

Black body source $\Rightarrow$ Wien's displacement law holds

$$\text{i.e. } \lambda_{\text{max}} T = 0.29 \text{ cm}$$

$$E = \frac{1.24 \times 10^{-6} \text{ eV}}{\lambda(\text{m})} = \frac{1.24 \times 10^{-7} \text{ keV}}{\lambda(\text{cm})}$$

$$E_{\text{max}} = 2.5 \text{ keV} \Rightarrow \lambda_{\text{max}} = 5 \times 10^{-8} \text{ cm}$$

$$\therefore T = \frac{0.29}{5 \times 10^{-8}} \text{ K} = 5.8 \times 10^6 \text{ K}$$

$$L = 10^{31} \text{ W} = 4\pi R^2 \sigma T^4 \Rightarrow R = \sqrt{\frac{L}{4\pi\sigma T^4}}$$

$$= \left(\frac{10^{31}}{4\pi \times 5.67 \times 10^{-8} \times (5.8 \times 10^6)^4} \right)$$

$$= 1.1 \times 10^5 \text{ m}$$

If luminosity arises from mass transfer, equate energy released with rest mass energy of transferred mass. (Not quite so efficient, but approx correct)

$$\text{In one year energy released, } E = L \times 3600 \times 24 \times 365.25$$

$$= 3.16 \times 10^{38} \text{ J}$$

$$\text{Equivalent rest mass} = \frac{E}{c^2} = 3.5 \times 10^{21} \text{ kg}$$

$$= \frac{3.5 \times 10^{21}}{2 \times 10^{30}} M_{\odot}$$

$$\text{Mass loss rate, } \dot{M} = 1.3 \times 10^{-9} M_{\odot} / \text{year}$$

$$2) \quad B_\lambda d\lambda = B_\nu d\nu$$

$$\Leftrightarrow B_\lambda = B_\nu(\lambda) \left| \frac{d\nu}{d\lambda} \right|$$

$$c = \nu \lambda$$

$$\Rightarrow \nu = \frac{c}{\lambda}$$

$$\Leftrightarrow d\nu = -\frac{c}{\lambda^2} d\lambda$$

$$= \frac{2h \left(\frac{c}{\lambda} \right)^3 \cdot \frac{1}{c^2} \cdot \frac{c}{\lambda^2}}{e^{\frac{hc}{\lambda kT}} - 1}$$

$$= \frac{2hc^2/\lambda^5}{e^{\frac{hc}{\lambda kT}} - 1}$$

$$B = \int_0^\infty B_\lambda d\lambda = \int_0^\infty \frac{2hc^2/\lambda^5}{e^{\frac{hc}{\lambda kT}} - 1} d\lambda$$

$$\text{Put } s = \frac{hc}{\lambda kT} \quad \lambda = \frac{hc}{k s T} \quad = \int_0^\infty \frac{2hc^2 (k s T)^5 \left(\frac{1}{hc} \right)^5 \frac{hc kT ds}{(kT)^2 s^2}}{e^s - 1}$$

$$\Rightarrow ds = -\frac{hc}{\lambda^2 kT} d\lambda$$

$$d\lambda = -\frac{\lambda^2 kT}{hc}$$

$$d\lambda = -\frac{hc kT}{(k s T)^2} ds$$

$$= \frac{2hc^2 k^5 T^5 hc kT}{(hc)^5 k^2 T^2} \int_0^\infty \frac{s^3 ds}{e^s - 1}$$

$$= C \int_0^\infty \frac{s^3 ds}{e^s - 1} T^4 \propto T^4$$

$$(3) \quad \tau = \frac{U}{L} = \frac{\frac{3R}{4\pi\sigma m_p} \frac{M}{R^2 T^4}}{L}$$

$$= \frac{3k M_{\odot} \left(\frac{M}{M_{\odot}} \right)}{4\pi\sigma m_p R_{\odot}^2 \left(\frac{R}{R_{\odot}} \right)^2 10^{21} \left(\frac{T}{10^7} \right)^3}$$

$$= \frac{3 \times 1.38 \times 10^{-23} \times 2 \times 10^{30} \left(\frac{M}{M_{\odot}} \right)}{4\pi \times 5.67 \times 10^{-8} \times 1.67 \times 10^{-27} \times (7 \times 10^8)^2 \times 10^{21} \left(\frac{T}{10^7} \right)^3}$$

$$= \frac{(142 \text{ seconds}) \left(\frac{M}{M_{\odot}} \right)}{\left(\frac{R}{R_{\odot}} \right)^2 \left(\frac{T}{10^7 \text{ K}} \right)^3}$$

$$(4) \quad (a) \quad F_G = \frac{GMm_p}{R^2} \quad (\text{ignoring } m_e)$$

$$(b) \quad \text{For photons, momentum} = \frac{E}{c}$$

$\Rightarrow$ rate of change of momentum

$$= \text{momentum radiated per second} = \frac{L}{c}$$

Probability of a photon being scattered by a particular electron

$$p = \frac{Q_T}{4\pi R^2}$$

$$\Rightarrow \text{radiation pressure transfers } F_R = \frac{L Q_T}{4\pi R^2 c}$$

$$F_R = F_G \Rightarrow \frac{L Q_T}{4\pi R^2 c} = \frac{GM m_p}{R^2}$$

$$\Leftrightarrow L = \frac{4\pi G M m_p c}{Q_T}$$

(5) By energy conservation we need $E > \varepsilon$

$$\begin{aligned} \frac{dJ}{d\varepsilon} &= \int_V \frac{dj}{d\varepsilon} dV \\ &= \int_V n_p \int_{\varepsilon}^{\infty} F(E) \frac{dQ_B(\varepsilon, E)}{d\varepsilon} dE dV \end{aligned}$$

Homogeneous plasma, so we can take out $\int n_p dV = n_p V$

$$\text{Also } \frac{dQ_B}{d\varepsilon} \approx \frac{Q_0 mc^2}{\varepsilon E}$$

$$\Rightarrow \frac{dJ}{d\varepsilon} = \frac{n_p V Q_0 mc^2}{\varepsilon} \int_{\varepsilon}^{\infty} \frac{F(E)}{E} dE \quad \text{as required}$$

Re-arranging ,

$$\varepsilon \frac{dJ}{d\varepsilon} = n_p V Q_0 mc^2 \int_{\varepsilon}^{\infty} \frac{F(E)}{E} dE$$

(5) contd

Differentiating both sides w.r.t. ε

$$\frac{d}{d\varepsilon} \left(\varepsilon \frac{dJ}{d\varepsilon} \right)_{\varepsilon=E} = - \frac{F(E)}{E} \cdot \frac{1}{n_p V Q_0 mc^2} \quad (\text{using Leibniz' rule})$$

$$\text{i.e.} \quad \left[\frac{dJ}{d\varepsilon} + \varepsilon \frac{d^2 J}{d\varepsilon^2} \right]_{\varepsilon=E} = \frac{-F(E)}{E n_p V Q_0 mc^2}$$

$$\Rightarrow F(E) = \frac{E}{n_p V Q_0 mc^2} \left[- \frac{dJ}{d\varepsilon} - \varepsilon \frac{d^2 J}{d\varepsilon^2} \right]_{\varepsilon=E}$$