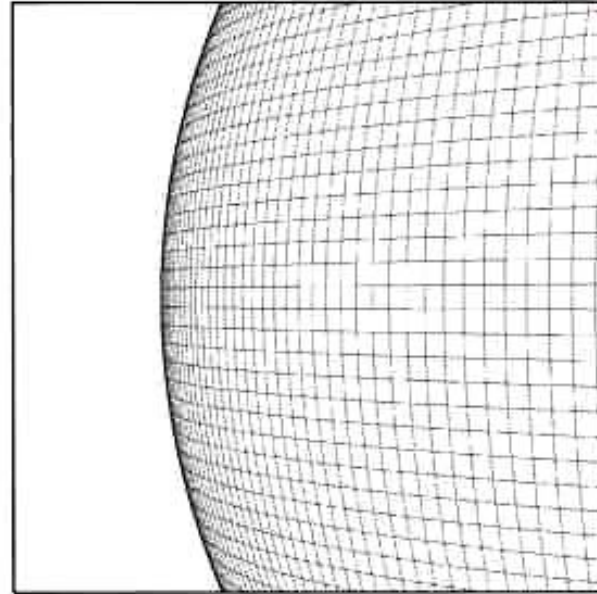
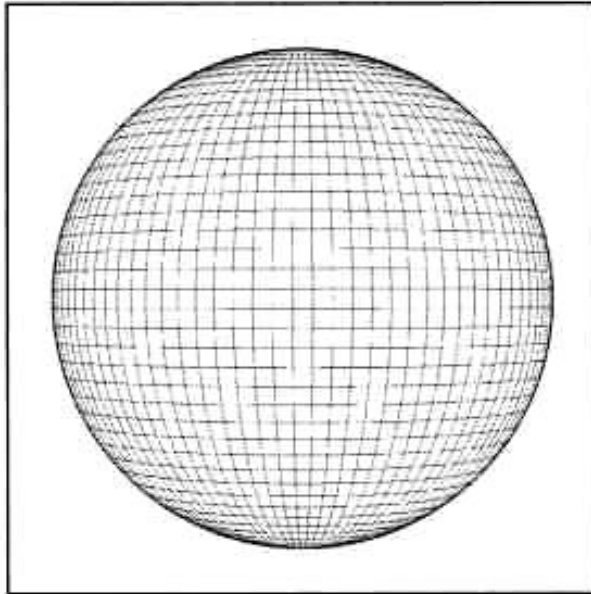


Present day ‘closeness’ of matter density to the critical density appears to require an incredible degree of ‘fine tuning’ in the very early Universe.

FLATNESS PROBLEM



Inflationary solution to the Flatness Problem

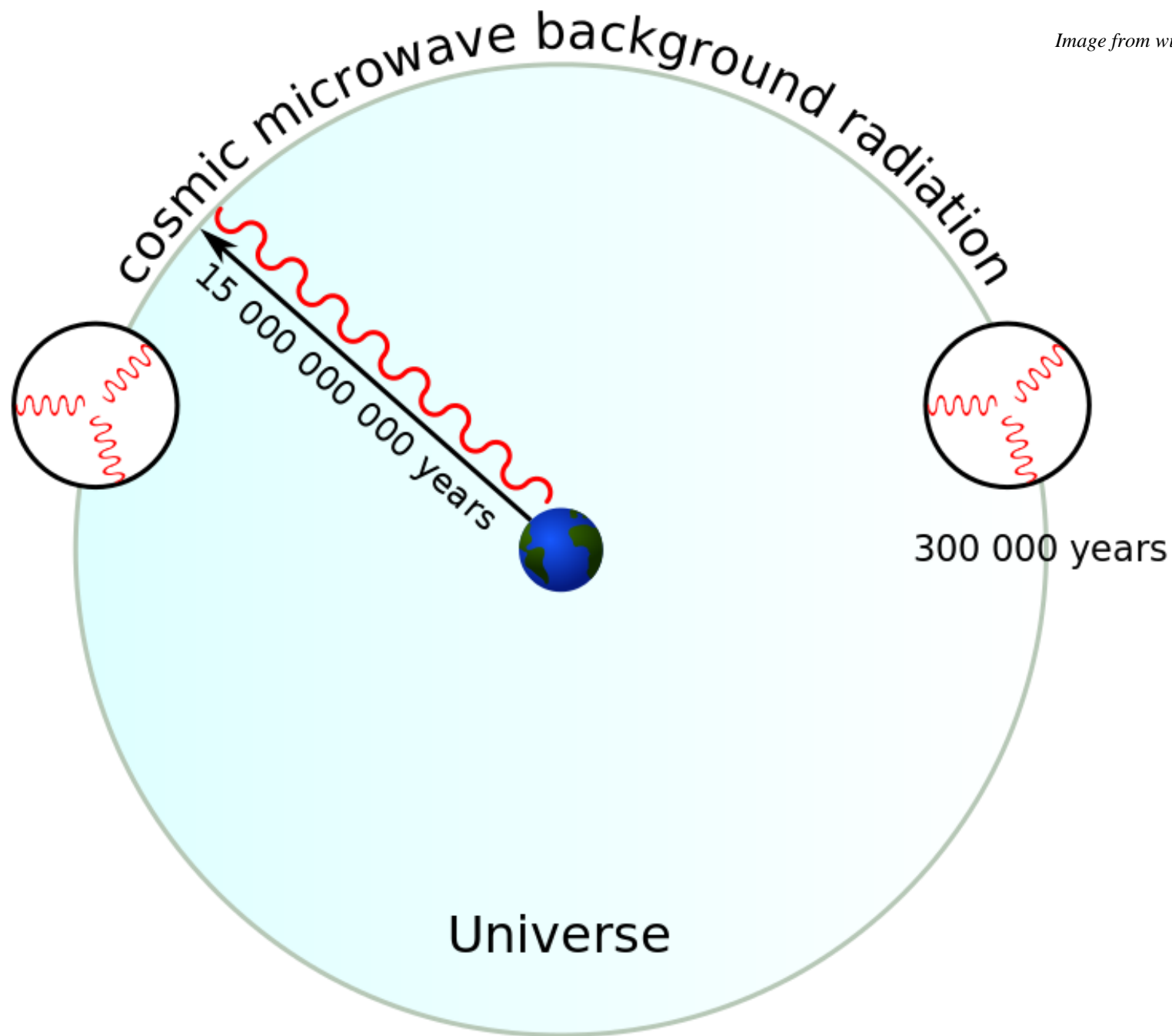


How do we explain the isotropy of the CMBR, when opposite sides of the sky were ‘causally disconnected’ when the CMBR photons were emitted?

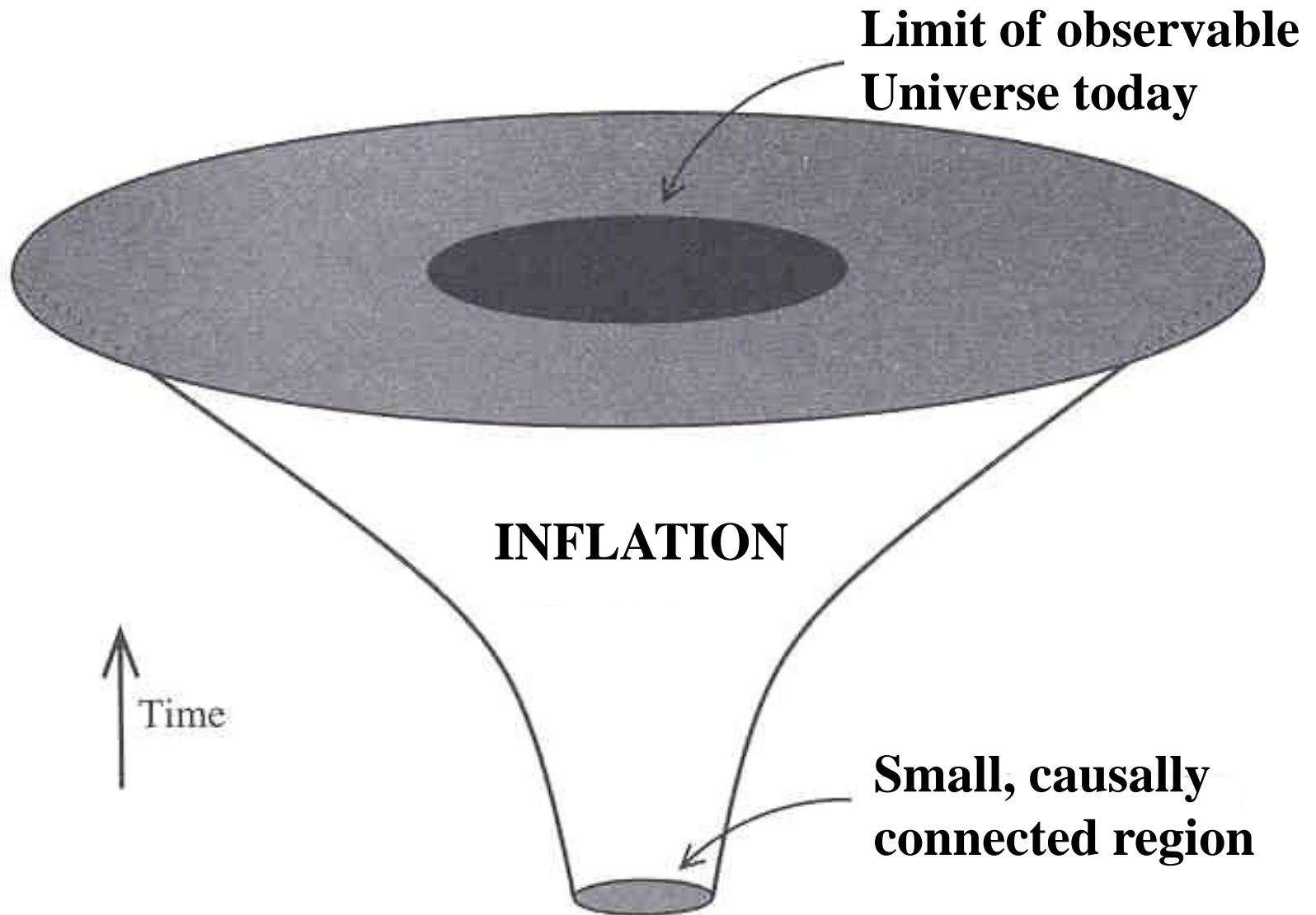
HORIZON PROBLEM

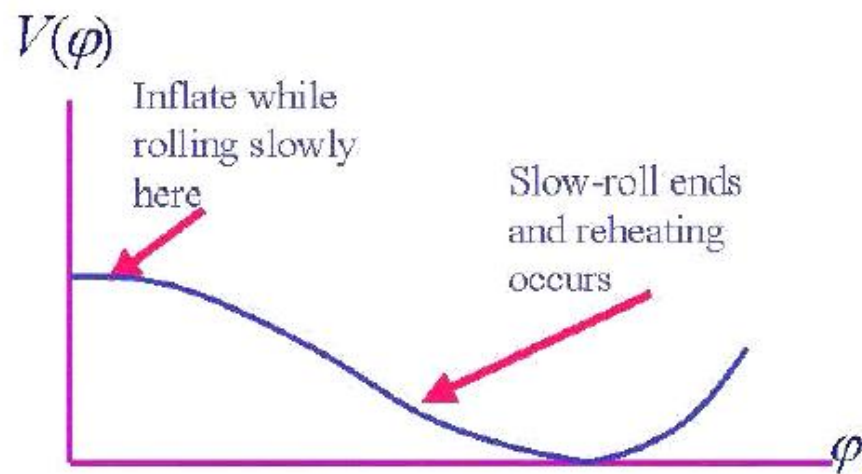
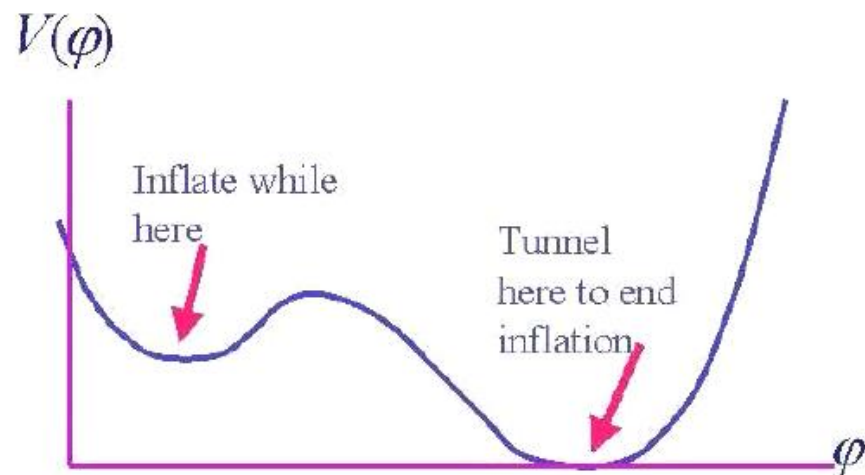


$T = 2.728 \text{ K}$



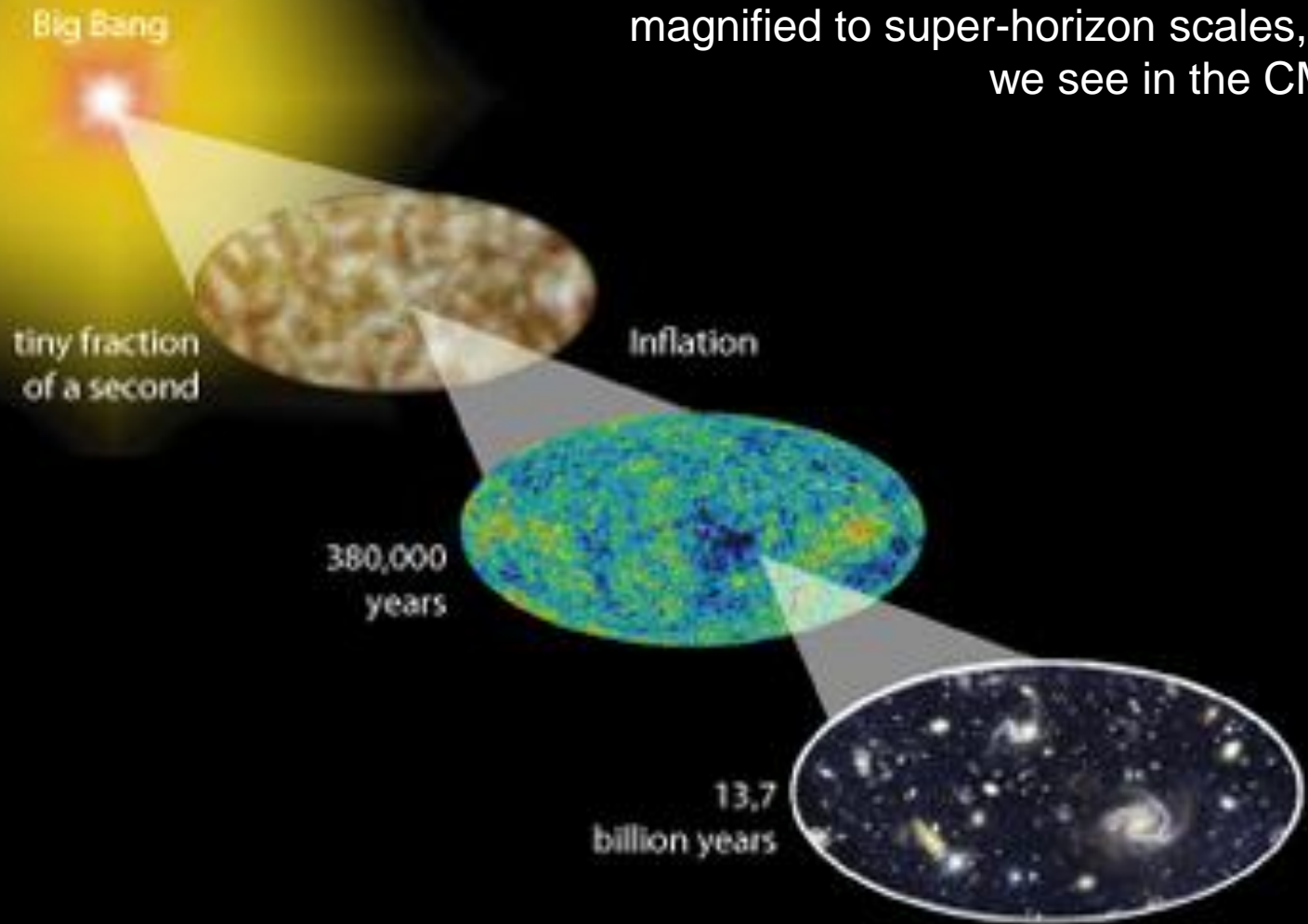
Inflationary solution to the Horizon Problem



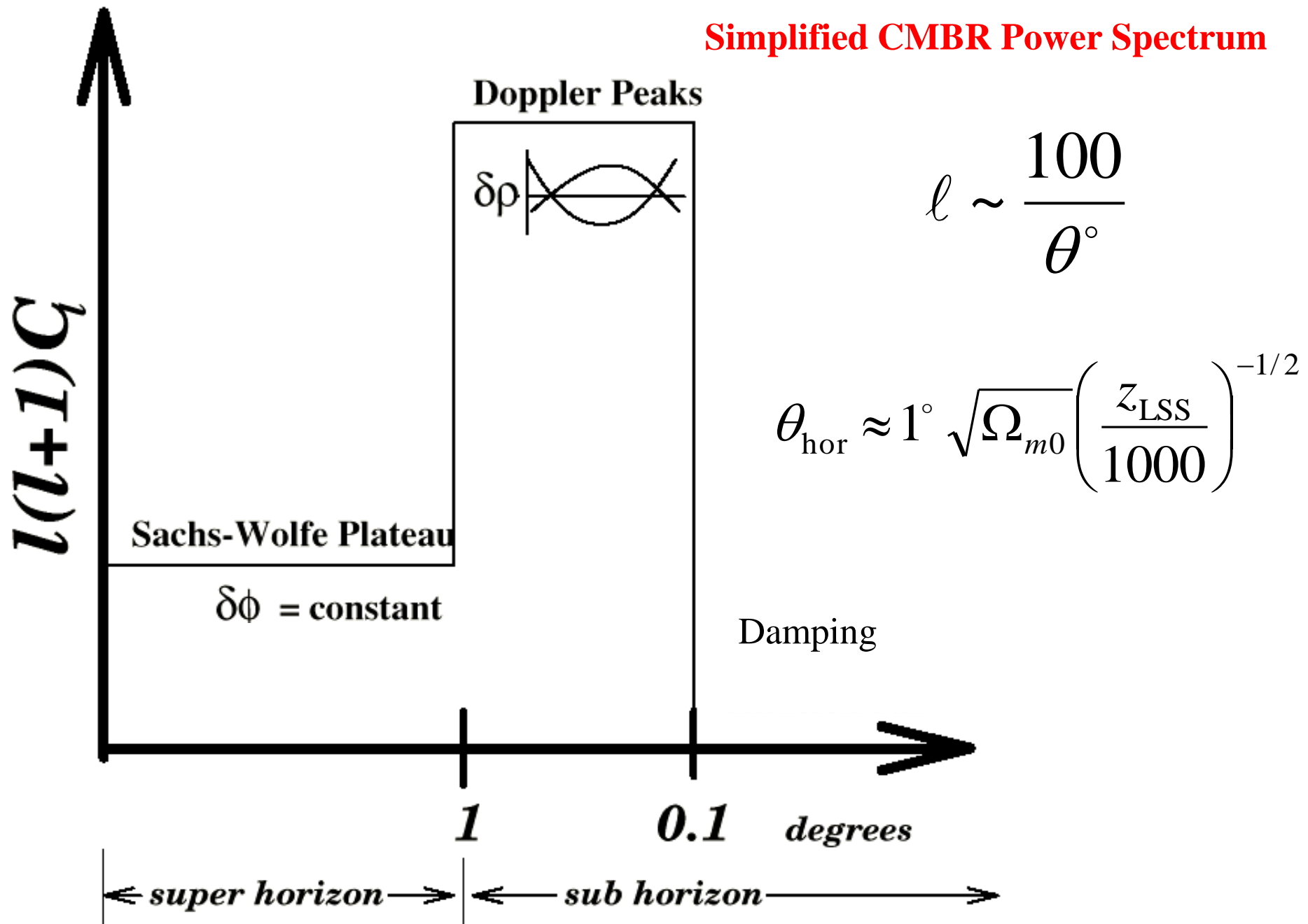


Inflation also provides a mechanism for generating CMBR fluctuations in the first place.

Primordial quantum fluctuations become magnified to super-horizon scales, that we see in the CMBR

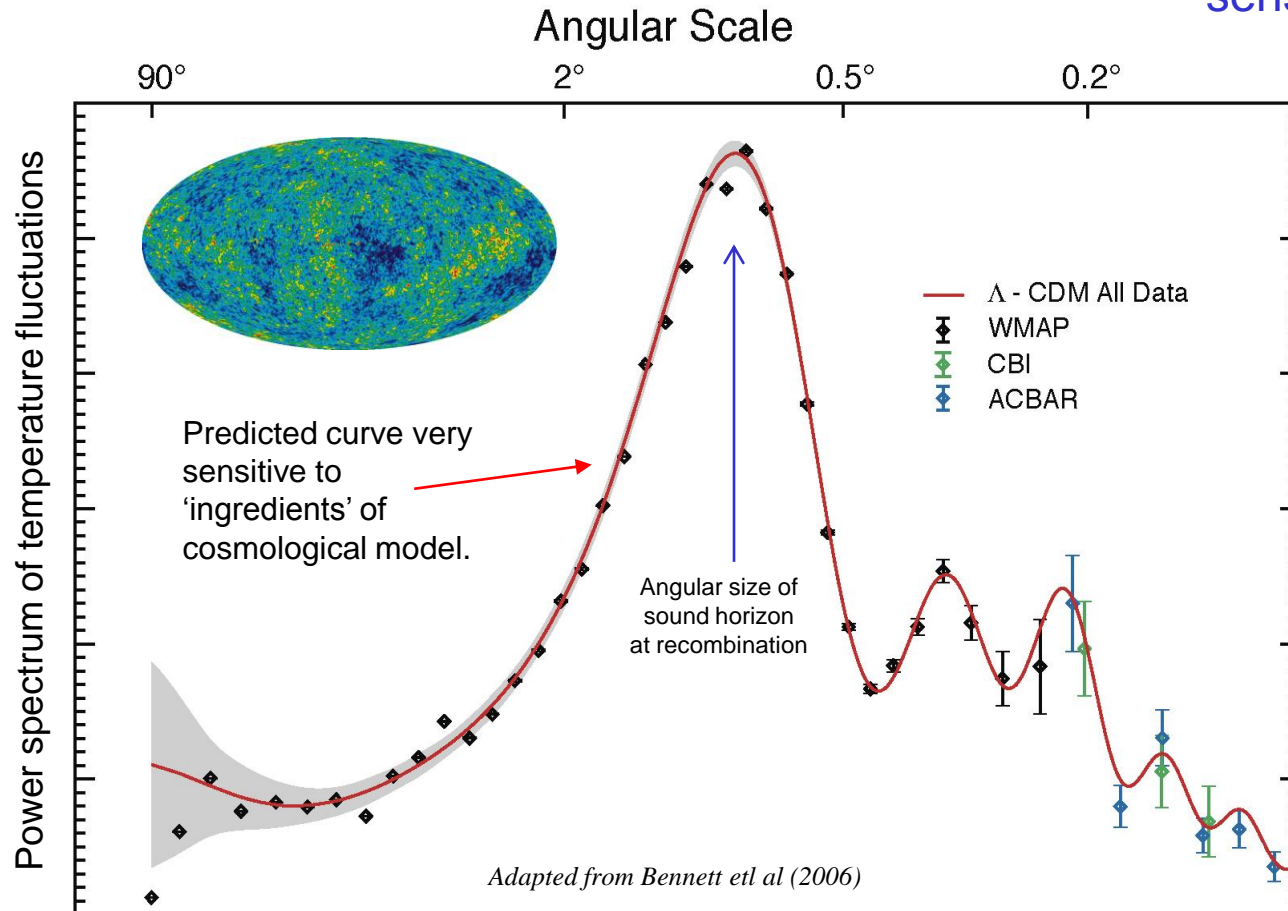


Simplified CMBR Power Spectrum



Cosmological parameters from the CMBR

Amplitude of Doppler peaks
sensitive to density of matter



Friedmann equation

$$H^2 = \left(\frac{\dot{R}}{R} \right)^2 = \frac{8\pi G\rho}{3} + \frac{\Lambda}{3} - \frac{k}{R^2}$$

For $\Lambda = 0$

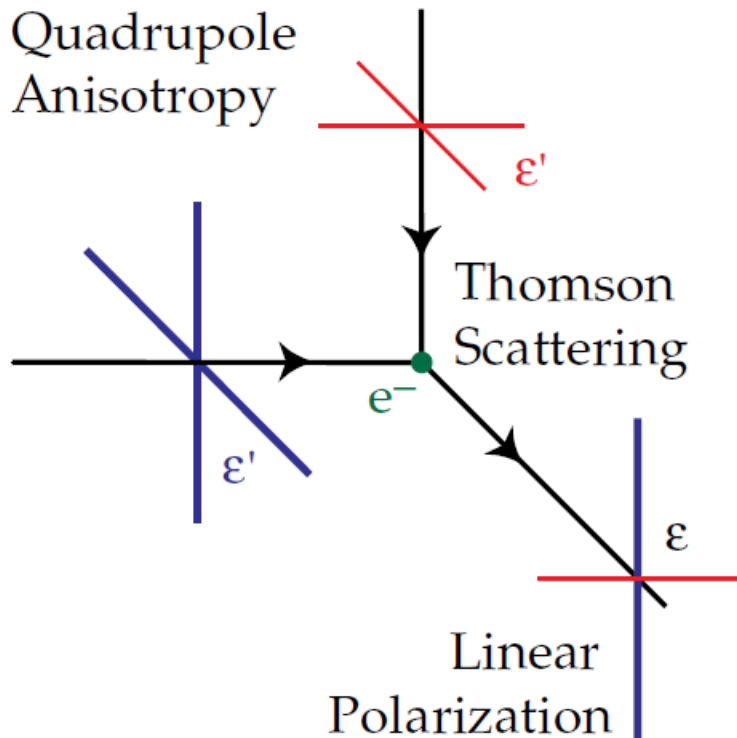
$$k = 0 \Leftrightarrow \rho = \left[\frac{8\pi G}{3H^2} \right]^{-1} = \rho_{\text{crit}}$$

$$\Omega_m = \frac{\rho}{\rho_{\text{crit}}} = \frac{8\pi G\rho}{3H^2}$$

i.e. $\rho \propto H^2$

CMBR spectrum tightly constrained but H_0 degenerate with other parameters

We can obtain information
from the CMBR
polarisation as well as
temperature maps.



arXiv:astro-ph/9706147v1 16 Jun 1997

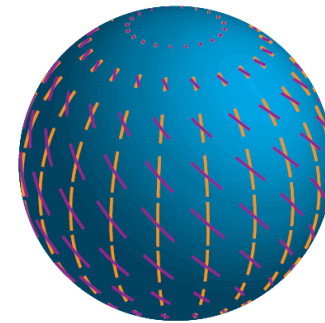
A CMB Polarization Primer

Wayne Hu¹

Institute for Advanced Study, Princeton, NJ 08540

Martin White

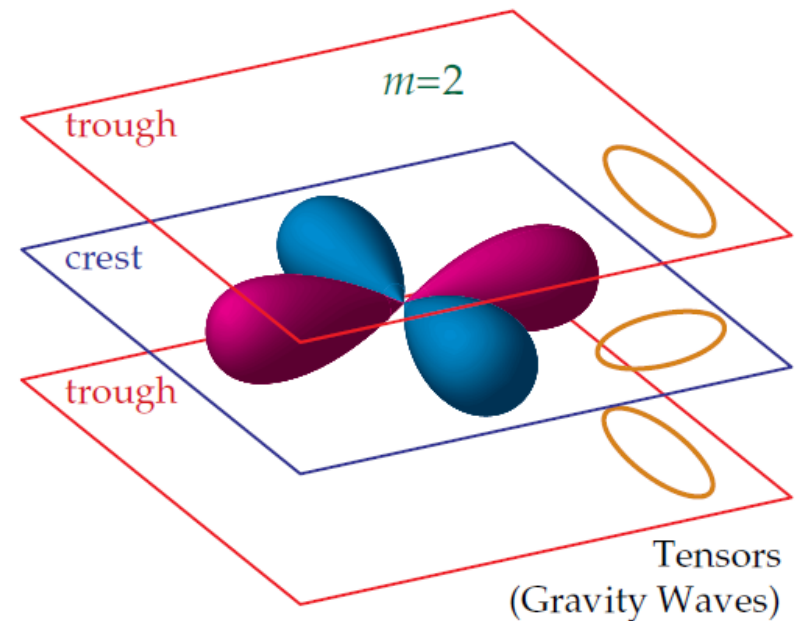
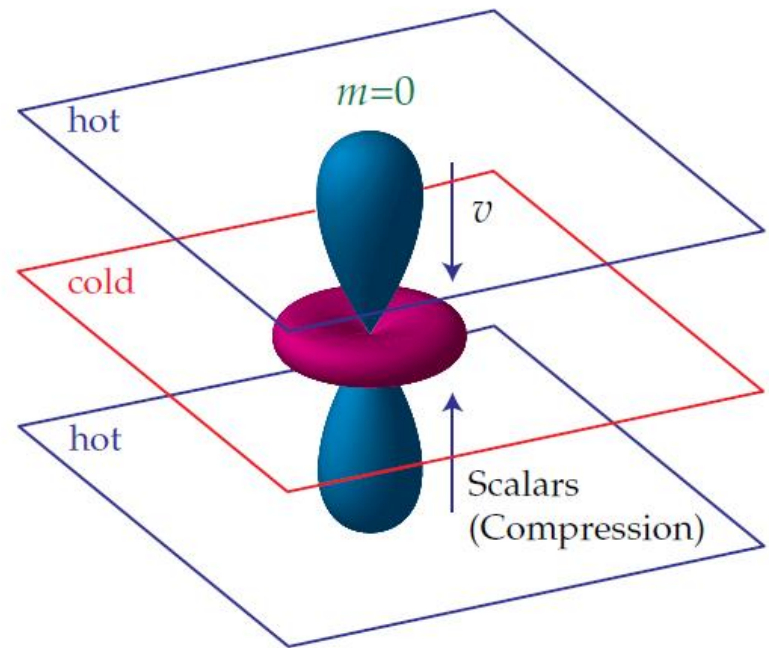
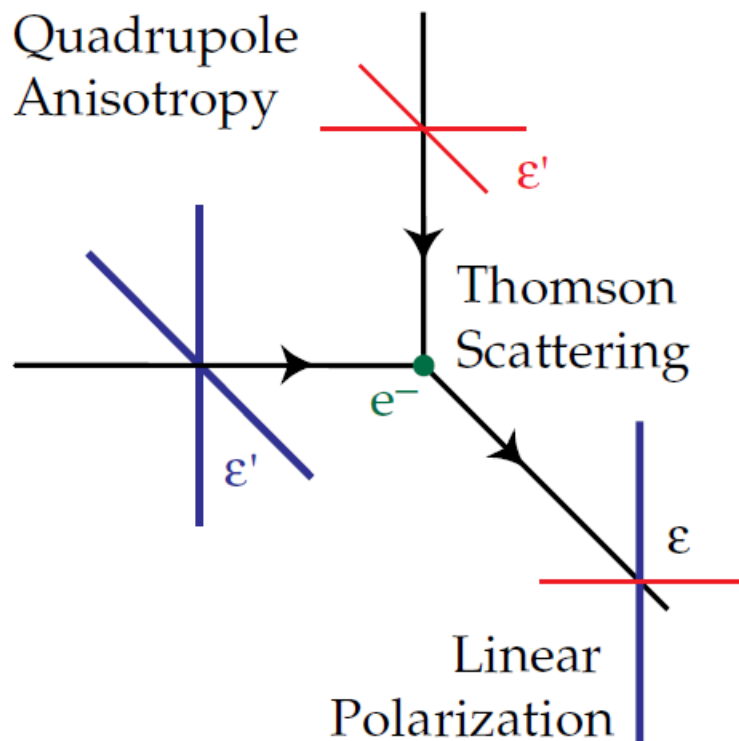
Enrico Fermi Institute, University of Chicago, Chicago, IL 60637



ABSTRACT

We present a pedagogical and phenomenological introduction to the study of cosmic microwave background (CMB) polarization to build intuition about the prospects and challenges facing its detection. Thomson scattering of temperature anisotropies on the last scattering surface generates a linear polarization pattern on the sky that can be simply read off from their quadrupole moments. These in turn correspond directly to the fundamental scalar (compressional), vector (vortical), and tensor (gravitational wave) modes of cosmological perturbations. We explain the origin and phenomenology of the geometric distinction between these patterns in terms of the so-called electric and magnetic parity modes, as well as their correlation with the temperature pattern. By its isolation of the last scattering surface and the various perturbation modes, the polarization provides unique information for the phenomenological reconstruction of the cosmological model. Finally we comment on the comparison of theory with experimental data and prospects for the future detection of CMB polarization.

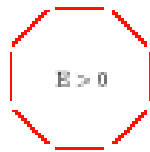
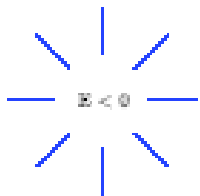
We can obtain information from the CMBR
polarisation as well as temperature maps.



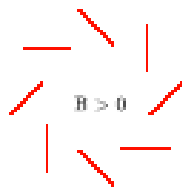
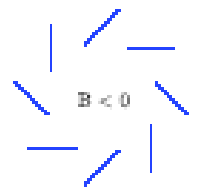
Because the Thomson scattering is anisotropic, the CMBR is **polarised**.

We can decompose the polarisation field into E and B modes.

$$P(\hat{n}) = \nabla E + \nabla \times B$$

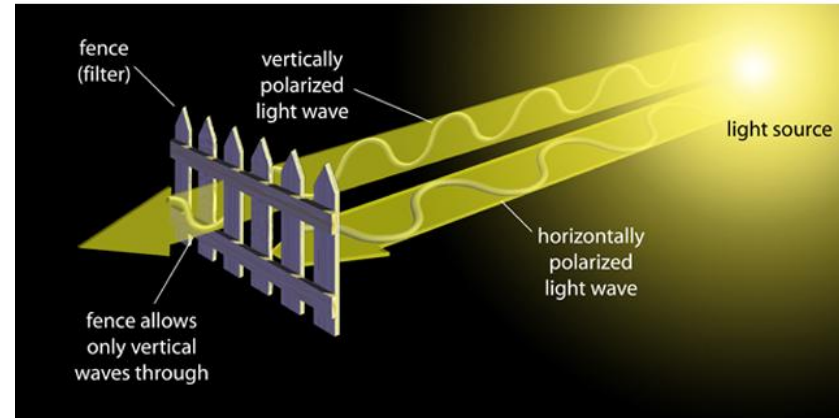


Grad

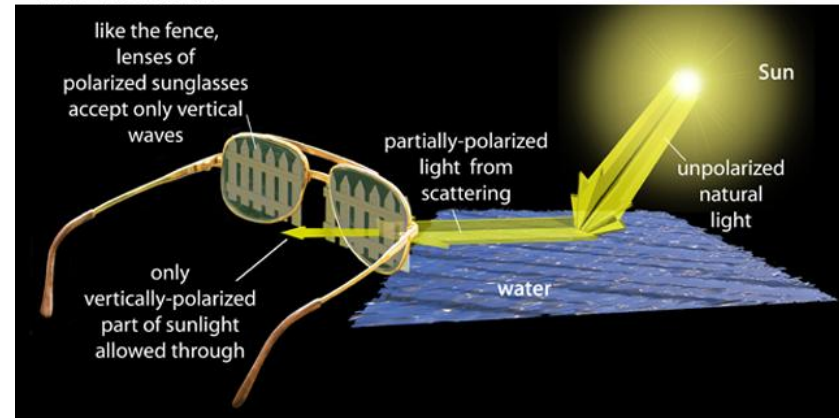


Curl

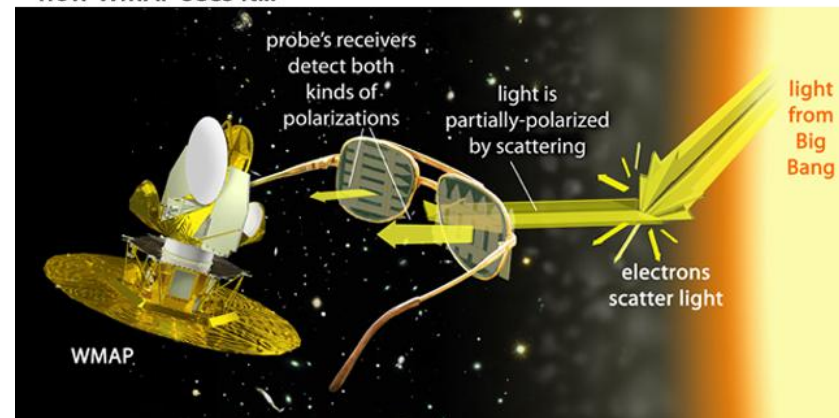
Polarization: How It Works



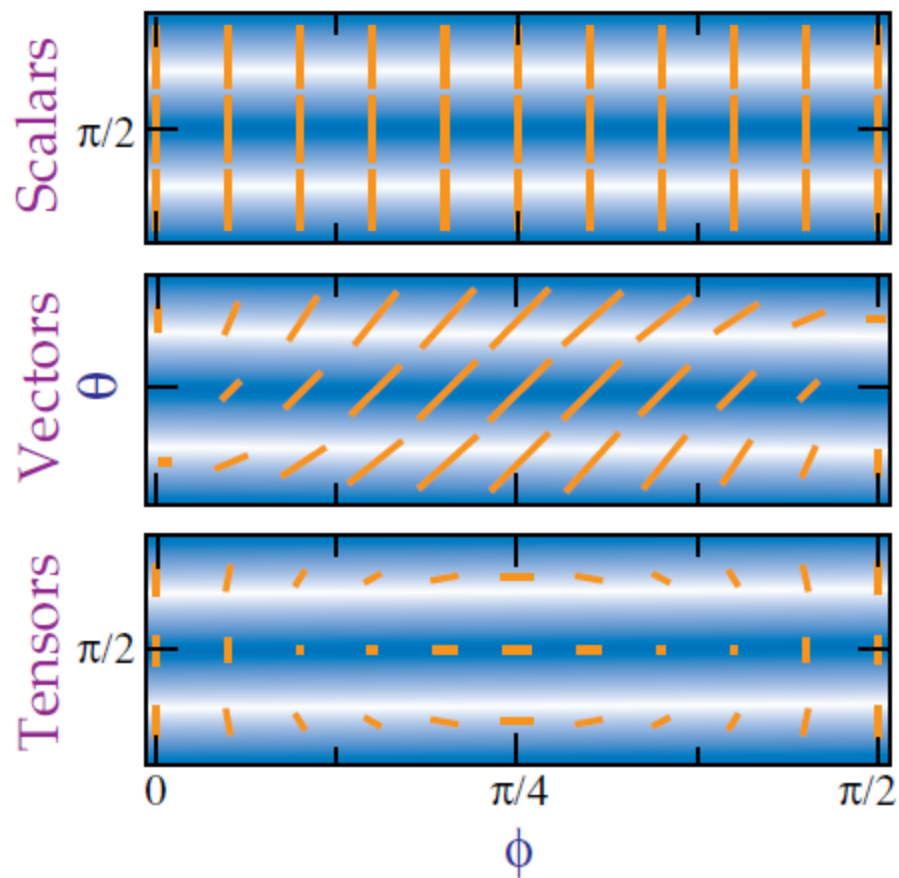
how we see it...



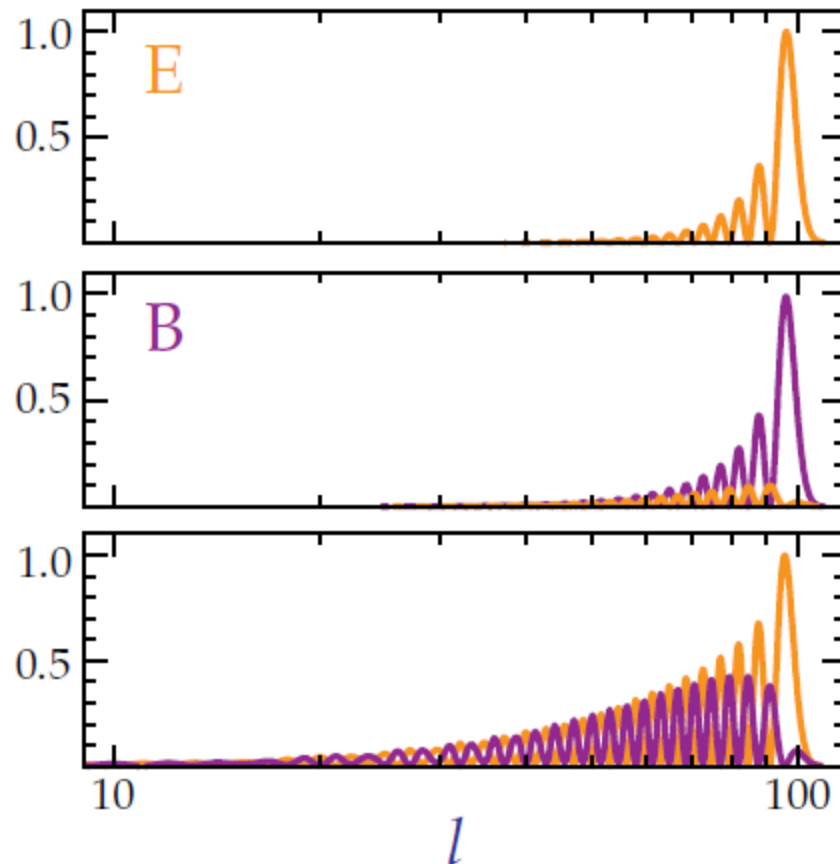
how WMAP sees it...

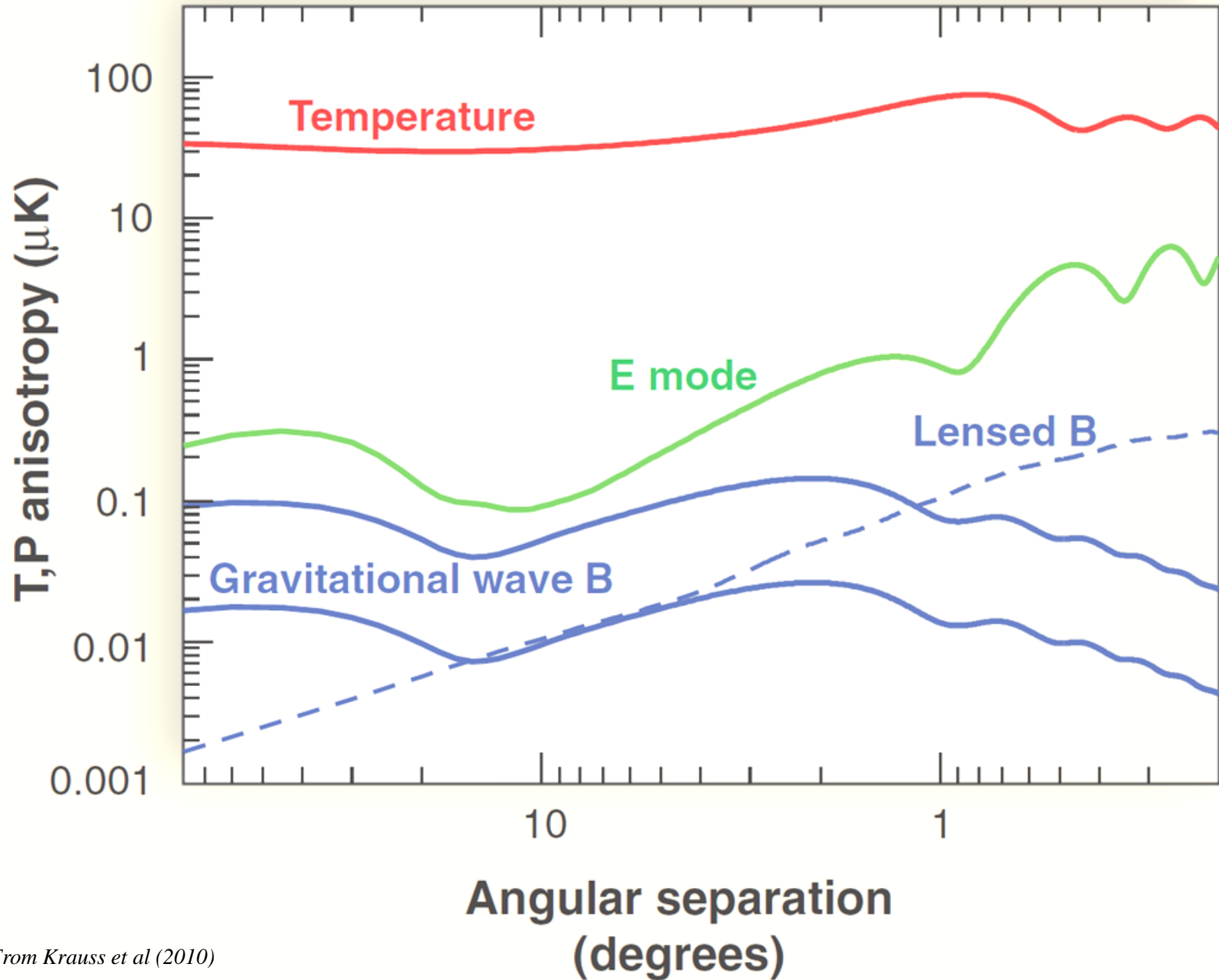


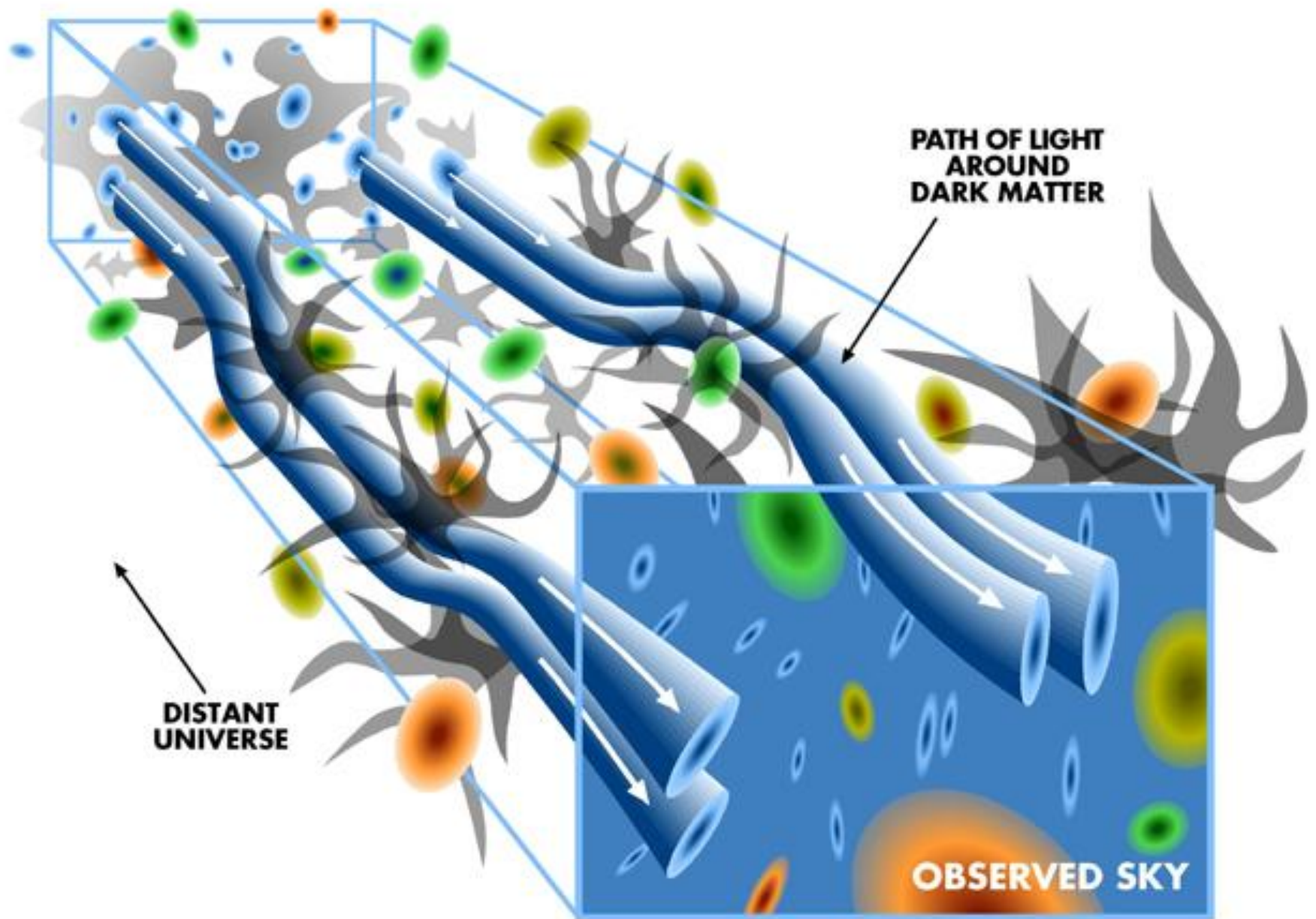
(a) Polarization Pattern



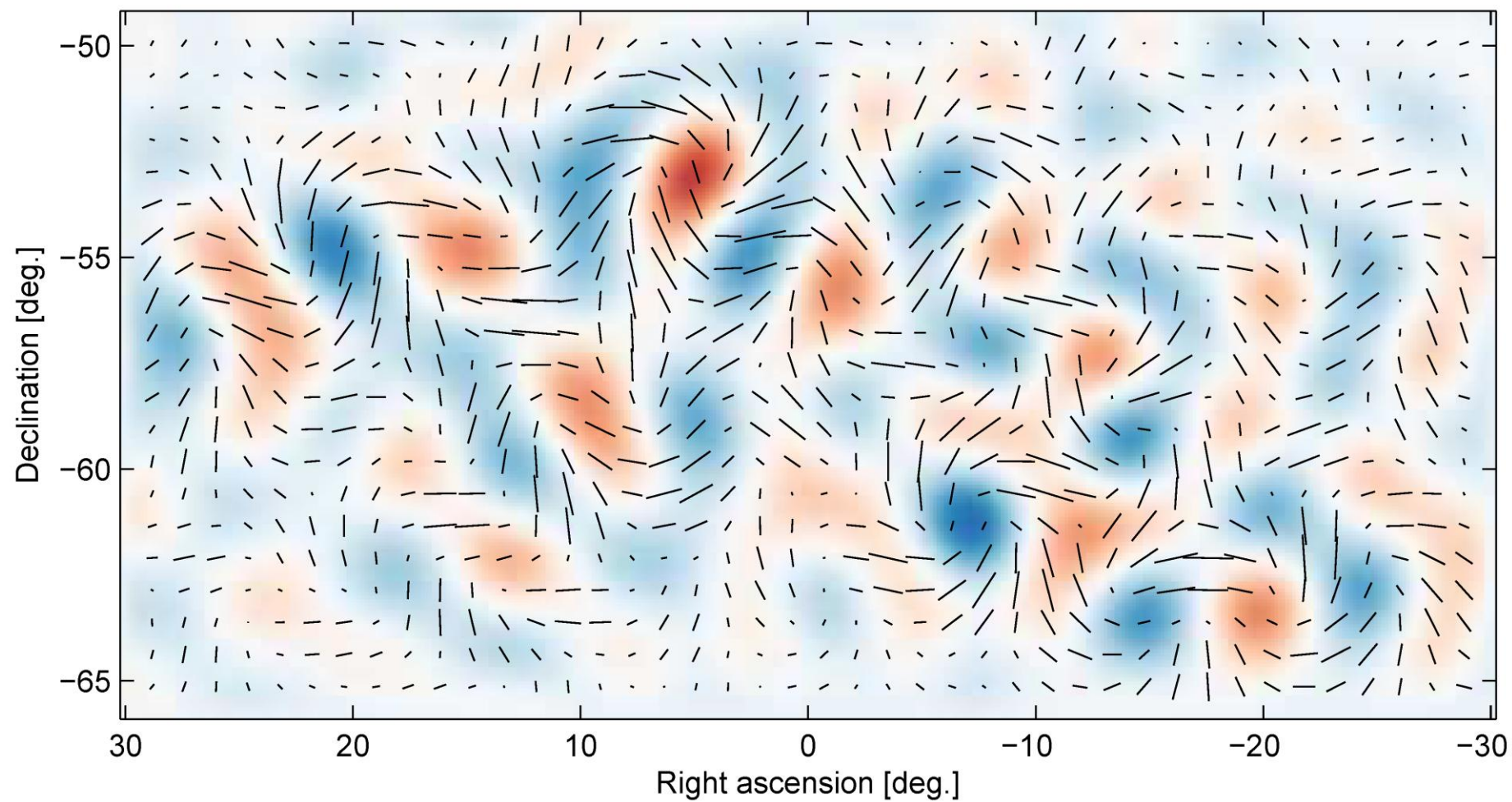
(b) Multipole Power



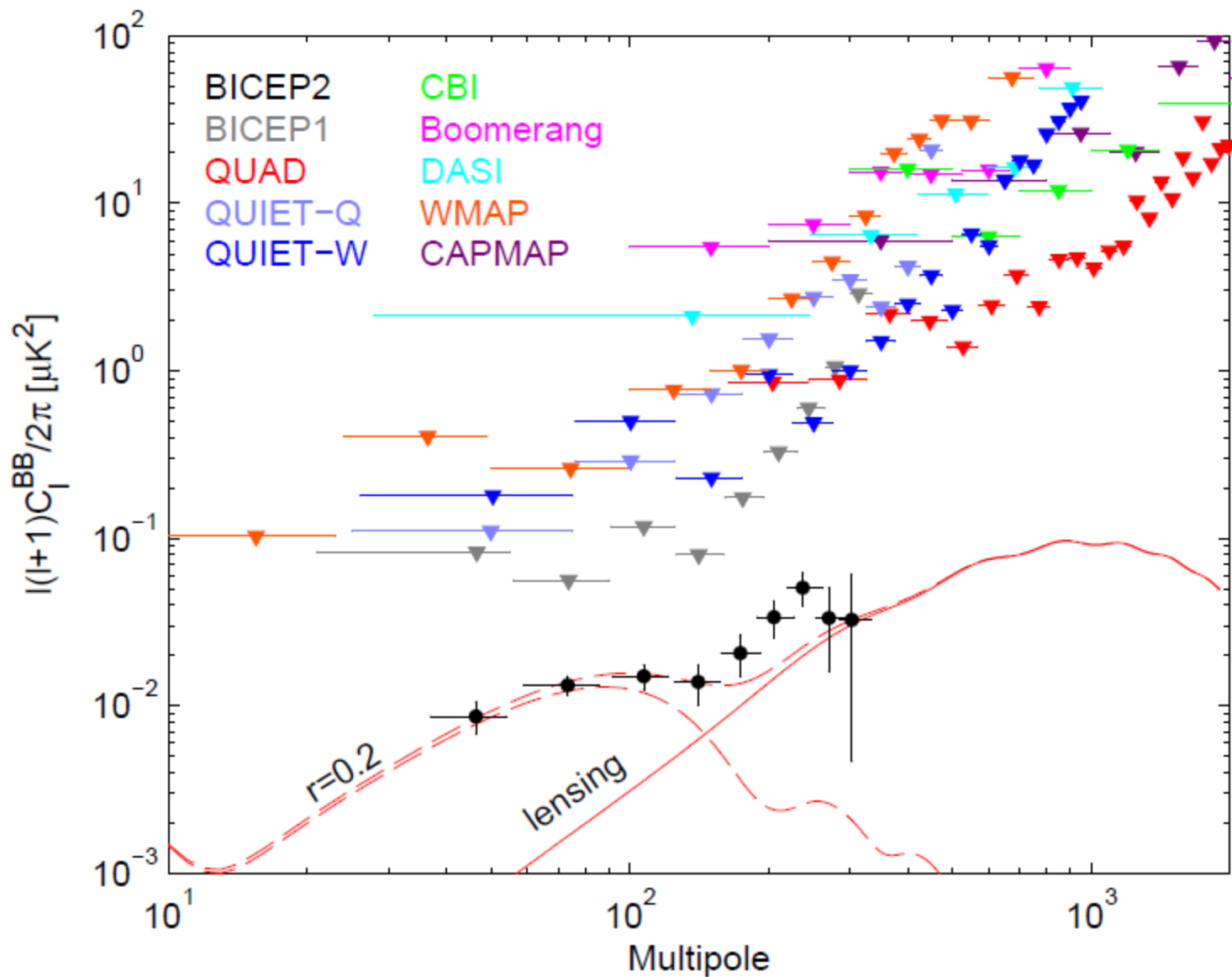




BICEP2 B-mode signal



Ade et al (2014) BICEP-2 results



What can we constrain with CMBR data?

It is usually a good approximation to take the power spectra as being power-laws with scale. So

Density perturbations $A_S^2(k) = A_S^2(k_0) \left[\frac{k}{k_0} \right]^{n_S - 1}$

Gravitational waves $A_T^2(k) = A_T^2(k_0) \left[\frac{k}{k_0} \right]^{n_T}$

Following Melchiorri (2008)

What can we constrain with CMBR data?

The 4 parameters are related to the inflaton potential and to its first two derivatives:

$$n_s - 1 \approx -\frac{m_{Pl}^2}{8\pi} \left(\frac{V_*'}{V_*} \right)^2 + \frac{m_{Pl}^2}{4\pi} \left(\frac{V_*'}{V_*} \right)'$$

$$n_T \approx -\frac{m_{Pl}^2}{8\pi} \left(\frac{V_*'}{V_*} \right)^2$$

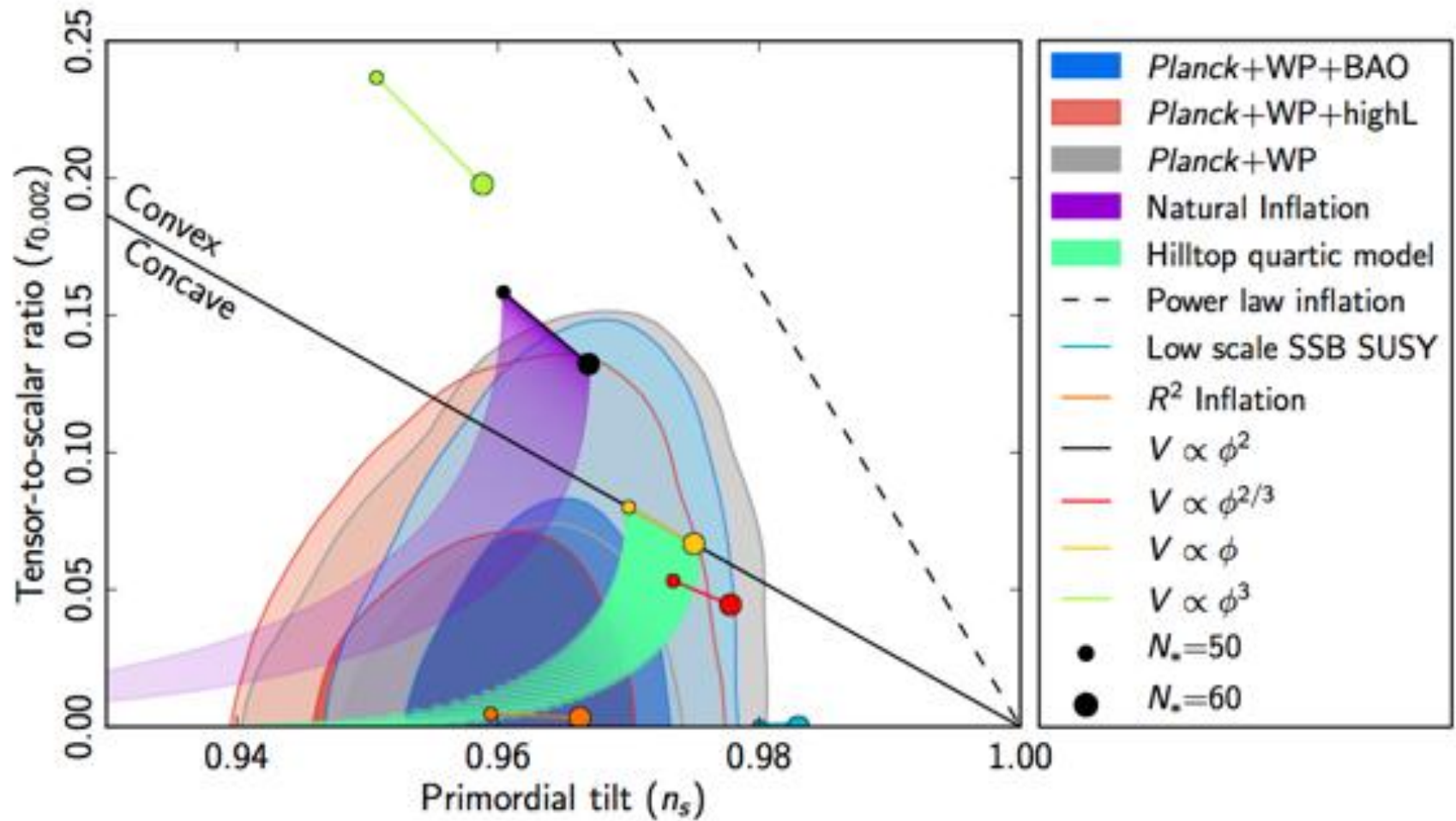
$$A_T \approx 0.61 \frac{V_*}{m_{Pl}^4}$$

$$A_s \approx -\frac{1}{7} \frac{A_T}{n_T}$$

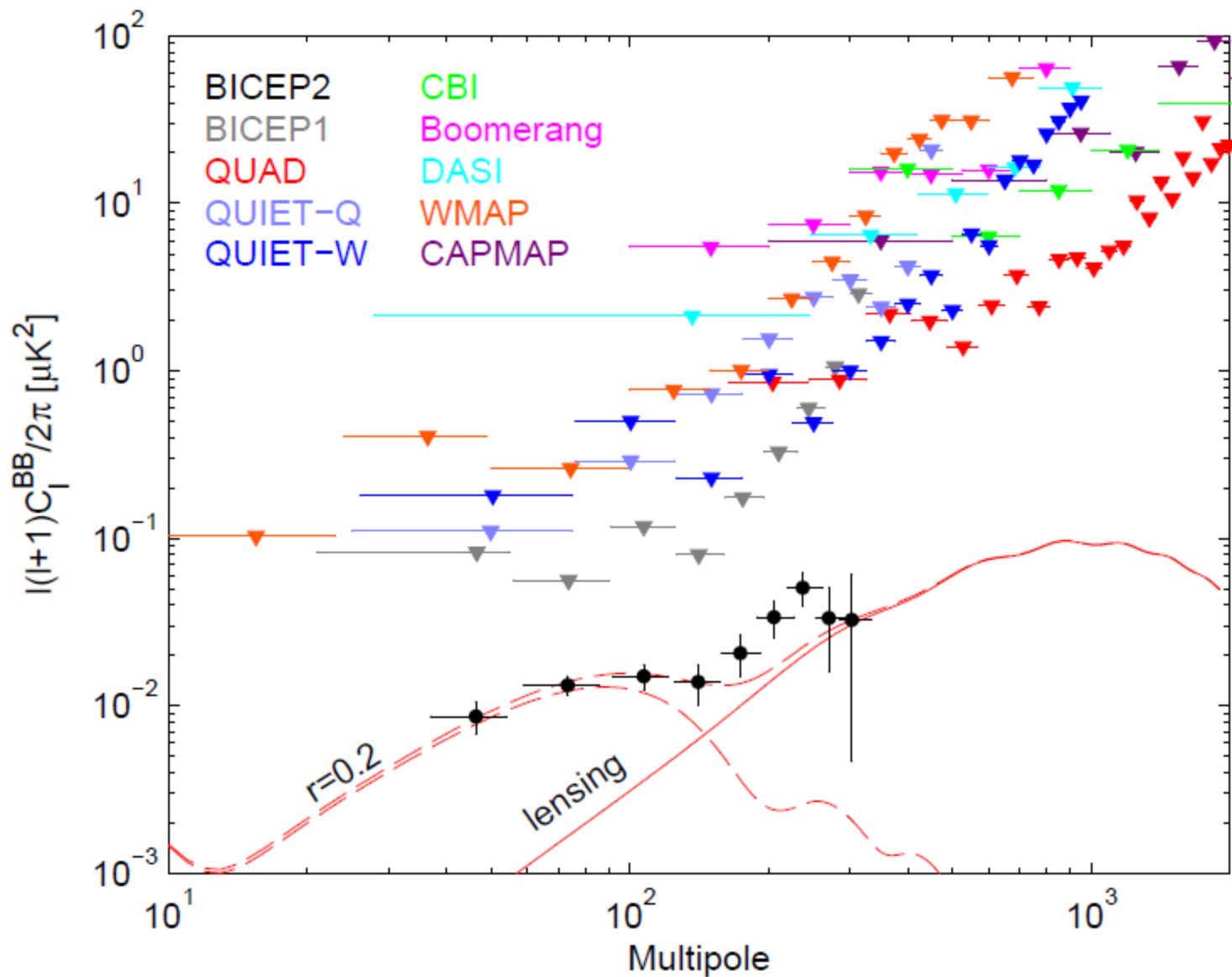
$$r = \frac{A_T}{A_s}$$

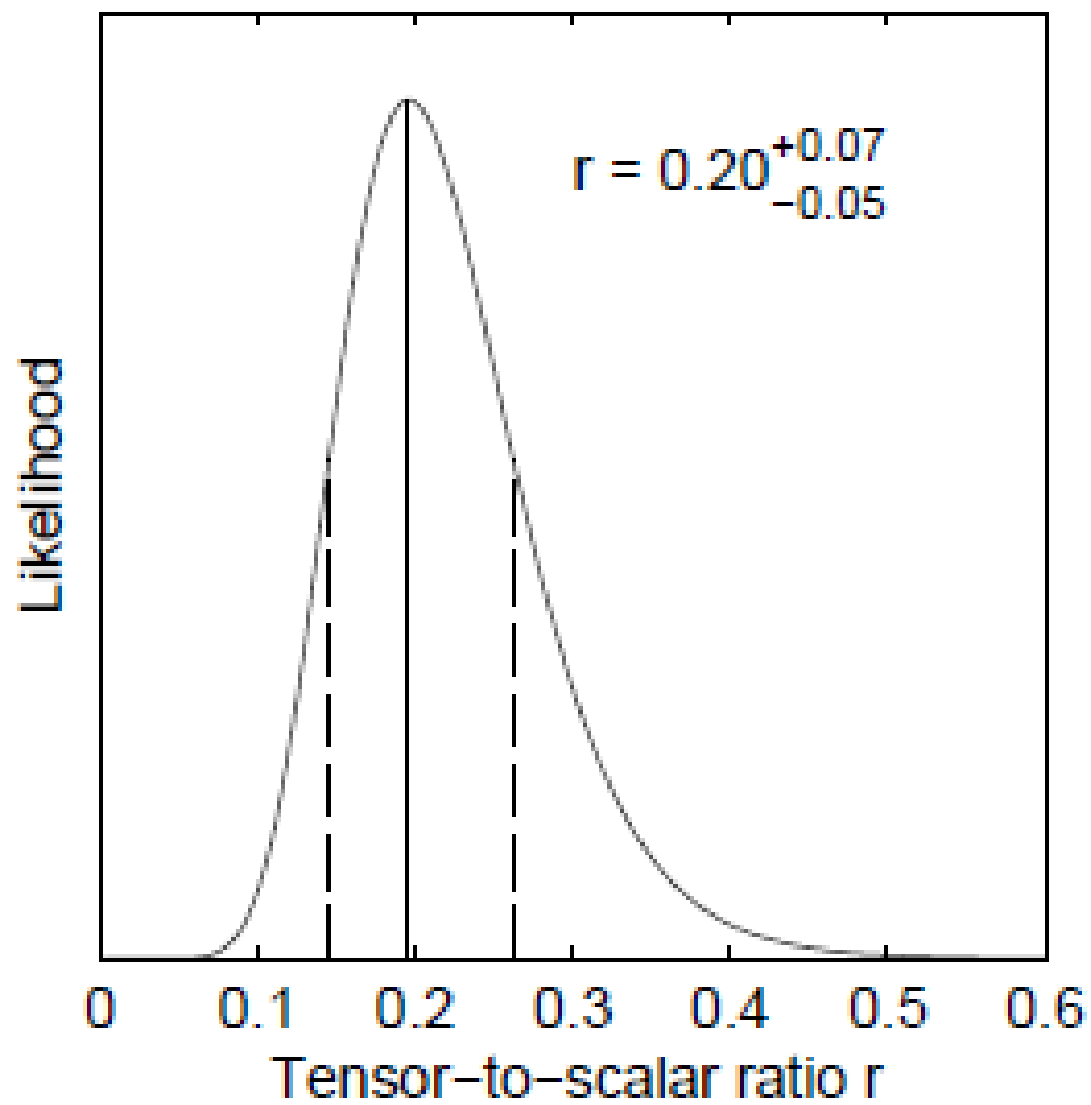
Following Melchiorri (2008)

Ade et al (2013) Planck results

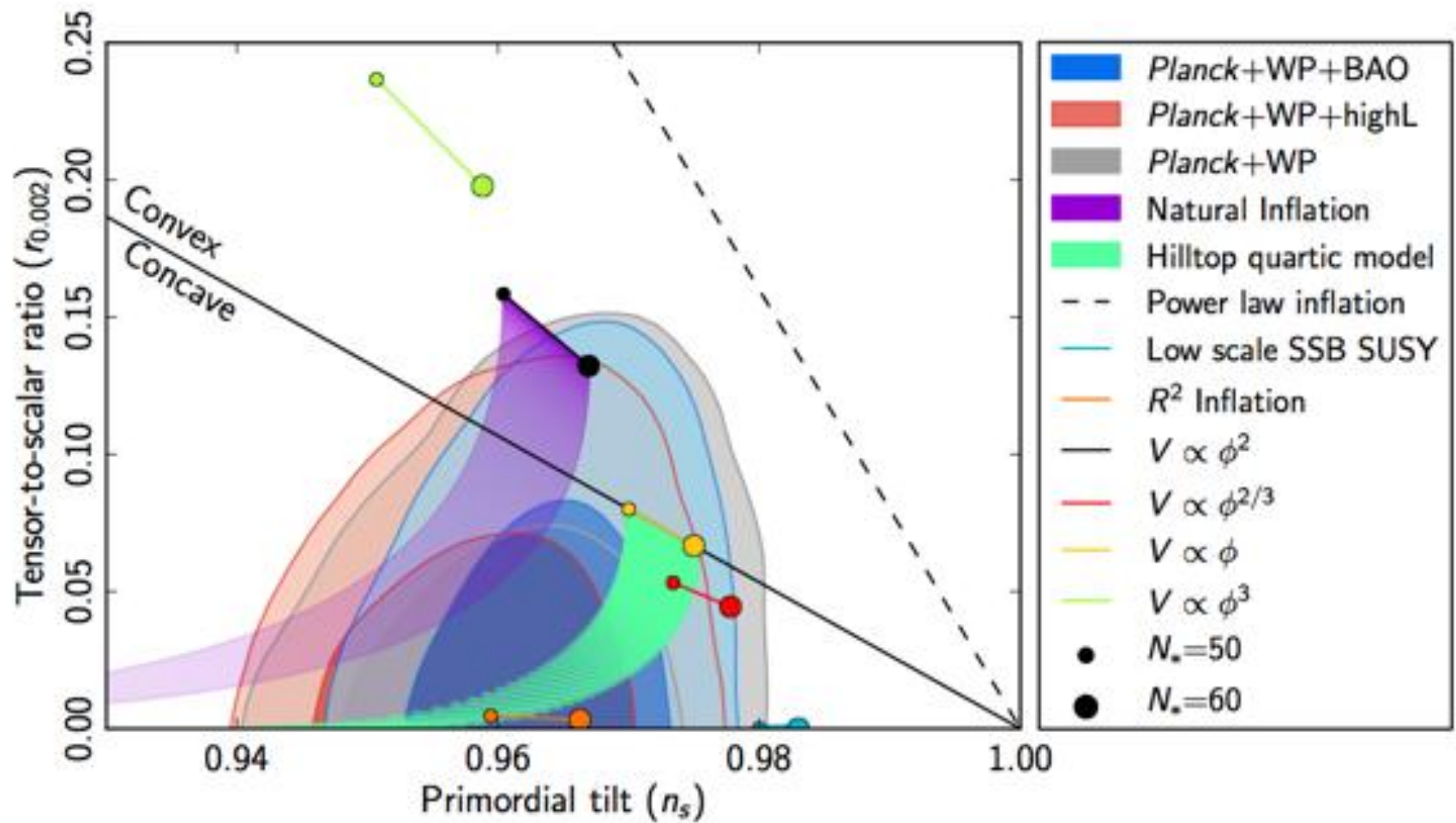


Ade et al (2014) BICEP-2 results

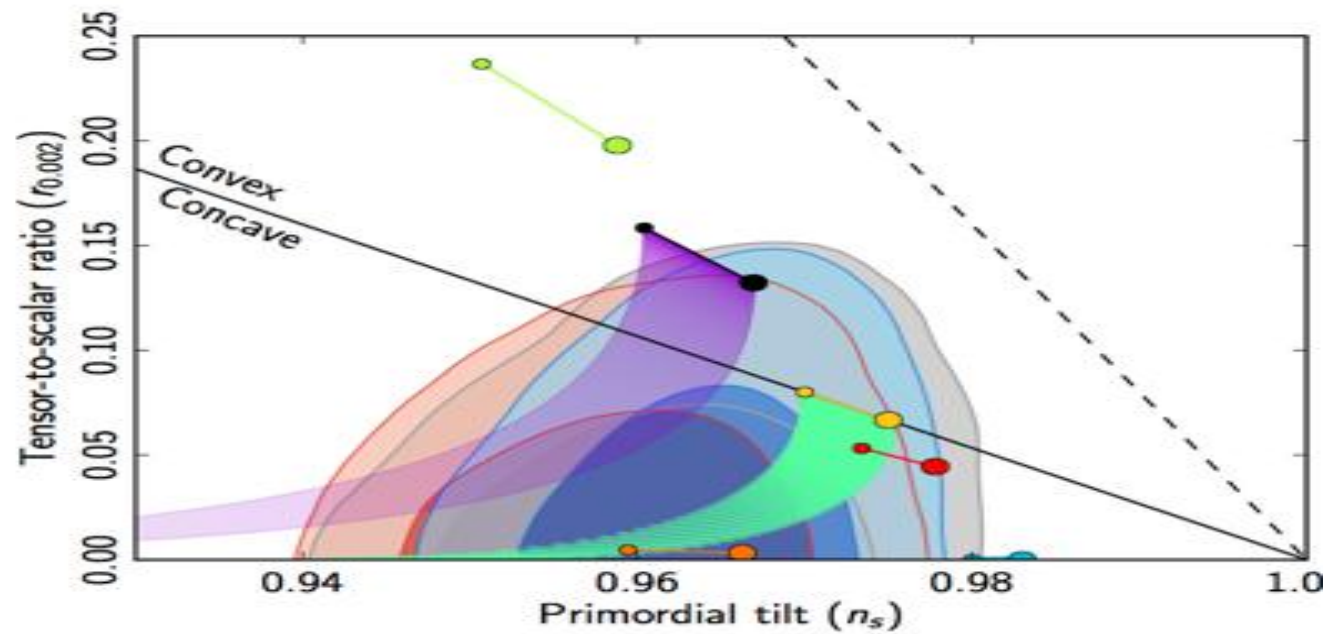




Ade et al (2013) Planck results

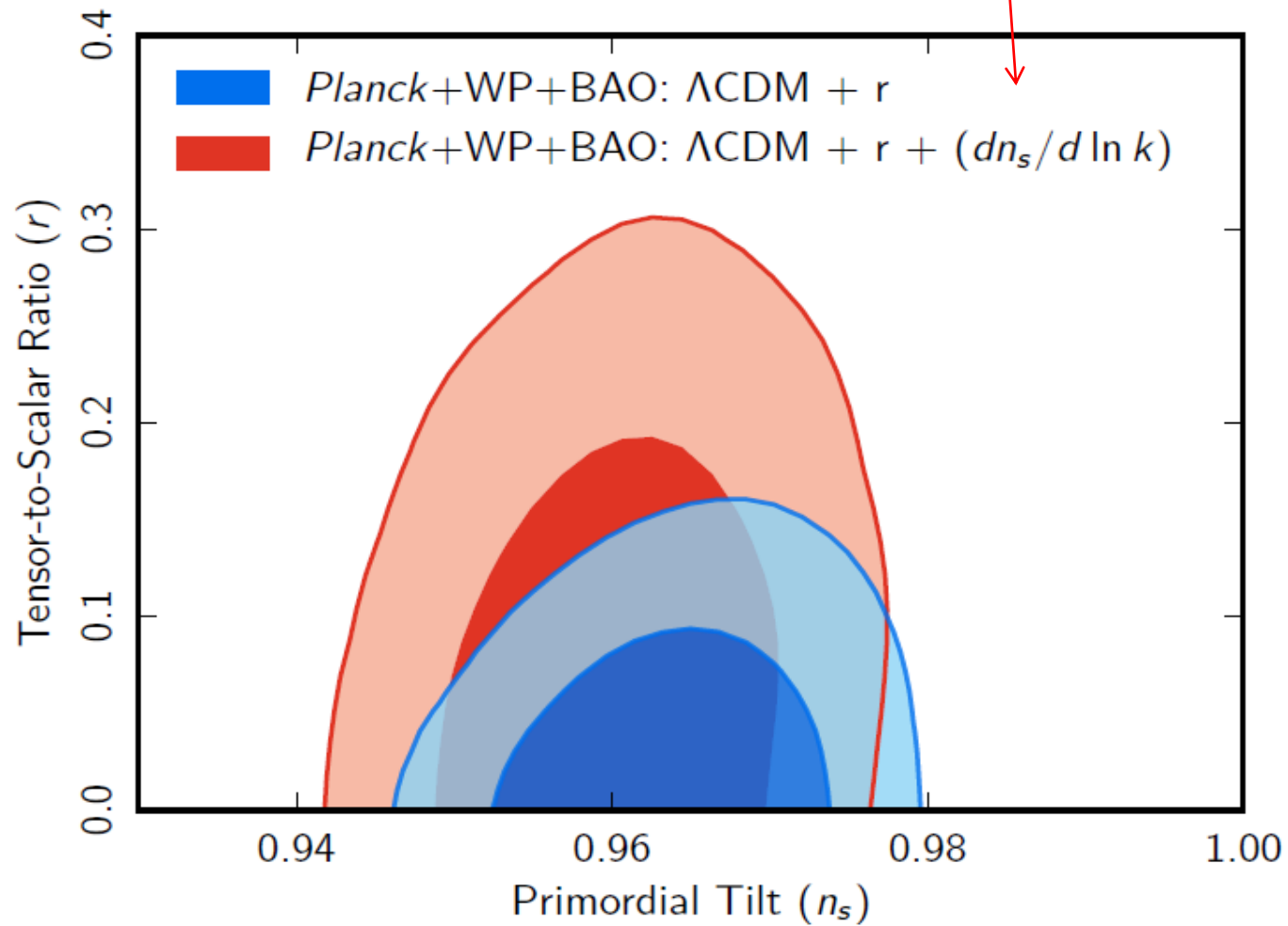


Ade et al (2013) Planck results

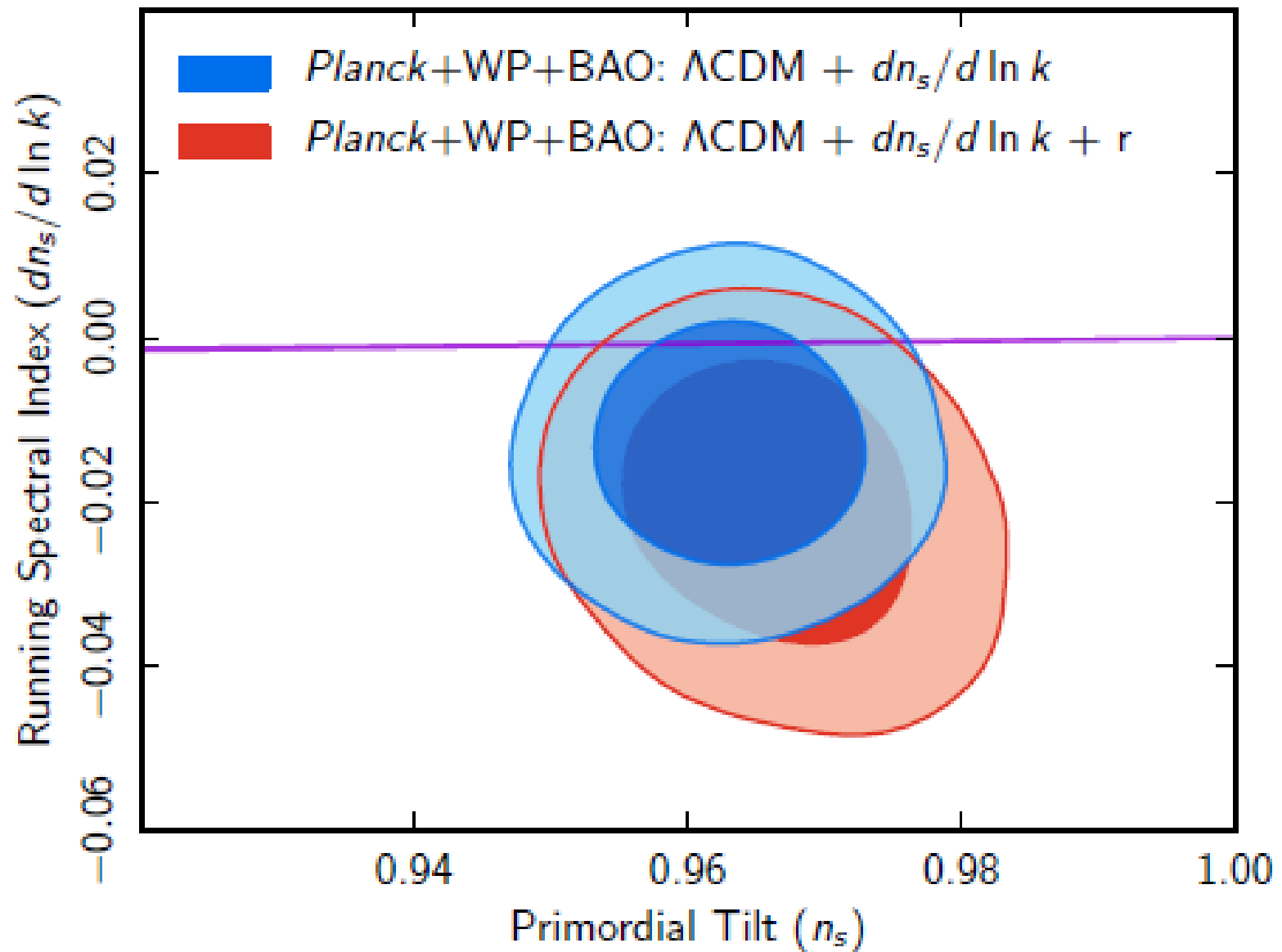


Ade et al (2013) Planck results

Allows for a running spectral index

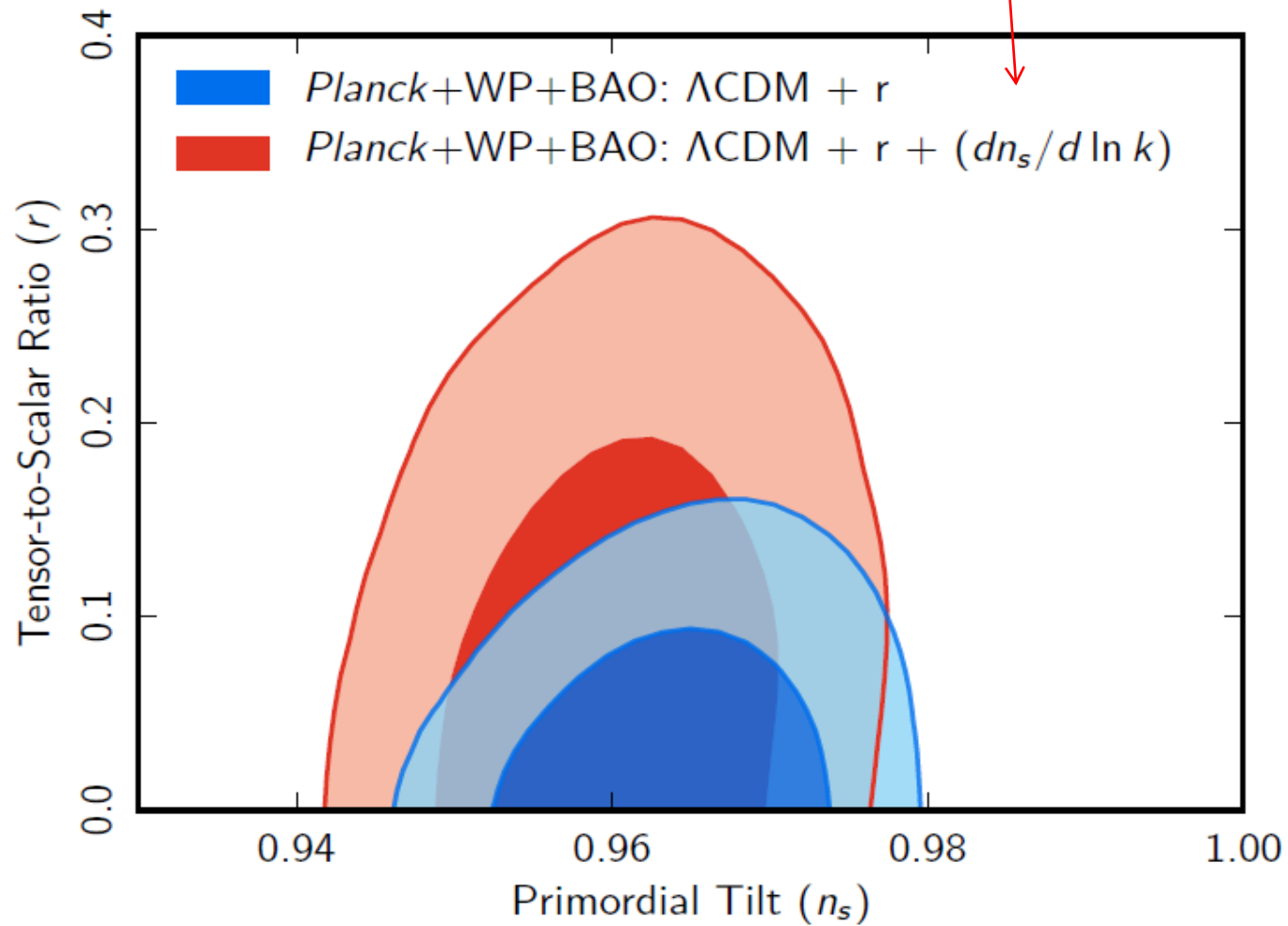


Ade et al (2013) Planck results

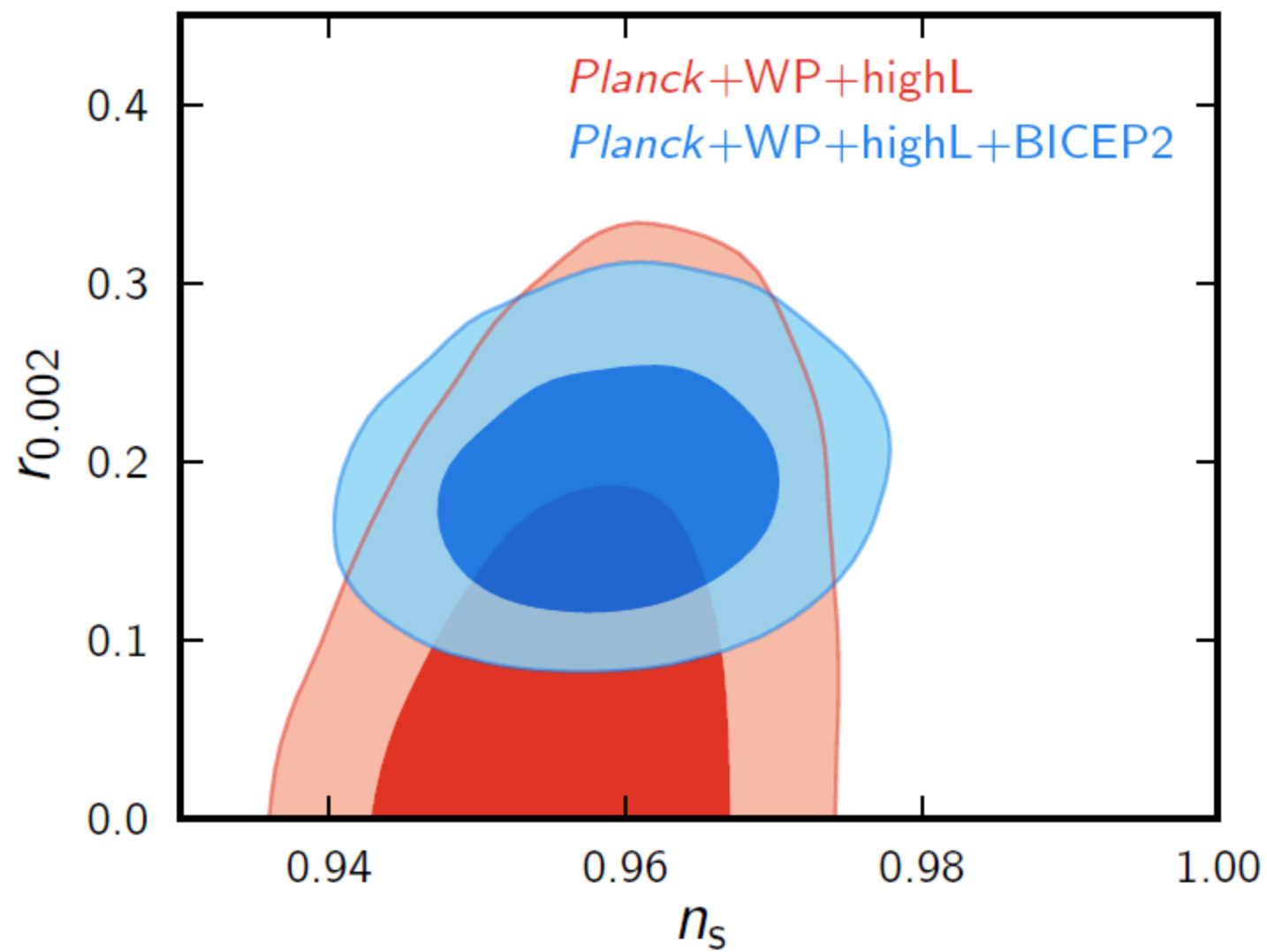


Ade et al (2013) Planck results

Allows for a running spectral index



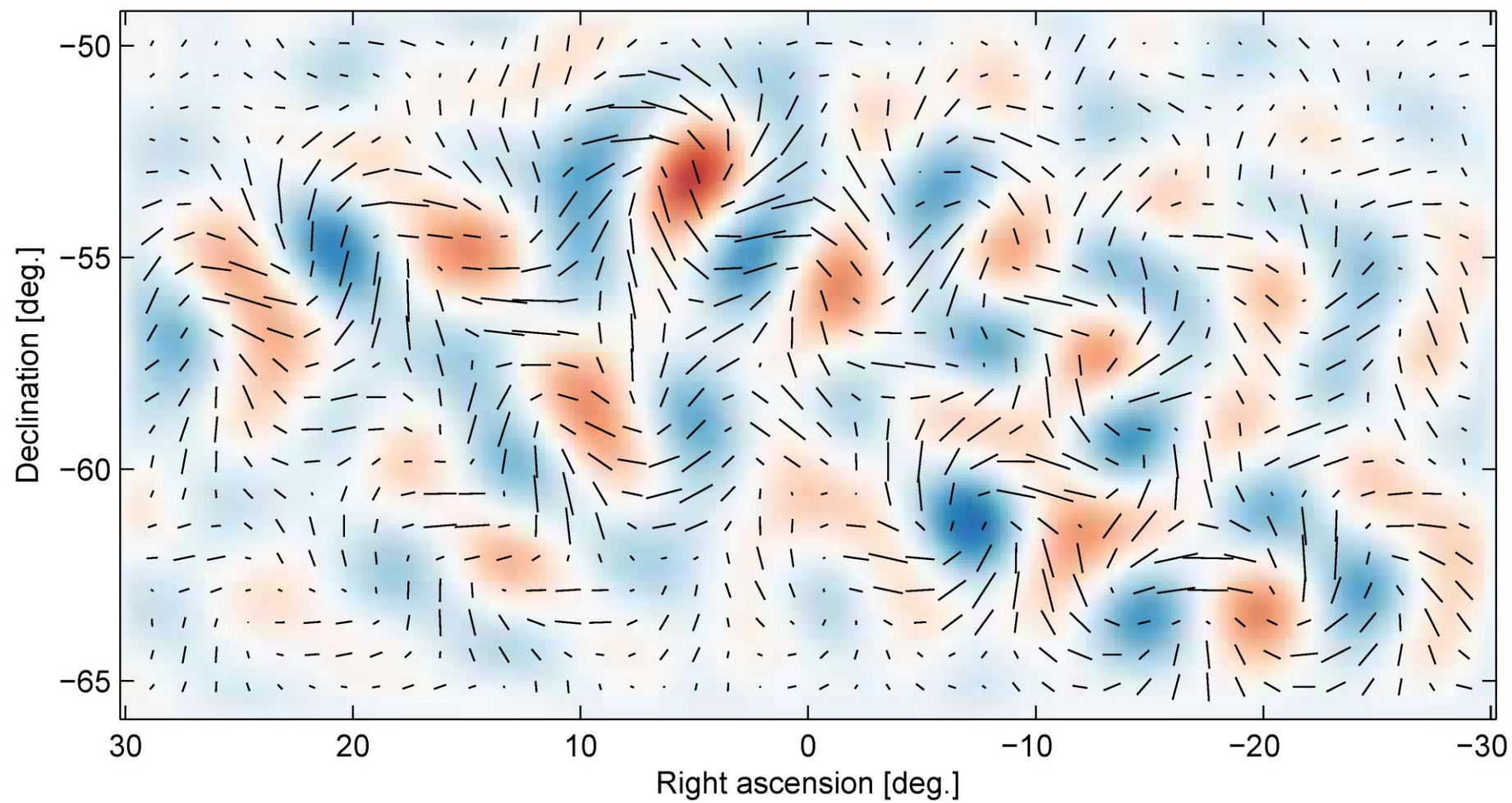
Ade et al (2014) BICEP-2 results



So...

- *Detection of B-mode polarisation looks secure. Is it primordial?....*

BICEP2 B-mode signal



So...

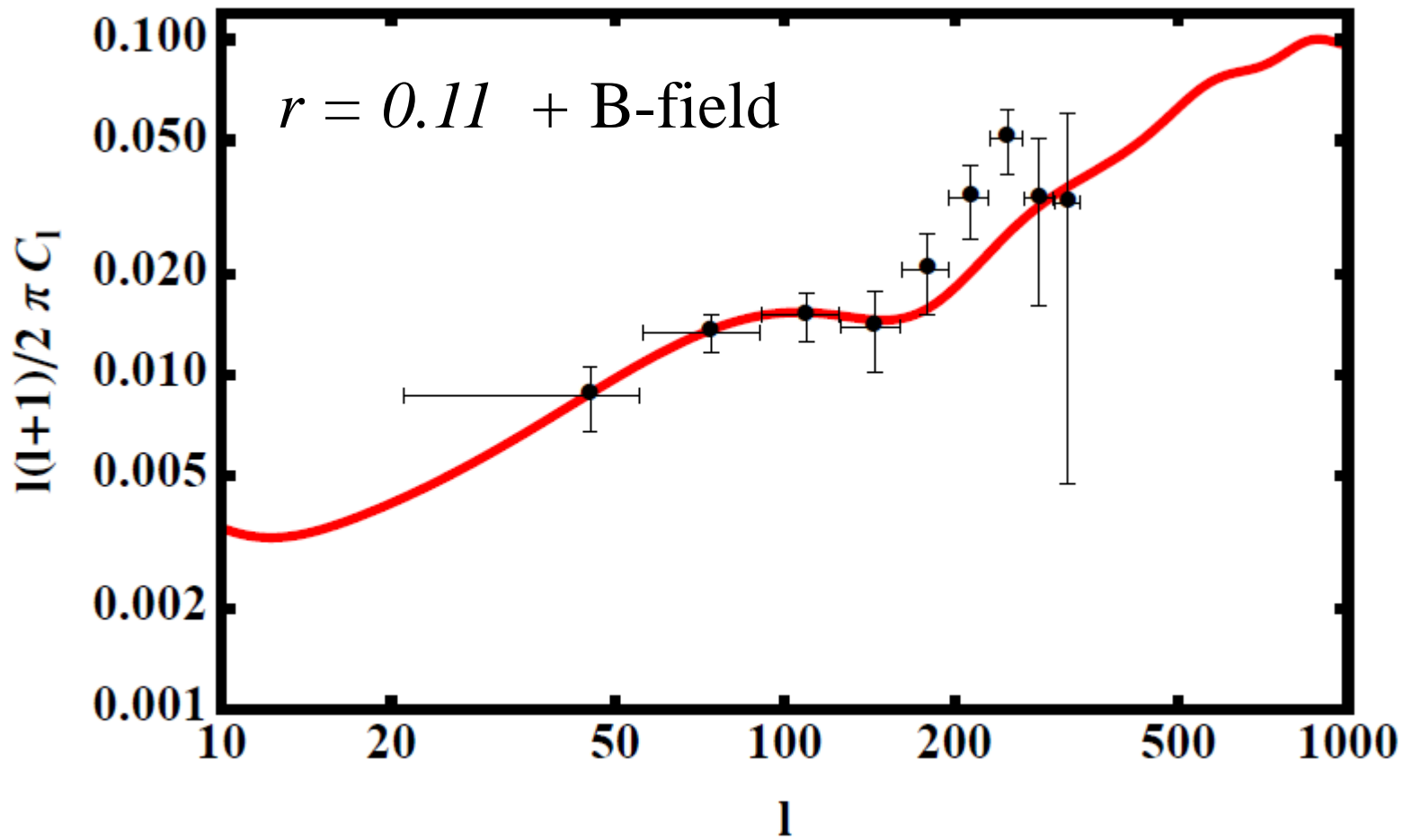
- *Detection of B-mode polarisation looks secure. Is it primordial?....*
- *Are we seeing gravitational waves from inflation?*

See Dent, Krauss & Mathur arxiv: 1403.5166

Isocurvature perturbations – not quite ruled out!

See Bonvin, Durrer & Maartens arxiv: 1403.6768

Some of B-mode signal from magnetic fields



Bonvin, Durrer & Maartens *arxiv: 1403.6768*



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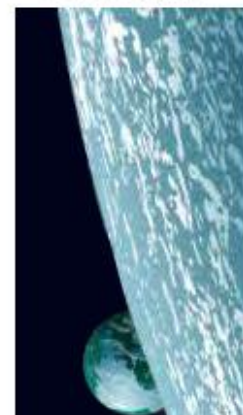
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Good Morning, Inflation! Hello, Multiverse!

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