

Using Albedo to Determine Changing Anisotropy in a Flare

Ewan Dickson

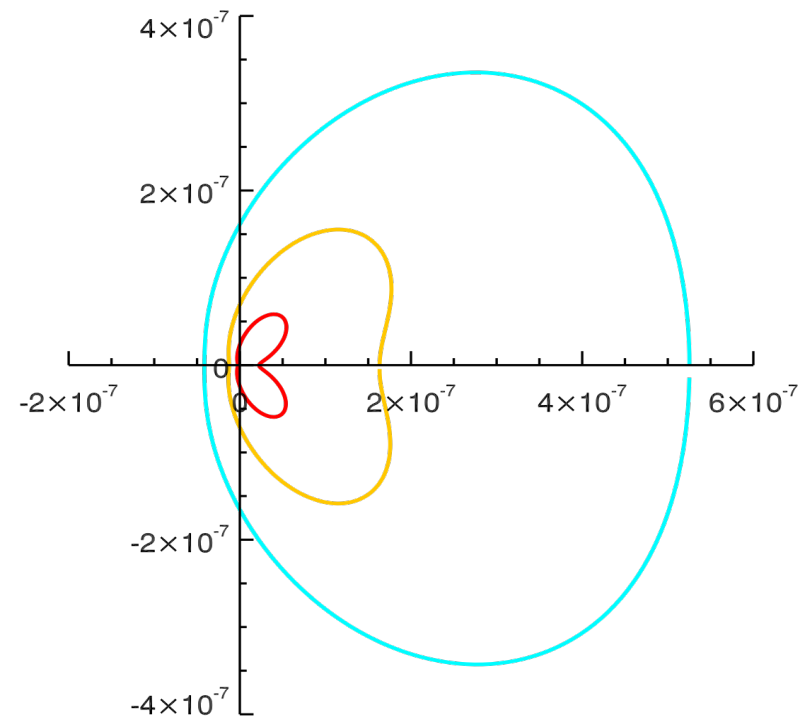
Eduard Kontar





- Observed X-ray spectrum must be related to emitting electron spectrum
- In general electron distribution will vary with angle.
- Many models assume strong downward beaming.
- However in when determining the electron distribution from observations an isotropic emission is assumed.

$$I(\epsilon) = \frac{\bar{n}V}{4\pi R^2} \int_0^\pi \int_\epsilon^\infty \bar{F}(E, \theta) Q(\epsilon, E, \theta) dE d\theta$$

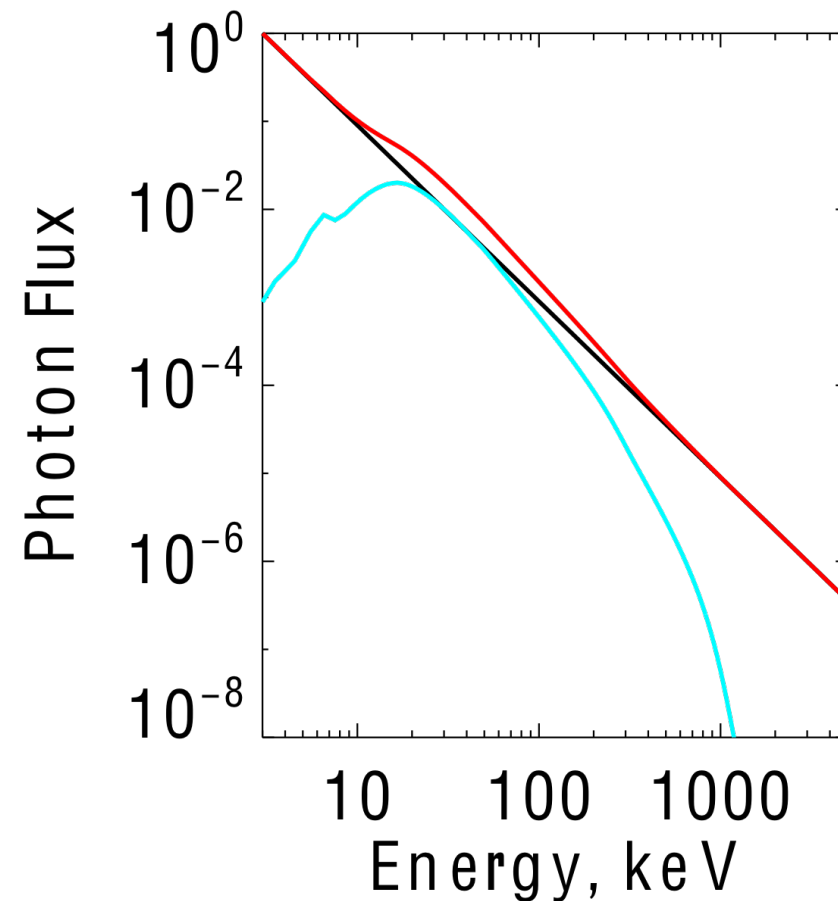


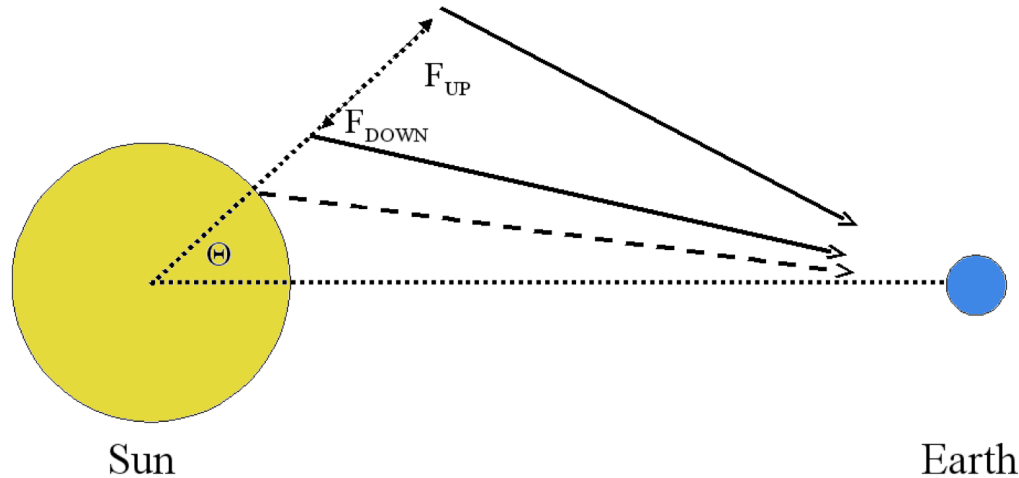


- One technique is to look at the centre to limb variation.
- Disadvantage is that variation can only be seen as an average over a large number of flares.
- More direct approach is the stereoscopic method. The disadvantage of this is the difficulty in cross calibrating, leading to large errors.
- Previous research on this problem suggests that the electron angular distribution is fairly isotropic with some studies showing directivity at higher energies



- Downward travelling X-rays can Compton backscatter low in the solar atmosphere and be observed at Earth.
- Albedo distorts primary X-ray spectrum.
- This effect will vary depending on the fraction of downward to upward going electrons.
- Direct estimate of downward going electrons





- Divide the electron flux into two components one going up away from the solar centre and one going down towards it.
- Angular dependant cross-section is averaged over two hemispheres
- Can be used to distinguish between highly beamed and isotropic cases.

$$\bar{Q}(\epsilon, E, \theta_0, \alpha) = \frac{1}{2\pi(1 - \cos(\alpha))} \int_{\phi=0}^{2\pi} \int_{\beta=0}^{2\pi} Q(\epsilon, E, \theta') \sin \beta d\beta d\phi$$

$$I = \begin{pmatrix} Q^F + A Q^B & Q^B + A Q^F \end{pmatrix} \begin{pmatrix} F_U \\ F_D \end{pmatrix}$$

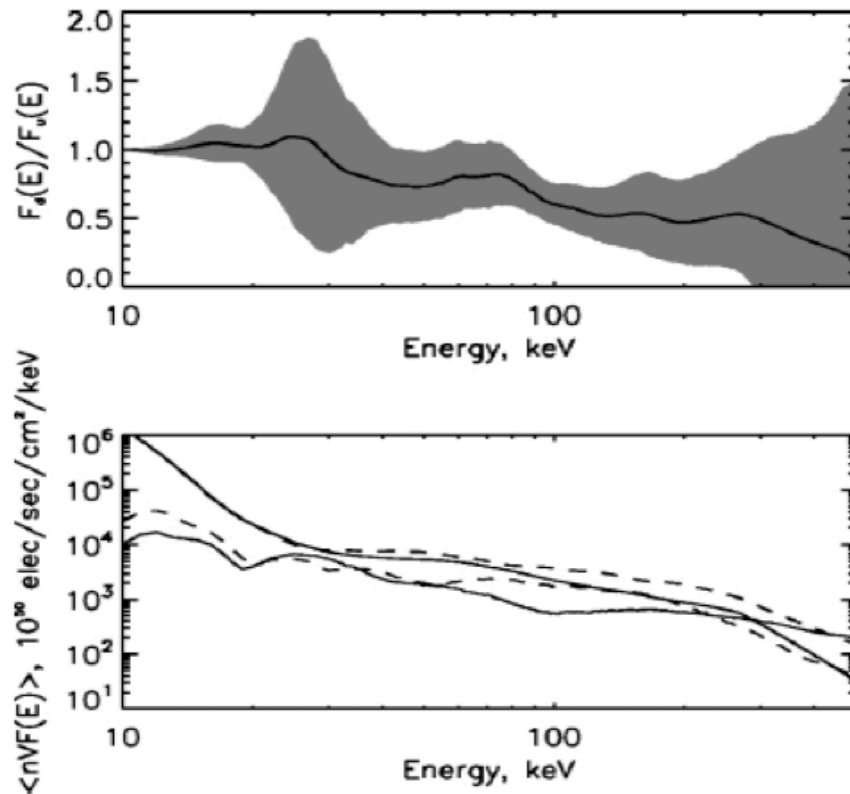


- Model independent method
- Direct inversion too contaminated by noise
- Use Tikhonov Regularisation
- Constraint is that electron spectrum is differentiable

$$\mathbf{I} = \mathbf{M}\mathbf{F}$$

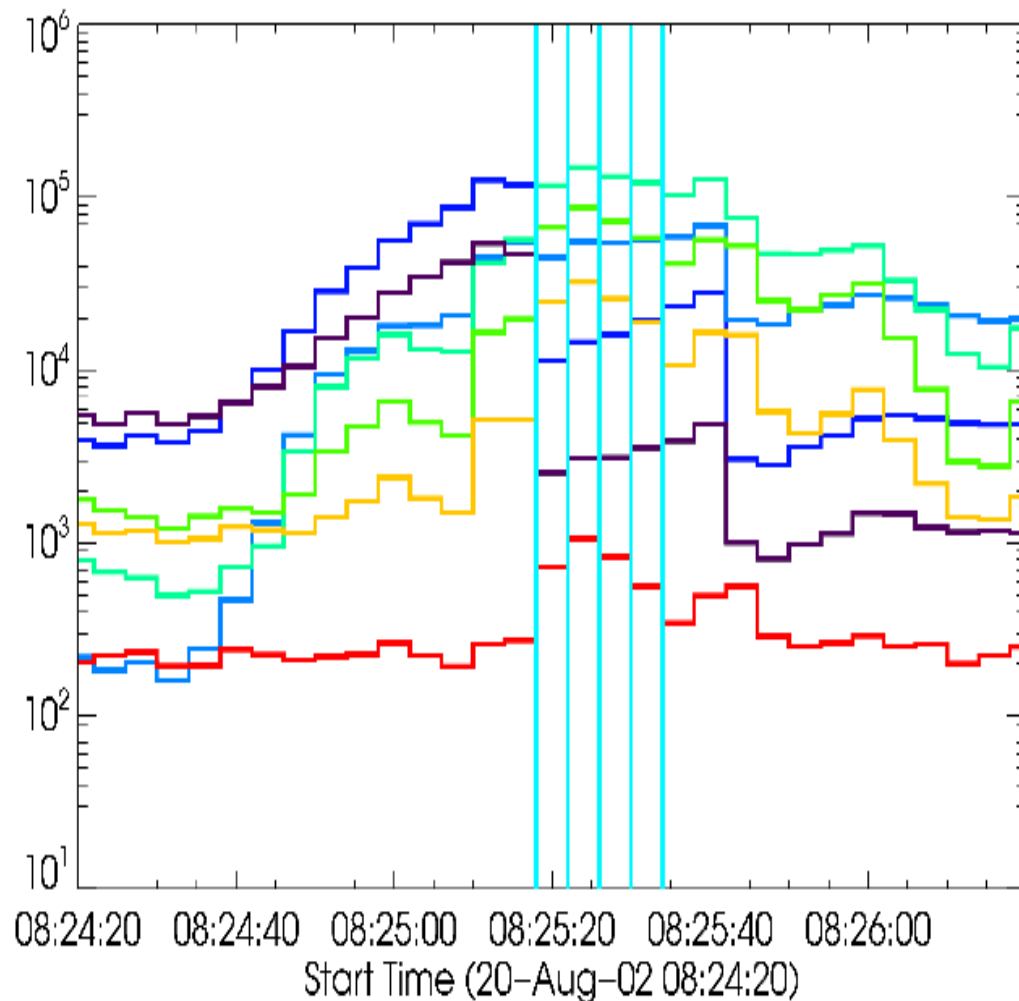
$$\|\mathbf{M}\bar{\mathbf{F}} - \mathbf{I}\|^2 = \min$$

$$\|\mathbf{M}\bar{\mathbf{F}} - \mathbf{I}\|^2 + \lambda\|\mathbf{L}\mathbf{F}\|^2 = \min$$

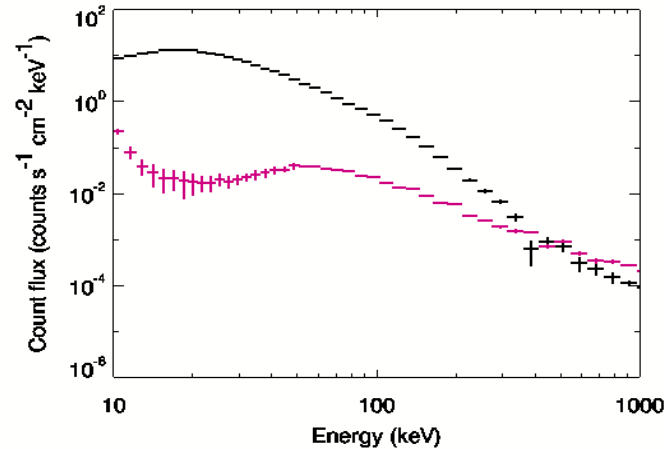
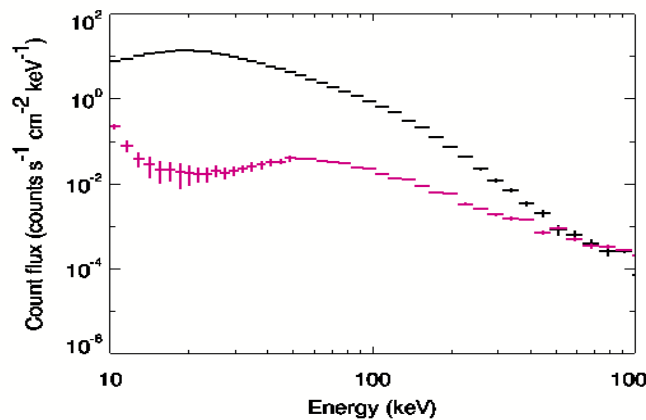
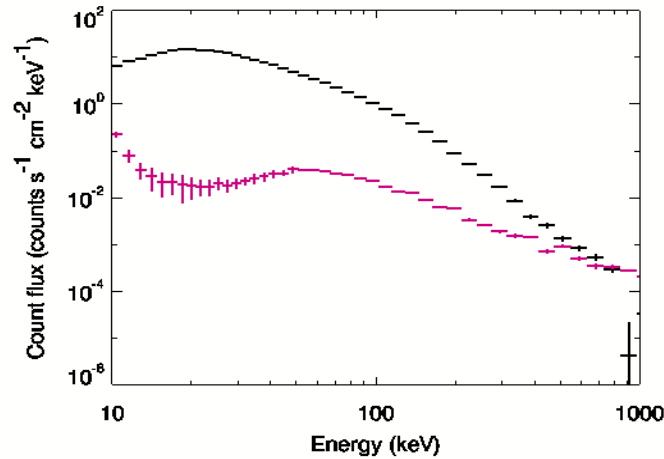
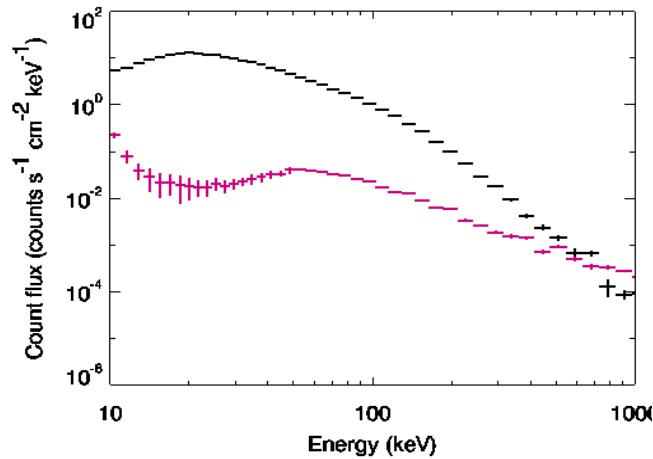


- This method was previously employed by Kontar et. Al on 2 fares detected by RHESSI
- Results suggested a distribution consistent with isotropic

Top : Plot of Anisotropy (F_{down}/F_{up}) against electron energy for flares full impulsive phase of the flare occurring on 20 Aug 2002
Bottom: Confidence bands for the two component electron flux (Kontar and Brown 2006)



- Is there any variation in anisotropy over the impulsive phase?
- Suitable flares need good statistics at high energies and be close to disk centre
- Four second time intervals over the impulsive peaks were studied



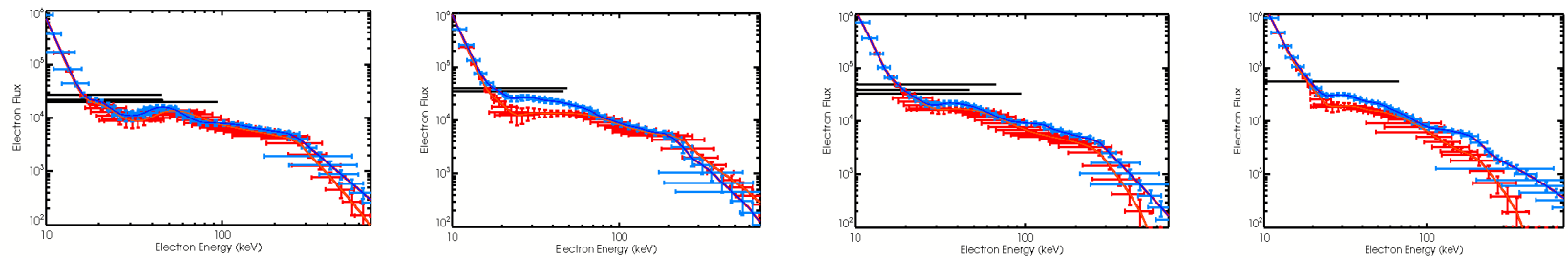
Plots of Count Flux against energy for each 4 second interval (black) with associated background (magenta)

- Photon counts accumulated over the impulsive phase.
- Pseudo logarithmic energy bins used.
- Energy range used – from 10 keV to maximum where counts are 3 sigma above background

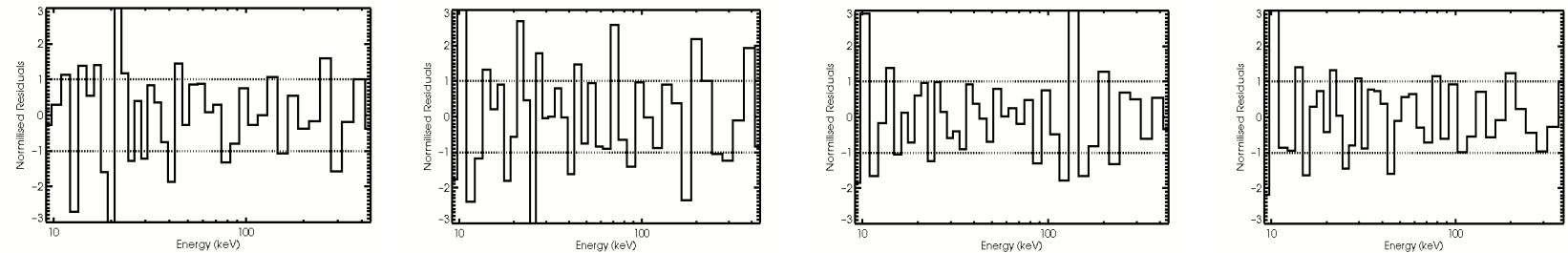


Regularized Inversion performed on count spectra for each time interval to determine 2D Electron spectra

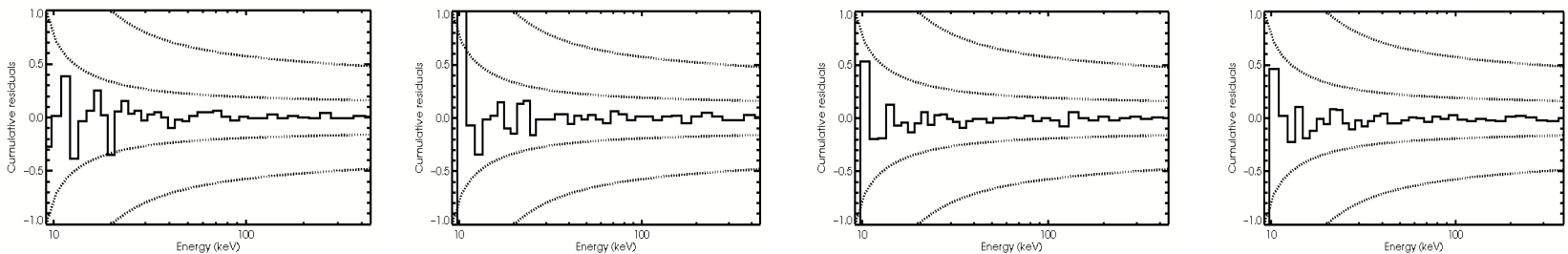
Electron Flux

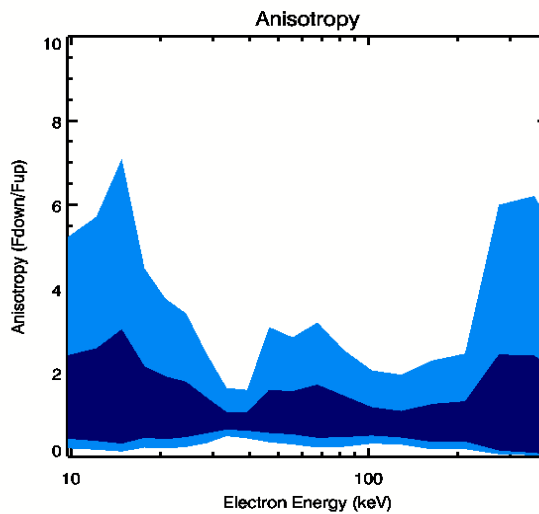
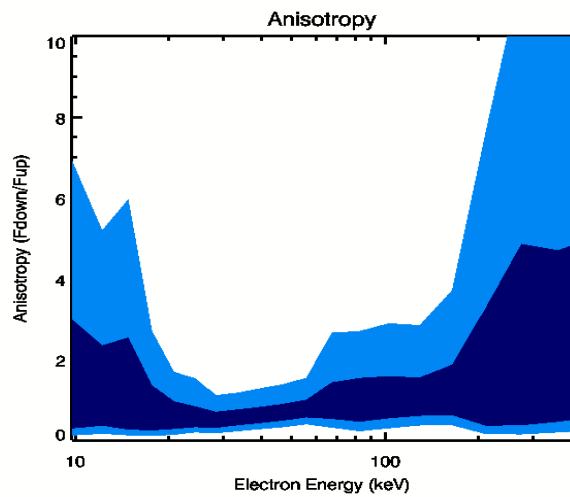
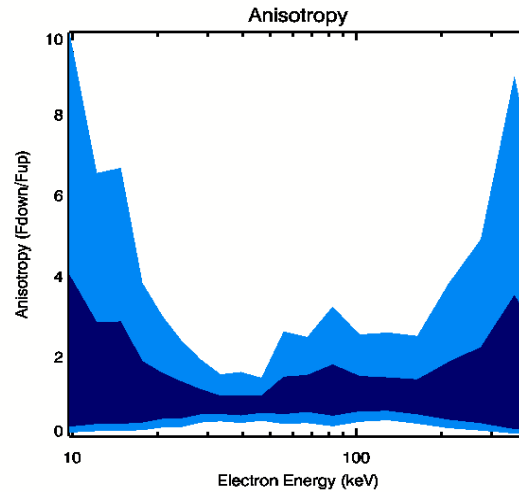
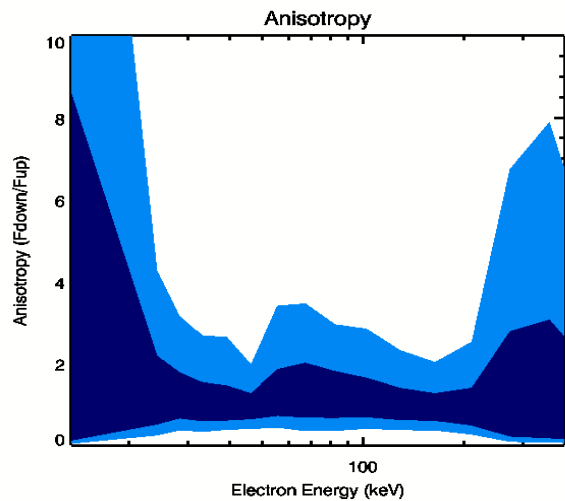


Normalized Residuals

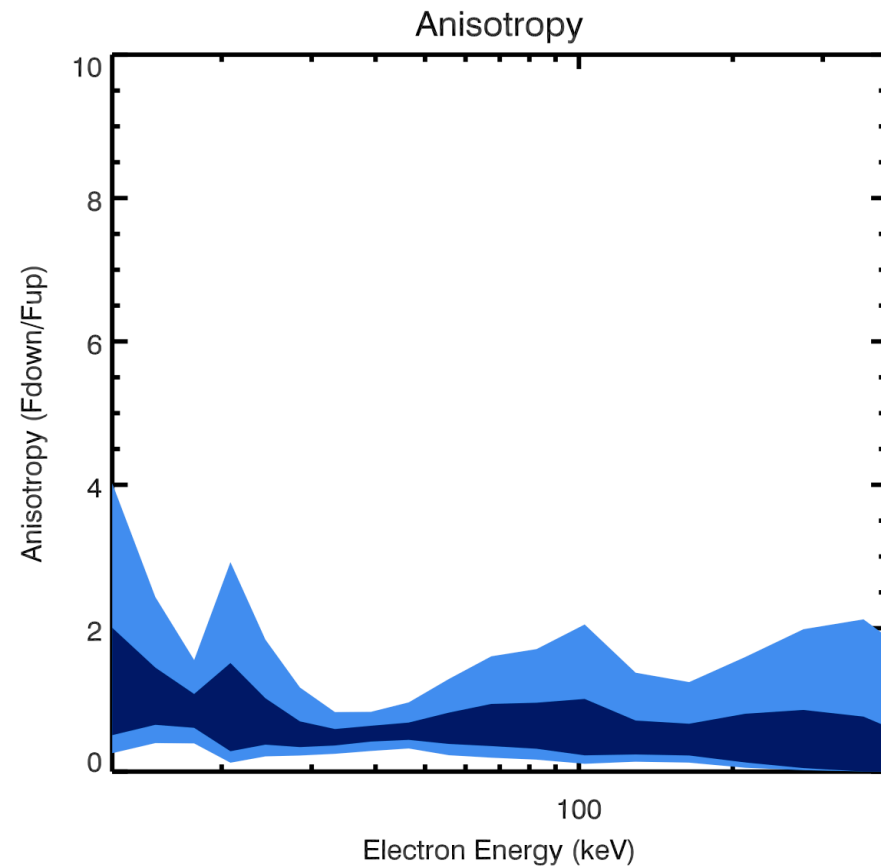
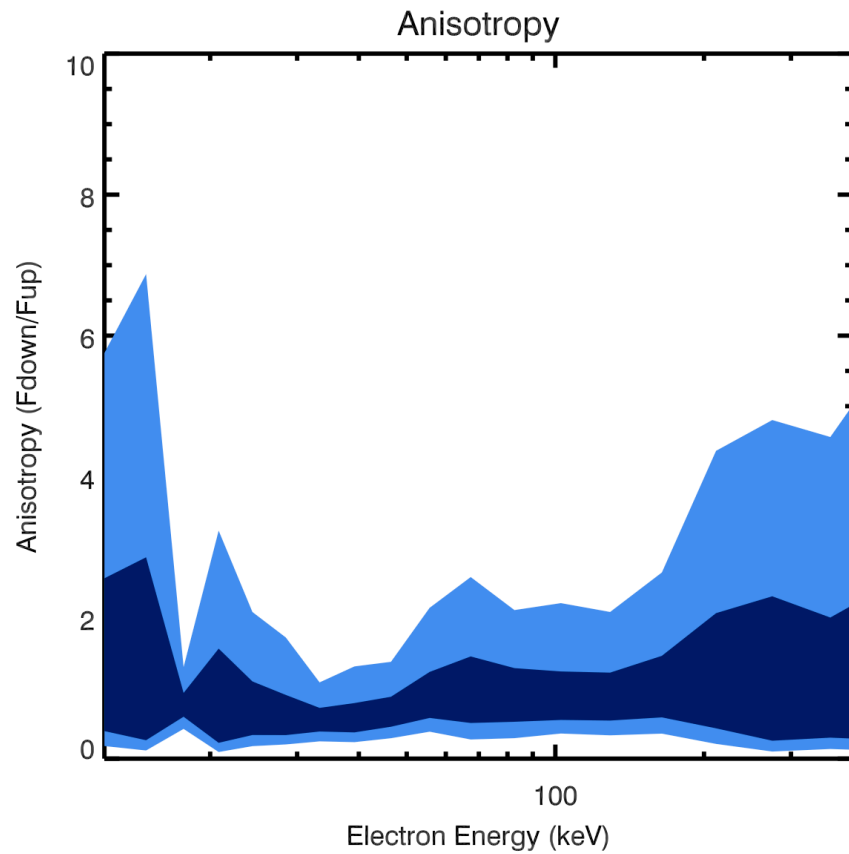


Cumulative Residuals





Confidence intervals – Anisotropy ($F_{\text{down}} / F_{\text{up}}$) against electron energy calculated using error estimates on electron spectra



- Longer time intervals give better statistics but changes in anisotropy over shorter timescales wont be apparent.



- This is consistent with studies which suggest that the electron spectrum is isotropic at low energies
- No evidence of variation on 4 second timescales.
- Where counts are strong enough could be possible to measure anisotropy for 2 second time intervals.