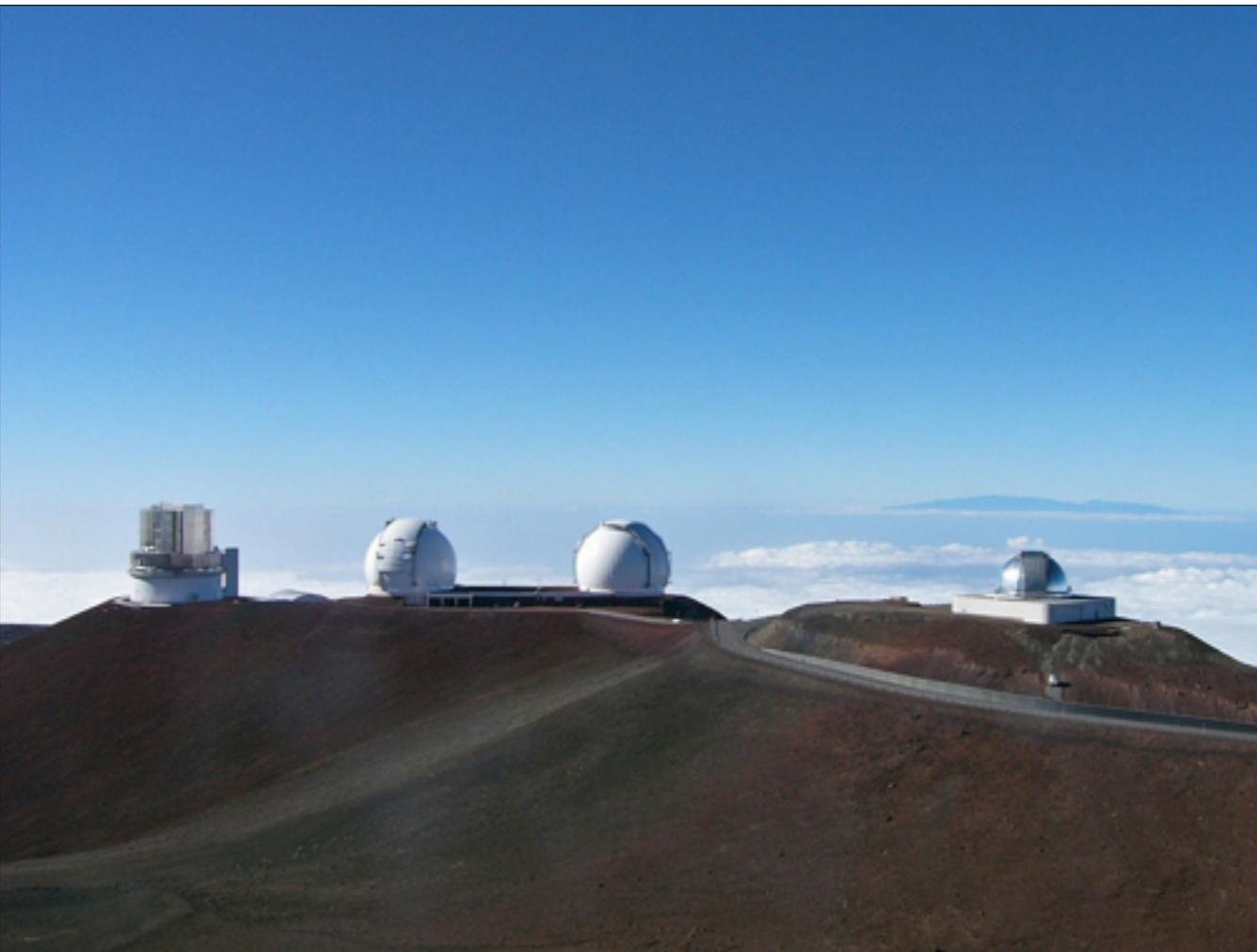


Classical mechanics from the mountaintop

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LEGES MOTUS

LEX I.

(*) Cioè sono percorrevate in state suo qualsiasi vel momenti uniformiter in direzione, non qualsiasi a tiribù impensis cogitare statim illam mutare.

PROJECTILIA perseverant in motibus suis, nisi quicquam a resistentiâ aeris retardanter, et vi gravitatis impellante ducantur. Tresque, cœsas partes coherentio perpetuo retrahunt sese a motibus rectilineis, non cessat rotari, nisi quicquam ab aere retardatur. Majora autem Planetarum et Constarum corpora motus suos et progressivos et circulares in spatio minus resistentiales factas conservant *dictum*.

LEX II

(c) Mutationem motu proportionaliter esse vi matrici impressae, et fieri secundum lineas rectas quod vis illa imprimitur.

Si vis aliqua motum quacunq; generet; dupla duplam, tripla triplam generabit, sive sinus et sensu, sive gradus et successore impressa fatur. Et hic motus fractionis in eundem sensum plasmon cum vi penerat.

(5) 24. Et hic primi legi quae (5) de-
stributio, nequitas omnes, modus esse natura-
lis et arboribus et mortaliis, adiutorum ex illis
victimam remansit, nec directione mortali-
bus aliq[ue]d distractio modi offenserat; Unde
etiam in aliis modis non solum omnes amicis,
quosdam et aliq[ue]s legatos retribuerantur.
Cum autem exp[er]iencia propter eum
tempore defensio, et aliis super aliis
tempore superponerentur, inveniebat, et in gra-
tia etiam deinceps magis organicas, neque
etiam solum corporis, sed etiam partis, quae in his
ostentabat superponere continuo, aliorum, ac
principali que major vel minus erit modi restitu-
tio, etiam res vel minus decessivum accipiente
corporis projecti voluntatis. Et hic iustus p[ro]p[ter]ea

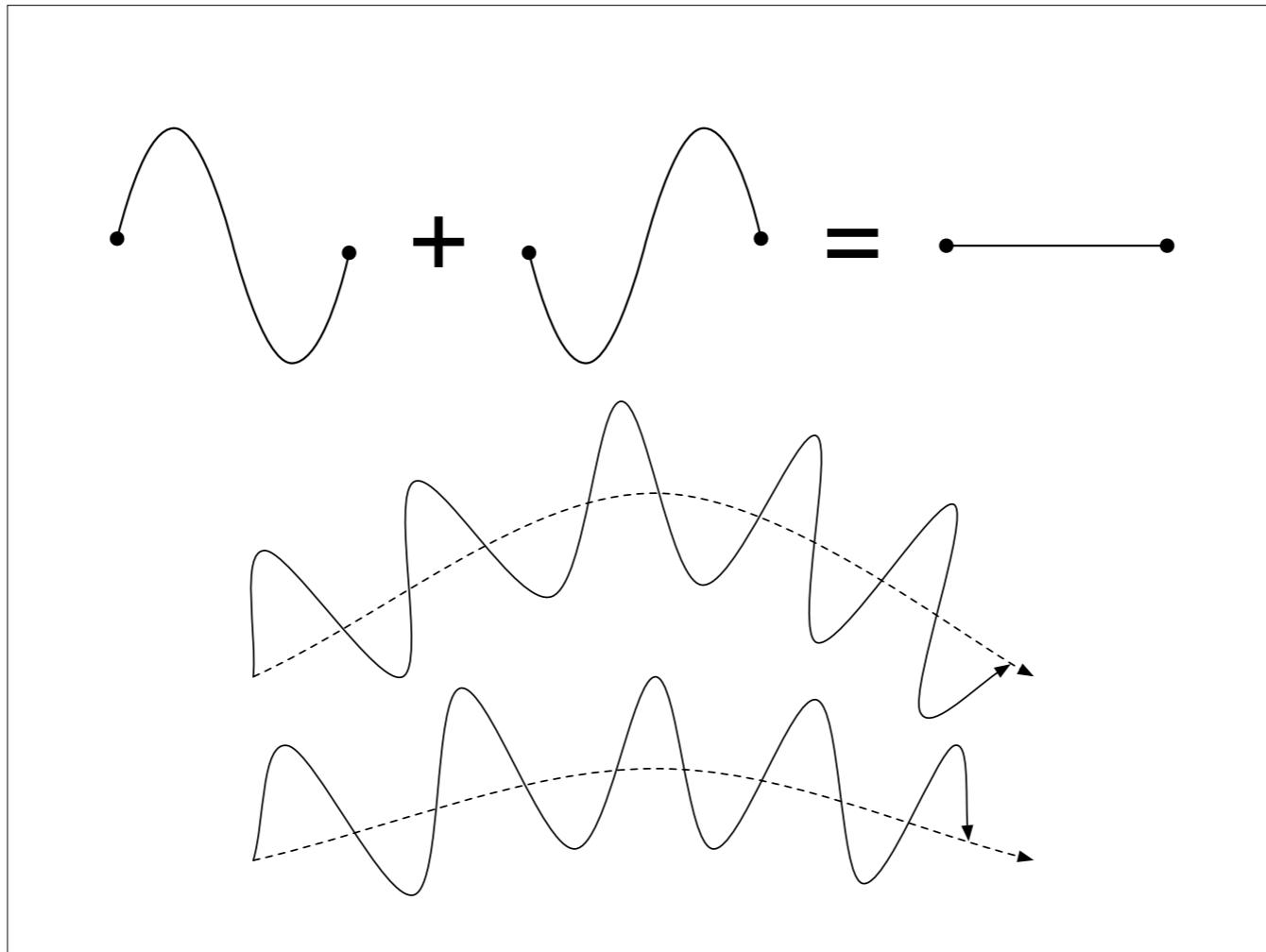
atricie deterministar) si corpus antea movebatur, motu ejus vel consipi-
tanti additur, vel contraria subducitur, vel oblique oblique adiicitur, et
cum eo secundum utrinque determinationem componitur.

(*) LEX III.

Actiones contrarias semper et aquales esse ratione: sive corporum duorum actiones in se mutuo semper esse aquales et in partes contrarias dirigi.

Quicquid premit vel trahit alterum, tantandis ab eo premitur vel trahitur. Si quis lapidem digitis premit, prematur et huius digitis a lapide. Si equus lapidem fusi alligatum trahit, retrahetur etiam et equus (ut ita dicimus) reseparatur in lapidem: nam fusionis stricte distantes eodem relaxandi se comatu urgebit equum versus lapidem, ac lapidem versus equum; tantumque impedit progressum unius quantum promovet progressum alterius. Si corpus aliquod in corpora aliad inspingens, rotam ejus vi sali quoniodocesque mutaverit, idem quoque viciniam in motu proprio eandem

— Corpse grisea in terra vicinia, subtili-
meli residuum, nata uniformiter aculeata
dissidens, et mato uniformiter rotundata
serruant. — Domestica Subtili meli
residuum idem ac cyprius corpus ponens
sive eodem illis in subtus plenum possum
tum in vertice, non in radice monte; et ante
ponens, seu sit monte (33) est minus in vige-
sime, et minus in radice. — Corpse grisea
corpora minus radice in vertice et in radice
monte permanens, manebit sicut radice vix
se levigata gravitate. — Corpse grisea corpo
ria in radice et vertice monte aquila spuma
hinc temporibus percurrent, solida et solidata
est, acutissimamente, sicut est experimenta
(12); ostendat non ligare ex levigatis, et per
difficile est, quod regis mato levigatis
ex levigatis, et levigatis mato levigatis
resistit (34). M. A. 1.



Fermat's principle

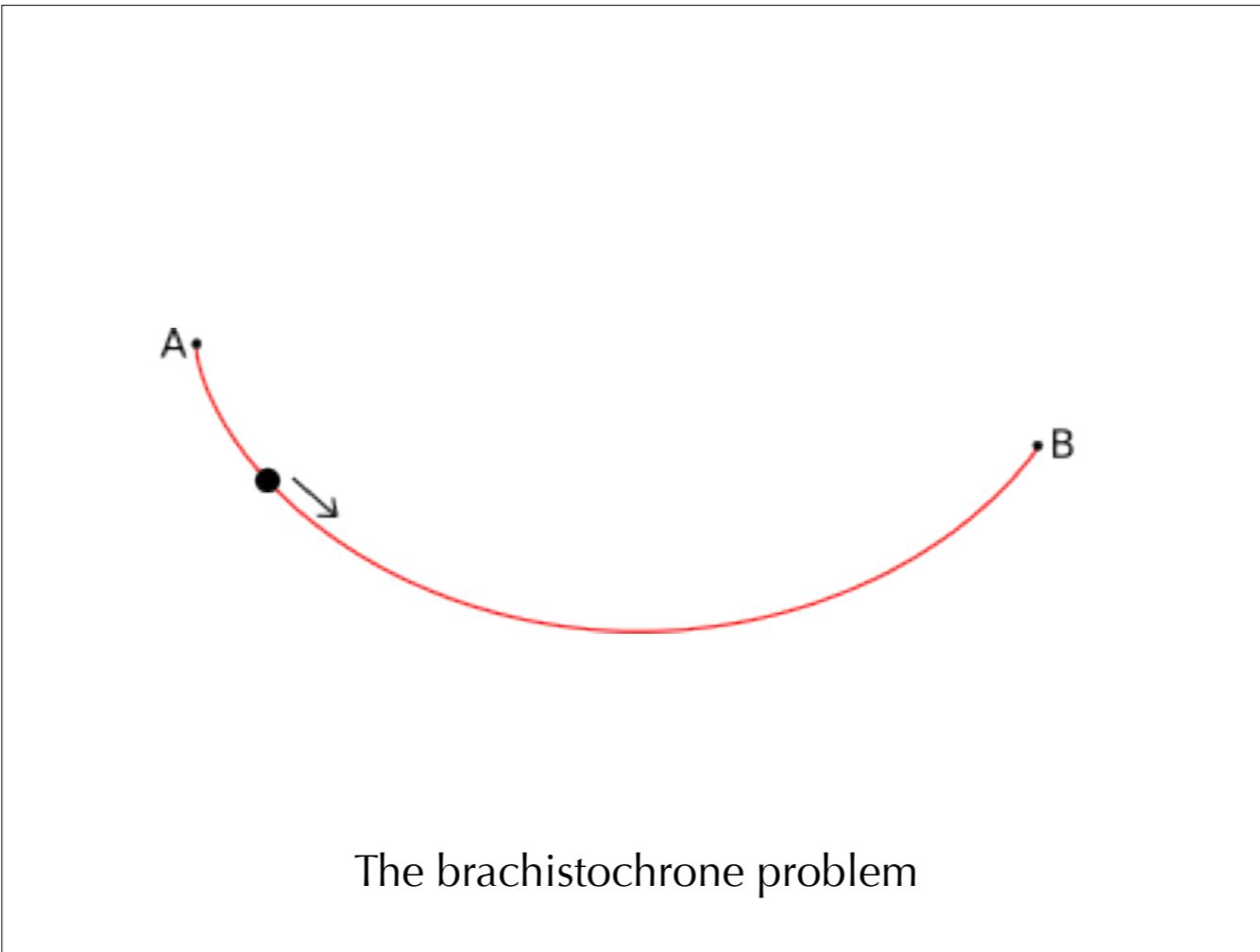


Travel time gets shorter as you go from the first to the second,
then it gets longer.
Ie minimum at the shortest path



Catenary

Can describe it in terms of balance of forces
but also minimising total potential energy



The brachistochrone problem

Brachistochrone problem

Very similar to the tautochrone problem, solved by Huygens in 1659

Johann Bernoulli, 1696 (first published ‘integral’)

Led to the calculus of variations

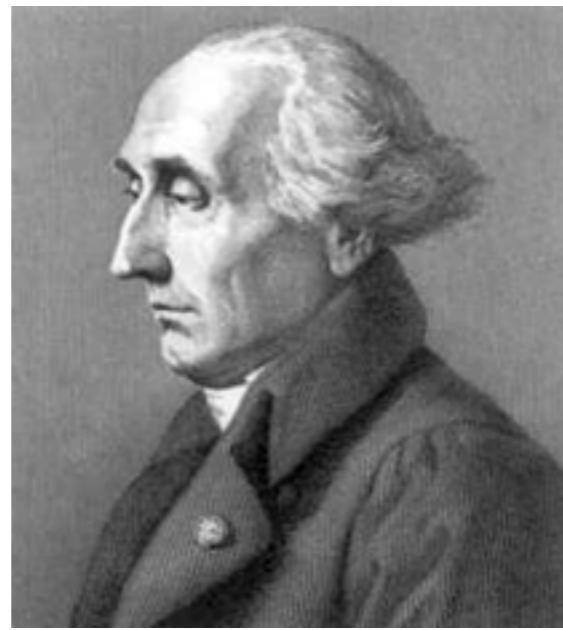
The Lagrangian

$$L = T - V$$

$$S = \int L \, dt$$

The action

Hamilton's principle



Joseph Louis Lagrange (1736–1813)
William Rowan Hamilton (1805–1865)

$$H = T + V$$

$$\dot{p}_i = -\frac{\partial H}{\partial q_i}$$

$$\dot{q}_i = +\frac{\partial H}{\partial p_i}$$

Describes motion in a 2n-dimensional space

But this tells us that only $2 \times 3 = 6$ numbers completely characterise motion in 3-d space.

Also, this is true for more than just particle mechanics.

Leads to Newton's equations, Maxwell's, Schrödinger's; electrical and acoustic systems.

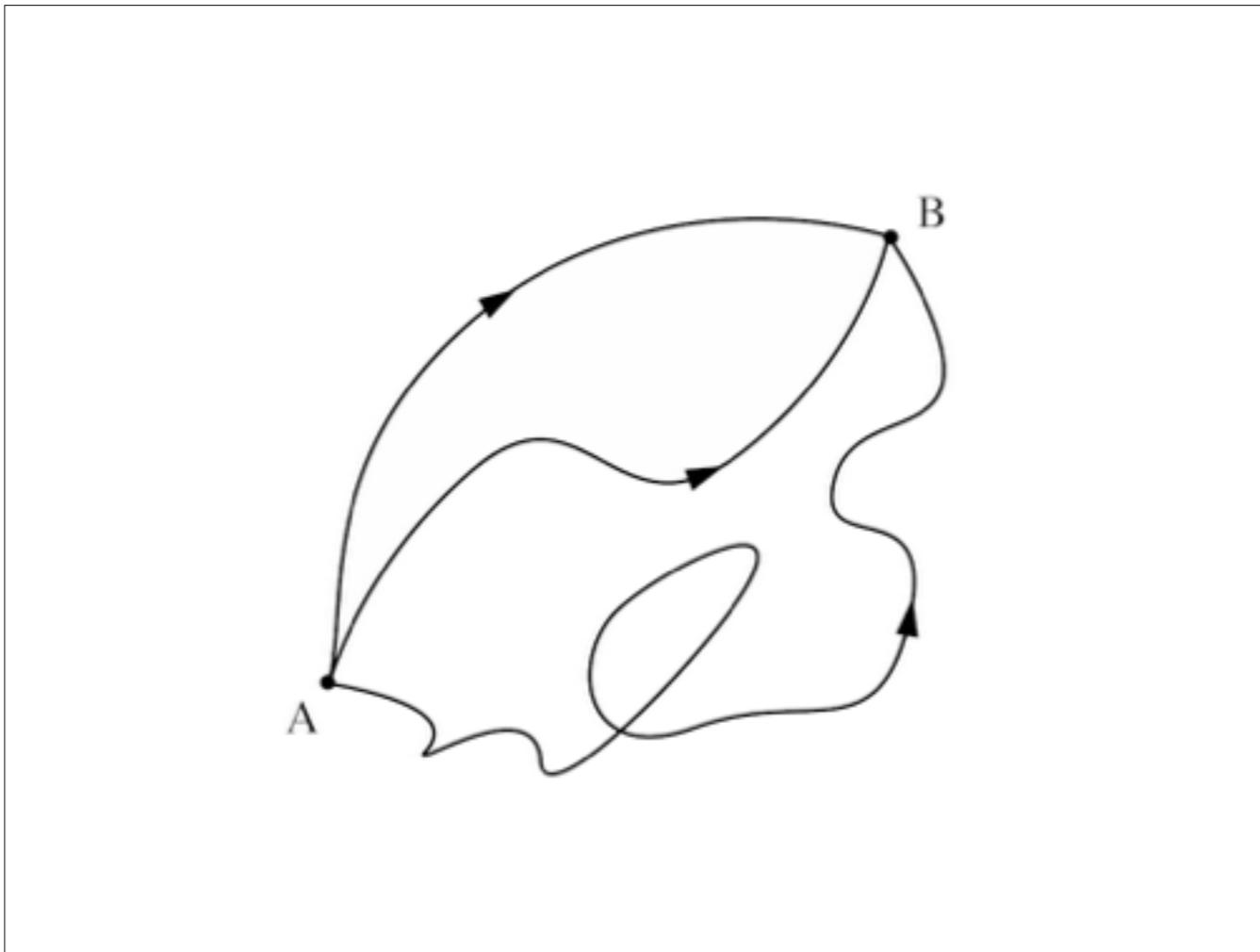
Not always useful for actual calculations, though

$$i\hbar \frac{\partial}{\partial t} \psi = \hat{H}\psi$$

$$\hat{H}\psi = E\psi$$

Schrödinger equation (one form of it)

Similarly, it's the classical theory of fields (based on continuous Lagrangians) which is used as the basis for quantum field theory



Feynman path-integral formulation
(building on work by Dirac)
Each path is weighted by $\exp(iS/\hbar)$

$$\dot{p}_i = -\frac{\partial H}{\partial q_i}$$

$$\dot{q}_i = +\frac{\partial H}{\partial p_i}$$

If H indep of q , then p is conserved



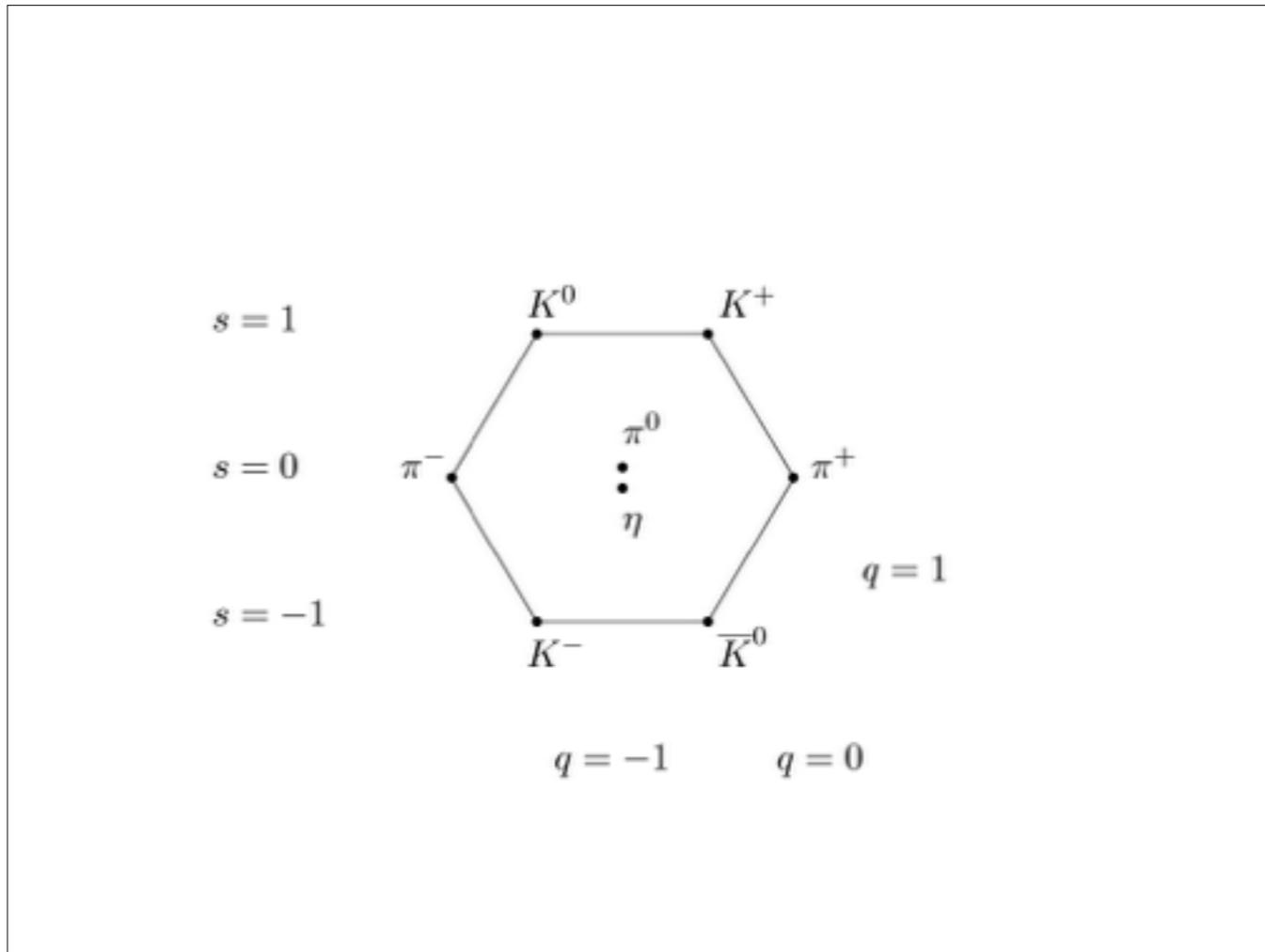
Emmy Noether (1882–1935)

Noether's theorem

For each symmetry of the Lagrangian,
there is a conserved quantity

$$\begin{aligned}
L = & -\frac{1}{2}\partial_\nu g_\mu^a \partial_\nu g_\mu^a - g_s f^{abc} \partial_\mu g_\mu^a g_\mu^b g_\nu^c - \frac{1}{4}g_s^2 f^{abc} f^{ade} g_\mu^b g_\nu^c g_\mu^d g_\nu^e + \frac{1}{2}ig_s^2 (\bar{q}_j^\sigma \gamma^\mu q_j^\sigma) g_\mu^a + \bar{G}^a \partial^2 G^a + \\
& g_s f^{abc} \partial_\mu \bar{G}^a g_\mu^c - \partial_\nu W_\mu^+ \partial_\nu W_\mu^- - M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2c_w^2} M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \\
& \frac{1}{2}\partial_\mu H \partial_\mu H - \frac{1}{2}m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2c_w^2} M \phi^0 \phi^0 - \beta_h [\frac{2M^2}{g^2} + \frac{2M}{g} H + \\
& \frac{1}{2}(H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-)] + \frac{2M^4}{g^2} \alpha_h - ig c_w [\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - Z_\nu^0 (W_\mu^+ \partial_\nu W_\mu^- - \\
& W_\mu^- \partial_\nu W_\mu^+) + Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+)] - ig s_w [\partial_\nu A_\mu (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - \\
& W_\mu^- \partial_\nu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+)] - \frac{1}{2}g^2 W_\mu^+ W_\nu^+ W_\mu^- + \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\mu^- + \\
& g^2 c_w^2 (Z_\mu^0 W_\mu^+ Z_\nu^0 W_\nu^- - Z_\mu^0 Z_\nu^0 W_\nu^+ W_\mu^-) + g^2 s_w^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\nu W_\mu^+ W_\nu^-) + \\
& g^2 s_w c_w [A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - 2A_\mu Z_\mu^0 W_\nu^+ W_\nu^-] - g\alpha [H^3 + H\phi^0 \phi^0 + 2H\phi^+ \phi^-] - \\
& \frac{1}{8}g^2 \alpha_h [H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2] - gM W_\mu^+ W_\mu^- H - \\
& \frac{1}{2}g \frac{M}{c_w^2} Z_\mu^0 Z_\mu^0 H - \frac{1}{2}ig [W_\mu^+ (\phi^0 \partial_\mu \phi^- - \phi^- \partial_\mu \phi^0) - W_\mu^- (\phi^0 \partial_\mu \phi^+ - \phi^+ \partial_\mu \phi^0)] + \frac{1}{2}g [W_\mu^+ (H \partial_\mu \phi^- - \\
& \phi^- \partial_\mu H) - W_\mu^- (H \partial_\mu \phi^+ - \phi^+ \partial_\mu H)] + \frac{1}{2}g \frac{1}{c_w} (Z_\mu^0 (H \partial_\mu \phi^0 - \phi^0 \partial_\mu H) - ig \frac{s_w^2}{c_w} M Z_\mu^0 (W_\mu^+ \phi^- - \\
& W_\mu^- \phi^+) + ig s_w M A_\mu (W_\mu^+ \phi^- - W_\mu^- \phi^+) - ig \frac{1-2c_w^2}{2c_w} Z_\mu^0 (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) + ig s_w A_\mu (\phi^+ \partial_\mu \phi^- - \\
& \phi^- \partial_\mu \phi^+) - \frac{1}{4}g^2 W_\mu^+ W_\mu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \frac{1}{4}g^2 \frac{1}{c_w} Z_\mu^0 Z_\mu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^-] - \\
& \frac{1}{2}g^2 \frac{s_w^2}{c_w} Z_\mu^0 \phi^0 (W_\mu^+ \phi^- + W_\mu^- \phi^+) - \frac{1}{2}ig^2 \frac{s_w^2}{c_w} Z_\mu^0 H (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \frac{1}{2}g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + \\
& W_\mu^- \phi^+) + \frac{1}{2}ig^2 s_w A_\mu H (W_\mu^+ \phi^- - W_\mu^- \phi^+) - g^2 \frac{s_w^2}{c_w} (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+ \phi^- - g^1 s_w^2 A_\mu A_\mu \phi^+ \phi^- - \\
& \bar{e}^\lambda (\gamma \partial + m_e^\lambda) e^\lambda - \bar{\nu}^\lambda \gamma \partial \nu^\lambda - \bar{u}_j^\lambda (\gamma \partial + m_u^\lambda) u_j^\lambda - \bar{d}_j^\lambda (\gamma \partial + m_d^\lambda) d_j^\lambda + ig s_w A_\mu [-(\bar{e}^\lambda \gamma^\mu e^\lambda) + \\
& \frac{2}{3}(\bar{u}_j^\lambda \gamma^\mu u_j^\lambda) - \frac{1}{3}(\bar{d}_j^\lambda \gamma^\mu d_j^\lambda)] + \frac{ig}{4c_w} Z_\mu^0 [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{e}^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) + (\bar{u}_j^\lambda \gamma^\mu (\frac{4}{3}s_w^2 - \\
& 1 - \gamma^5) u_j^\lambda) + (\bar{d}_j^\lambda \gamma^\mu (1 - \frac{8}{3}s_w^2 - \gamma^5) d_j^\lambda)] + \frac{ig}{2\sqrt{2}} W_\mu^+ [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) e^\lambda) + (\bar{u}_j^\lambda \gamma^\mu (1 + \gamma^5) C_{\lambda\kappa} d_j^\kappa)] + \\
& \frac{ig}{2\sqrt{2}} W_\mu^- [(\bar{e}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{d}_j^\kappa C_{\lambda\kappa}^\dagger \gamma^\mu (1 + \gamma^5) u_j^\lambda)] + \frac{ig}{2\sqrt{2}} \frac{m_\lambda^2}{M} [-\phi^+ (\bar{\nu}^\lambda (1 - \gamma^5) e^\lambda) + \phi^- (\bar{e}^\lambda (1 + \\
& \gamma^5) \nu^\lambda)] - \frac{g}{2} \frac{m_\lambda^2}{M} [H (\bar{e}^\lambda \gamma^\mu) + i\phi^0 (\bar{e}^\lambda \gamma^5 e^\lambda)] + \frac{ig}{2M\sqrt{2}} \phi^+ [-m_d^\kappa (\bar{u}_j^\lambda C_{\lambda\kappa} (1 - \gamma^5) d_j^\kappa) + m_u^\lambda (\bar{u}_j^\lambda C_{\lambda\kappa} (1 + \\
& \gamma^5) d_j^\kappa)] + \frac{ig}{2M\sqrt{2}} \phi^- [m_d^\kappa (\bar{d}_j^\lambda C_{\lambda\kappa}^\dagger (1 + \gamma^5) u_j^\kappa) - m_u^\kappa (\bar{d}_j^\lambda C_{\lambda\kappa}^\dagger (1 - \gamma^5) u_j^\kappa)] - \frac{g}{2} \frac{m_\lambda^2}{M} H (\bar{u}_j^\lambda u_j^\lambda) - \\
& \frac{g}{2} \frac{m_\lambda^2}{M} H (\bar{d}_j^\lambda d_j^\lambda) + \frac{ig}{2} \frac{m_\lambda^2}{M} \phi^0 (\bar{u}_j^\lambda \gamma^5 u_j^\lambda) - \frac{ig}{2} \frac{m_\lambda^2}{M} \phi^0 (\bar{d}_j^\lambda \gamma^5 d_j^\lambda) + \bar{X}^+ (\partial^2 - M^2) X^+ + \bar{X}^- (\partial^2 - \\
& M^2) X^- + \bar{X}^0 (\partial^2 - \frac{M^2}{c_w^2}) X^0 + \bar{Y} \partial^2 Y + ig c_w W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \partial_\mu \bar{X}^+ X^0) + ig s_w W_\mu^+ (\partial_\mu \bar{Y} X^- - \\
& \partial_\mu \bar{X}^+ Y) + ig c_w W_\mu^- (\partial_\mu \bar{X}^- X^0 - \partial_\mu \bar{X}^0 X^+) + ig s_w W_\mu^- (\partial_\mu \bar{X}^- Y - \partial_\mu \bar{Y} X^+) + \\
& ig c_w Z_\mu^0 (\partial_\mu \bar{X}^+ X^+ - \partial_\mu \bar{X}^- X^-) + ig s_w A_\mu (\partial_\mu \bar{X}^+ X^+ - \partial_\mu \bar{X}^- X^-) - \frac{1}{2}gM [\bar{X}^+ X^+ H + \\
& \bar{X}^- X^- H + \frac{1}{c_w^2} \bar{X}^0 X^0 H] + \frac{1-2c_w^2}{2c_w} ig M [\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-] + \frac{1}{2c_w} ig M [\bar{X}^0 X^- \phi^+ - \\
& \bar{X}^0 X^+ \phi^-] + ig M s_w [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \frac{1}{2}ig M [\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0]
\end{aligned}$$

The standard model of particle physics



Eightfold way – meson octet.
Murray Gell-Mann

