

Orbit determination for extra-solar planetary systems.

Lecture I

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Outline

- 1 Extra-solar planets
- 2 Observations
- 3 Modeling observations
- 4 Fitting orbits.

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Early and false discoveries

- The Barnard star planetary system.
- Peter van de Kamp observations 1938–1982 (about 10000 plates).
- Final estimate: two planets with masses: $M_1 = 0.7M_J$ and $M_2 = 0.5M_J$ and periods $P_1 = 12\text{y}$ and $P_2 = 20\text{y}$, respectively; (Vistas in Astronomy, 26:141–157, 1982).

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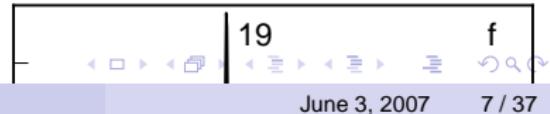
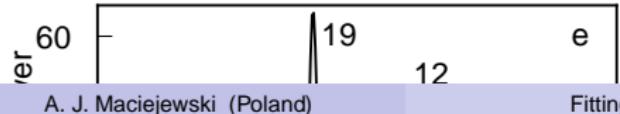
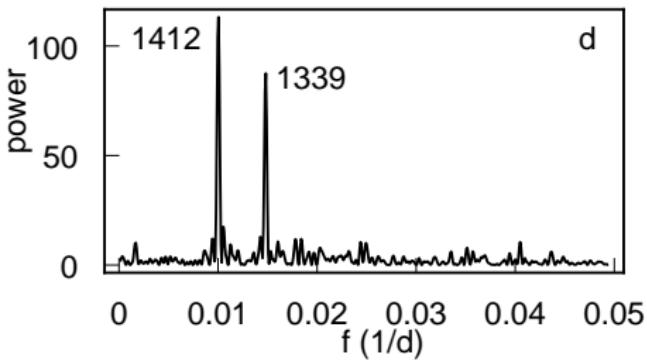
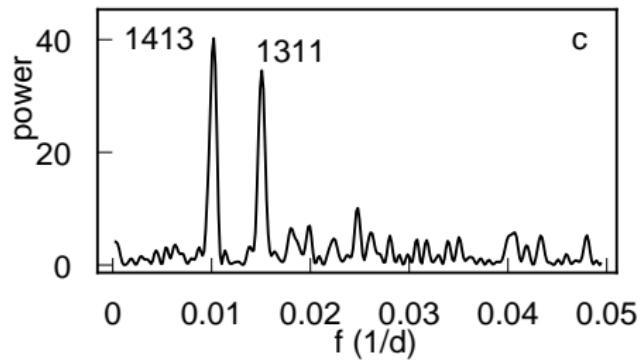
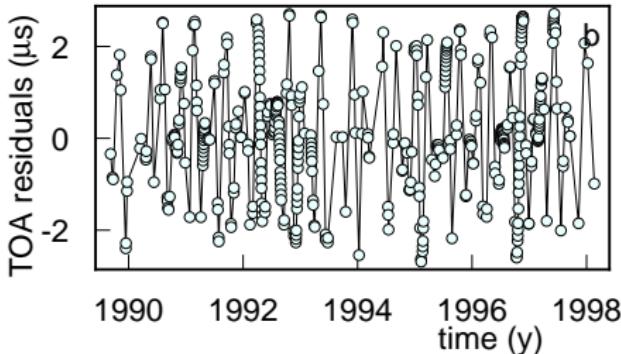
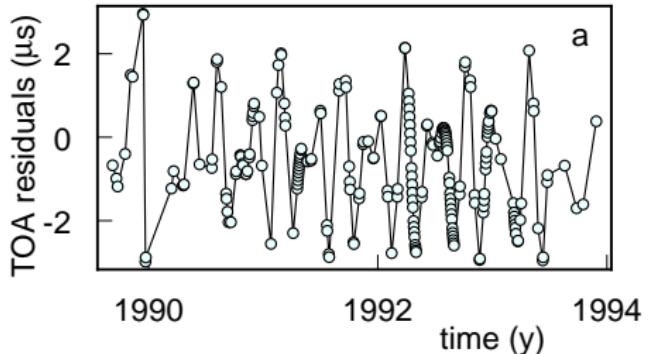
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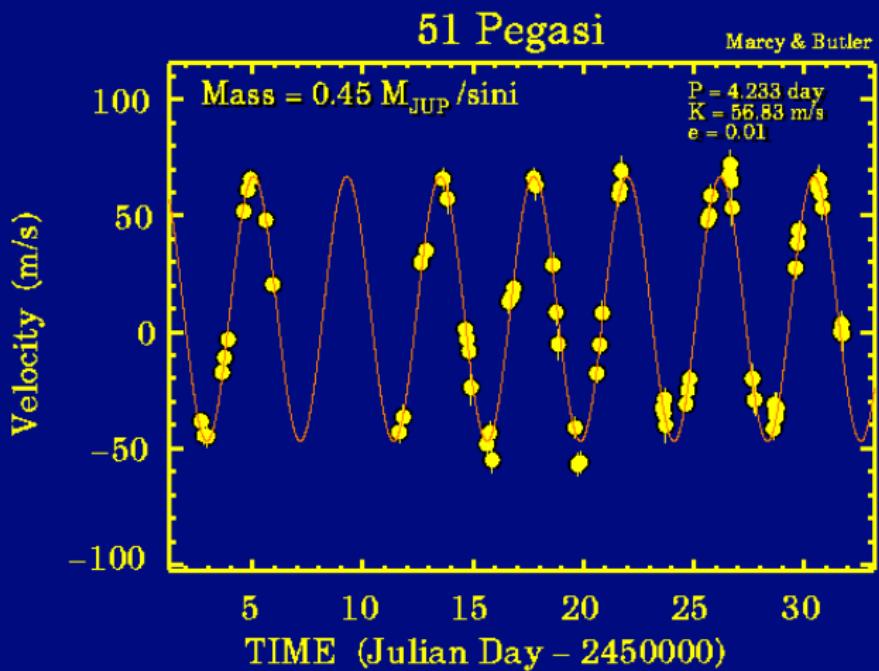
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51Peg planet

Mayor M. and Queloz D., 1995, *A Jupiter-mass companion to a solar-type star*, Nature, **378**, 355.

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Actually (24 May, 2007)

- Radial velocities: 224 planets, 23 multiple planetary systems.
- Transit: 20 planets.
- Microlensing: 4 planets.
- Imaging: 4 planets.
- Timing: 4 planets, 1 multiple planetary system.
- All: 236 planets, 24 multiple planetary systems.

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- Kepler.
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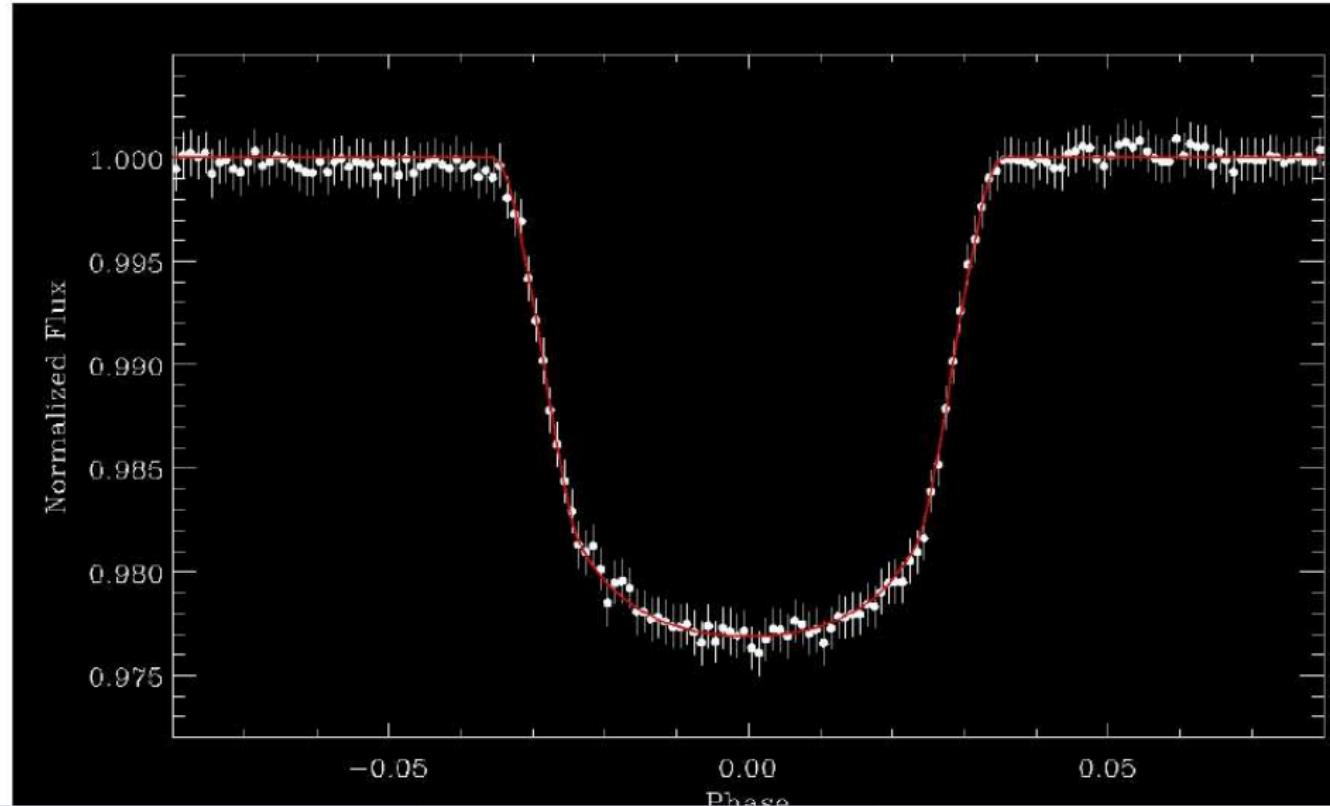
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Future

Tomorrow—it is today but tomorrow. —St. Mrożek



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Detection methods

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- Pulsars timing.
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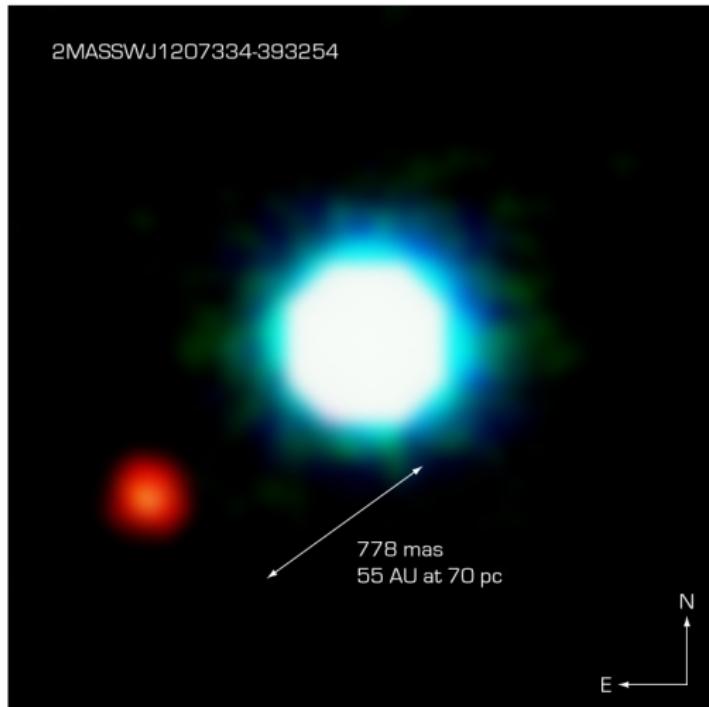
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Imaging



NACO Image of the Brown Dwarf Object 2M1207 and GPCC

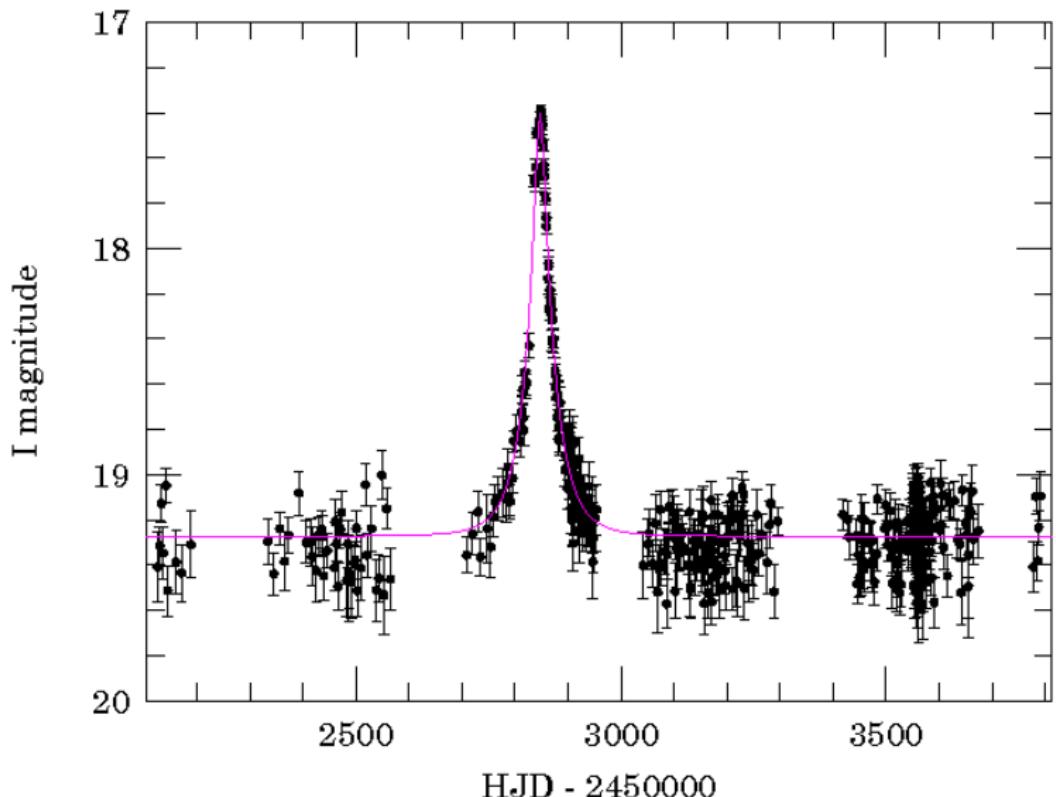
ESO PR Photo 26a/04 (10 September 2004)

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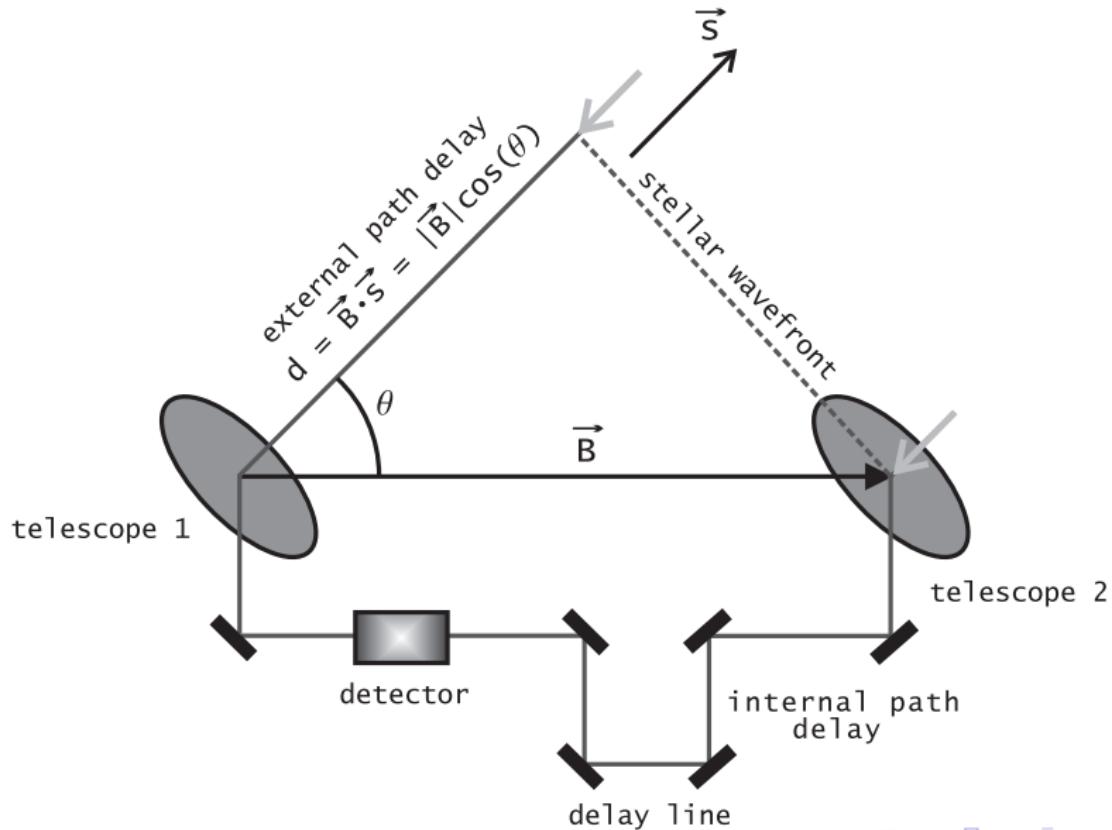


Microlensing

OGLE-2003-BLG-235



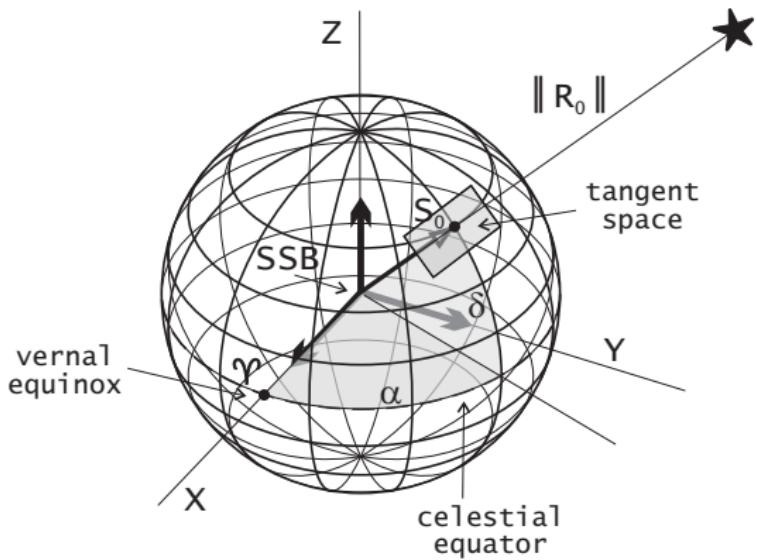
Space interferometry



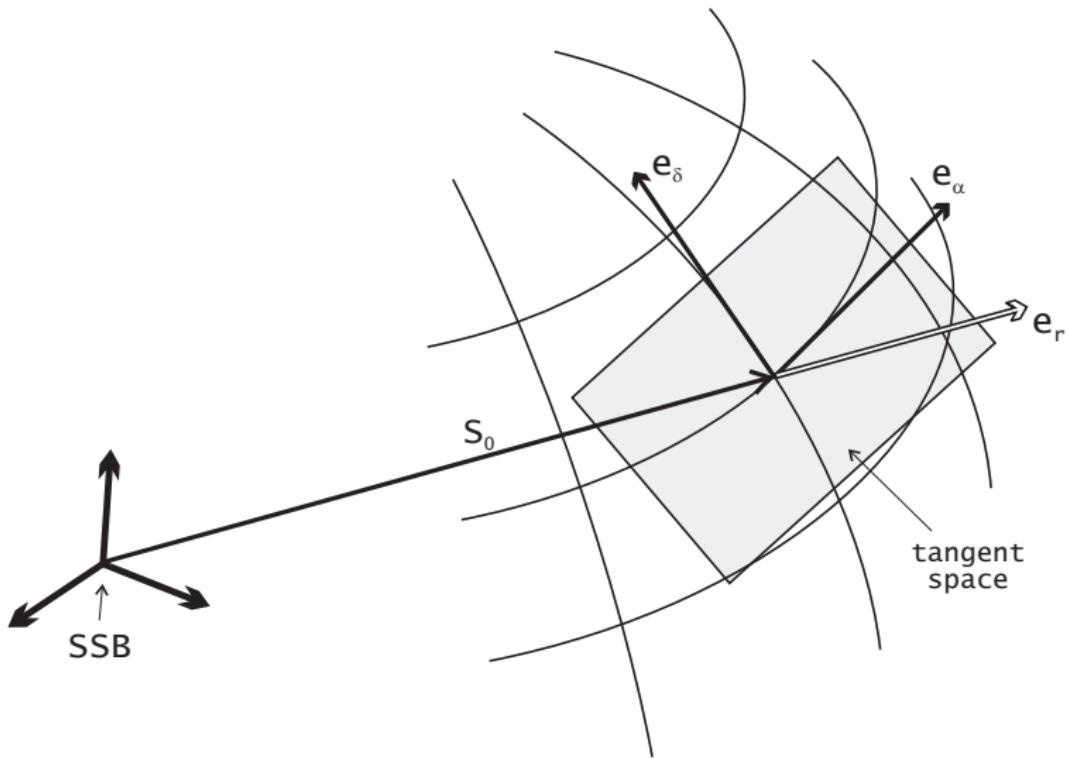
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Geometry



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Basic idea

If $\mathbf{V}^*(t)$ is the barycentric velocity of the star and \mathbf{e}_r the unit vector from observer to the star planetary system then the observed radial velocity is given by

$$v_r(t) := \mathbf{e}_r \cdot \mathbf{V}^*(t).$$

Thus, if $\mathbf{e}_3 = \mathbf{e}_r$, then $v_r(t) = V_3^*$, and $v_r(t)$ does not depend how we choose \mathbf{e}_1 and \mathbf{e}_2 .

One planet.

The mass center

$$m\mathbf{R} + m_\star\mathbf{R}^\star = \mathbf{0}, \quad m\mathbf{V} + m_\star\mathbf{V}^\star = \mathbf{0}$$

$$\mathbf{r} = \mathbf{R} - \mathbf{R}^\star, \quad \mathbf{v} = \mathbf{V} - \mathbf{V}^\star,$$

$$\mathbf{R}^\star = -\frac{m}{m_\star}\mathbf{R} = -\frac{m}{m_\star + m}\mathbf{r}$$

$$\mathbf{V}^\star = -\frac{m}{m_\star}\mathbf{V} = -\frac{m}{m_\star + m}\mathbf{v}$$

$$\ddot{\mathbf{R}} = -\frac{\mu_r}{R^3}\mathbf{R}, \quad \mu_b = \frac{Gm_\star^3}{(m_\star + m)^2}$$

$$\ddot{\mathbf{r}} = -\frac{\mu_r}{r^3}\mathbf{r}, \quad \mu_r = G(m_\star + m),$$

Kepler orbit

$$\mathbf{R}(t) = \frac{a(1 - e^2)}{1 + e \cos(\nu)} [\cos(\nu) \mathbf{P} + \sin(\nu) \mathbf{Q}],$$

$$\mathbf{V}(t) = \frac{an}{\sqrt{1 - e^2}} [-\sin(\nu) \mathbf{P} + (e + \cos(\nu)) \mathbf{Q}],$$

$$\mathbf{P} = \mathbf{I} \cos(\omega) + \mathbf{m} \sin(\omega), \quad \mathbf{Q} = -\mathbf{I} \sin(\omega) + \mathbf{m} \cos(\omega),$$

$$\mathbf{I} = \begin{bmatrix} \cos \Omega \\ \sin \Omega \\ 0 \end{bmatrix}, \quad \mathbf{m} = \begin{bmatrix} -\cos i \sin \Omega \\ \cos i \cos \Omega \\ \sin i \end{bmatrix},$$

$$\tan \frac{\nu}{2} = \sqrt{\frac{1+e}{1-e}} \tan \frac{E}{2}, \quad E - e \sin E = n(t - T_P), \quad n^2 a^3 = \mu.$$

Final formula

$$v_r(t) = -K [\cos(\nu(t) + \omega) + e \cos \omega] + v_0,$$

$$K = \frac{\sigma a n \sin i}{\sqrt{1 - e^2}}, \quad \sigma = \frac{m}{m_*}.$$

Parameters

$$\mathbf{p} = (K, n, e, \omega, T_p, v_0)$$

More than one planet.

- Keplerian models.
- Gravitational N -body problem.

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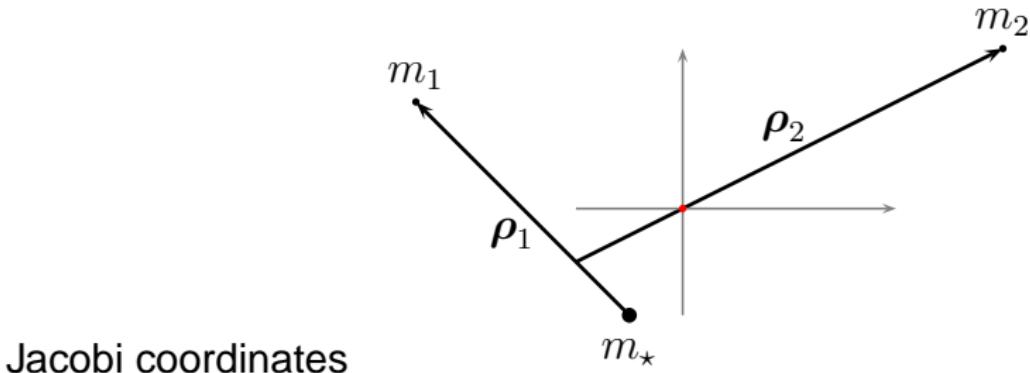
More than one planet. Keplerian models.

Just add

$$v_r(i) = \sum_{k=1}^n v_r^{(k)} + v_0,$$

$$v_r^{(k)} = -K_k [\cos(\nu_k(t) + \omega_k) + e_k \cos \omega_k], \quad k = 1, \dots, n.$$

Yet another Keplerian model.



$$v_r = -K_1 f_1(t) - K_2 f_2(t),$$

$$K_k = \frac{\sigma_k a_k n_k \sin i_k}{\sqrt{1 - e_k^2}}, \quad f_k(t) = \cos(\nu_k(t) + \omega_k) + e_k \cos \omega_k$$

$$\sigma_1 = \frac{m_1}{m_* + m_1}, \quad \sigma_2 = \frac{m_2}{m_* + m_1 + m_2}$$

$$\mu_1 = G(m_* + m_1), \quad \mu_1 = G(m_* + m_1 + m_2)$$

Gravitational N -body model.

Barycentric reference frame

$$\left. \begin{aligned} \dot{\mathbf{R}}_i &= \mathbf{v}_i, \\ m_i \dot{\mathbf{v}}_i &= - \sum_{j=0}^n' \frac{Gm_i m_j}{R_{ij}^3} (\mathbf{R}_i - \mathbf{R}_j), \quad \mathbf{v}_i := \dot{\mathbf{R}}_i \end{aligned} \right\}, \quad i = 0, \dots, n,$$
$$\sum_{i=0}^n m_i \mathbf{R}_i = \mathbf{0}, \quad \sum_{i=0}^n m_i \dot{\mathbf{v}}_i = \mathbf{0}$$
$$v_r(t) = - \sum_{i=1}^n \frac{m_i}{m_0} v_{i,3}.$$

Gravitational N -body model. Parametrisation.

- ➊ Osculating elements: $(m_k, \mathbf{a}_k, \mathbf{e}_k, i_k, \omega_k, \Omega_k, T_{\text{p},k})$, for $k = 1, \dots, n$,
and $\Omega_1 = 0!$
- ➋ Initial condition: $(m_k, \mathbf{R}_k(t_0), \mathbf{V}_k(t_0))$, for $k = 1, \dots, n$, and fix e.g.
 $\mathbf{R}_{1,1} = 0!$

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An opinion.

An astrophysicist can fit a sinusoid to two observations and moreover he claims that this fit is very good.

— Ryszard Szczerba, astrophysicist

Classical χ^2 method.

One-dimensional observations $x(t)$.

- Observations: x_1, \dots, x_N , $x_i = x(t_i)$
- Errors of individual observations: $\sigma_1, \dots, \sigma_N$.
- Model: $\hat{x} = \hat{x}(t, \mathbf{p})$, $\mathbf{p} = (p_1, \dots, p_k) \in \mathbb{R}^k$.
-

$$\chi^2(\mathbf{p}) = \frac{1}{N} \sum_{i=1}^N \left(\frac{x_i - \hat{x}(t_i, \mathbf{p})}{\sigma_i} \right)^2,$$

- degrees of freedom: $\nu = N - k - 1$,

$$\chi_\nu^2(\mathbf{p}) = \frac{1}{\nu} \sum_{i=1}^N \left(\frac{x_i - \hat{x}(t_i, \mathbf{p})}{\sigma_i} \right)^2.$$

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Classical χ^2 method.

Problem: find $\mathbf{p}_0 \in \mathbb{R}^k$ such that

$$\chi^2(\mathbf{p}_0) = \min_{\mathbf{p} \in \mathbb{R}^k} \chi^2(\mathbf{p}).$$

How to find the global minimum?

Find all local minima and choose the smallest one!

Facile dictu difficile factu!

Local minimum

$$\mathbf{p}_0 = \min_{\mathbf{p} \in U} F(\mathbf{p}), \quad U \in \mathbb{R}^k;$$

Condition: $\|\mathbf{p} - \mathbf{p}_0\|$ is small. Then

$$F(\mathbf{p}) = F(\mathbf{p}_0) + \nabla F(\mathbf{p}_0) \mathbf{x} + \frac{1}{2} \mathbf{x}^T \nabla^2 F(\mathbf{p}_0) \mathbf{x} + \dots,$$

$$\mathbf{x} = \mathbf{p} - \mathbf{p}_0, \quad A := \left[\frac{\partial^2 F}{\partial p_i \partial p_j} (\mathbf{p}_0) \right].$$

Quadratic approximation

- ➊ \mathbf{p}_1 close to \mathbf{p}_0 ;
- ➋ $\mathbf{p}_2 = \mathbf{p}_1 - A^{-1}\mathbf{y}$; $\mathbf{y} = \nabla F(\mathbf{p}_1)$, $A = \nabla^2 F(\mathbf{p}_1)$.
- ➌ $\|\mathbf{p}_1 - \mathbf{p}_2\| < \epsilon$?
- ➍ $\mathbf{p}_1 = \mathbf{p}_2$ and go 1.

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Problems

- Expensive calculations of $\nabla F(\mathbf{p})$ and/or $\nabla^2 F(\mathbf{p})$. Solution: choose the simplex method.
- Initial approximation. Monte Carlo, periodograms, or pseudo-global genetic algorithm.
- Natural constraints.